

Lattice Calculation of Neutron and Proton EDMs

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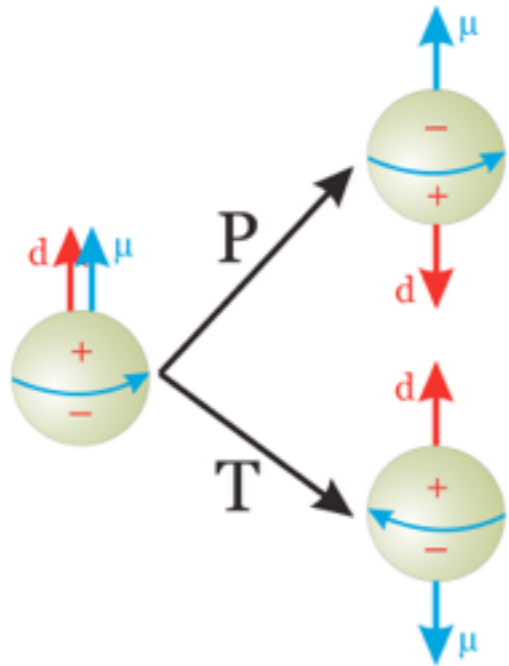


Symmetry Tests in Nuclei and Atoms
Kavli Institute for Theoretical Physics,
Santa Barbara, Sep 19-23, 2016

Outline

- Lattice basics
- nEDM induced by θ -term
- nEDM induced by quark chromo-EDM
- EDM in Background Electric Field

Neutron and Proton EDMs from quark-gluon CPv



$$\vec{d}_N = d_N \frac{\vec{S}}{S}$$

$$\mathcal{H} = -\vec{d}_N \cdot \vec{E}$$

Motivations to search for new CP-odd interactions

- Extensions of SM
- Required for baryogenesis
- Strong CP problem

Lattice QCD : connect quark/gluon-level effective operators to hadron/nuclei matrix elements and interactions

$$\mathcal{L}_{eff} = \sum_n \frac{c_n}{\Lambda^{d_n-4}} \mathcal{O}_n^{(d_n)}$$

$$\begin{cases} \mathcal{L}^{(4)} &= \theta \frac{g^2}{32\pi^2} G\tilde{G} \\ \mathcal{L}^{(5)} &= \sum_q [d_q \bar{q}(F \cdot \sigma)\gamma_5 q + \tilde{d}_q \bar{q}(G \cdot \sigma)\gamma_5 q] \\ \dots & \end{cases}$$



$$\begin{aligned} & d_{n,p} \\ & F_3^{n,p}(Q^2) \end{aligned}$$

Hadron Structure in Lattice QCD

Lattice Field Theory \Leftrightarrow Numerical evaluation of the Path Integral

$$\langle q_x \bar{q}_y \dots \rangle = \int \mathcal{D}(\text{Glue}) \int \mathcal{D}(\text{Quarks}) e^{-S_{\text{Glue}} - \bar{q}(\not{D} + m)q} [q_x \bar{q}_y \dots]$$

Grassmann integration

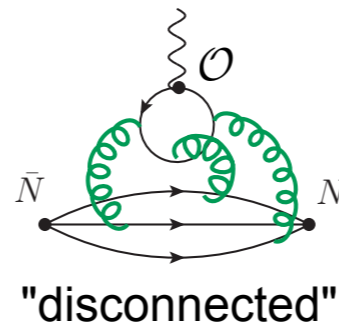
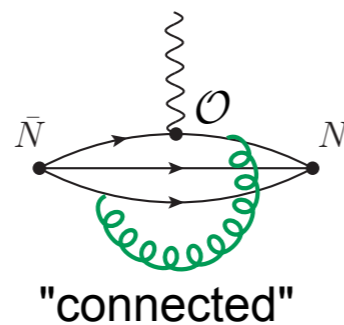
$$= \int \mathcal{D}(\text{Glue}) e^{-S_{\text{Glue}}} \text{Det}(\not{D} + m) [(\not{D} + m)^{-1}_{x,y} \dots]$$

Hybrid Monte Carlo

$(\not{D} + m) \cdot q = 0$
quark motion in
gluon background

Hadron Matrix Elements:

$$C_{3\text{pt}}^{\mathcal{O}}(T) = \langle N(T) \mathcal{O}(\tau) \bar{N}(0) \rangle =$$

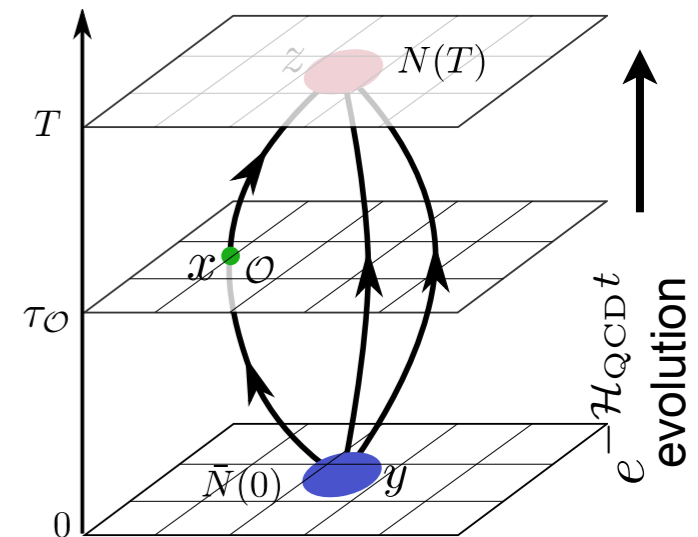


$$\langle N(T) \mathcal{O}(\tau) N(0) \rangle = \sum_{n,m} Z_m e^{-E_n(T-\tau)} \langle n | \mathcal{O} | m \rangle e^{-E_m \tau} Z_n^*$$

$$\xrightarrow{T \rightarrow \infty} Z_{00} e^{-M_N T} \left[\langle P' | \mathcal{O} | P \rangle + \underbrace{\mathcal{O}(e^{-\Delta E_{10} T}, e^{-\Delta E_{10} \tau}, e^{-\Delta E_{10}(T-\tau)})}_{\text{excited states}} \right]$$

Ground state
form factors

*Excited states contribute to correlators
and may (and do) bias results*



Each quark line = $(\not{D} + m)^{-1} \cdot \psi$

CP-odd Interaction on a Lattice

- Linearizing in CP-odd interaction, e.g. with θ -term

$$e^{-S_{QCD} - i\theta Q} = e^{-S_{QCD}} [1 - i\theta Q + O(\theta^2)]$$

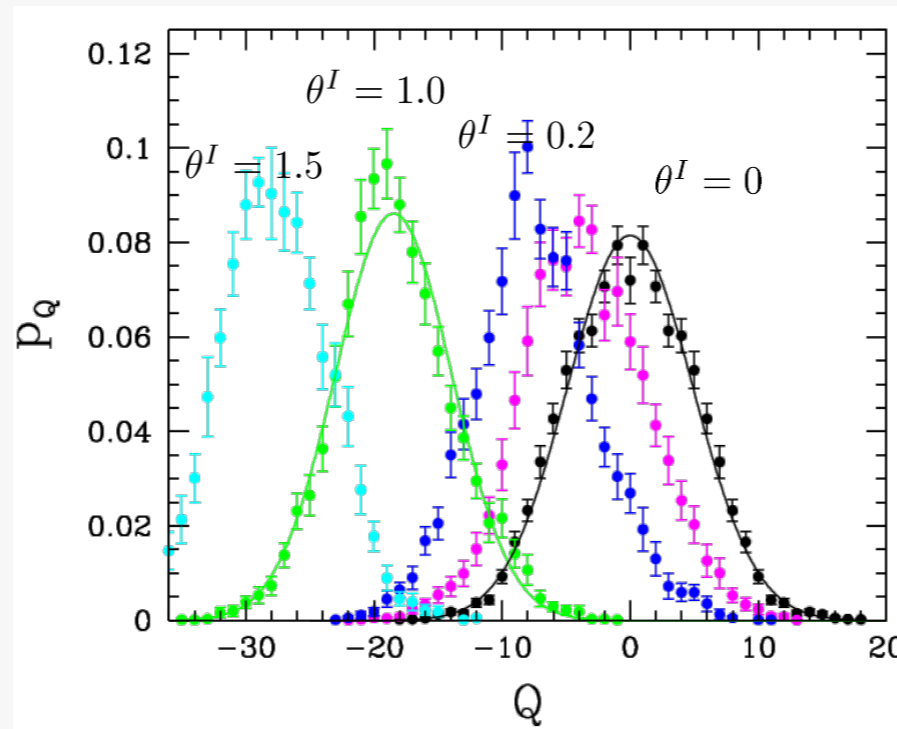
$$\langle \mathcal{O} \dots \rangle_{CP} = \langle \mathcal{O} \dots \rangle_{CP-even} - i\theta \langle Q \cdot \mathcal{O} \dots \rangle_{CP-even} + O(\theta^2)$$

- Simulating with CP-odd term(s)

$$\langle \mathcal{O} \dots \rangle_{\theta} \sim \int \mathcal{D}U e^{-S - \theta^I Q} (\mathcal{O} \dots)$$

continued to Imag. θ
to avoid sign problem

[T.Izubuchi et al (2007) ;
R.Horsley et al (2008) ;
F.K.Guo et al (2015)]

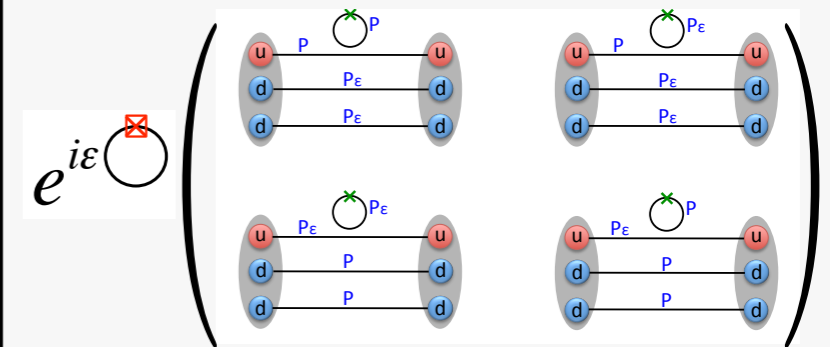


- + better sampling of $Q \neq 0$
- linearity needs check
- need new ensemble

- Reweighting with CP-odd term(s)

$$\bar{q} [\not{D} + m_q + i\epsilon(G \cdot \sigma)\gamma_5] q$$

quark operator with cEDM
[T.Bhattacharya et al(LANL)]



EDM from Spectrum vs. Form Factors

- Nucleon spectrum in the background electric field [S.Aoki et al '89 ; E.Shintani et al '06]

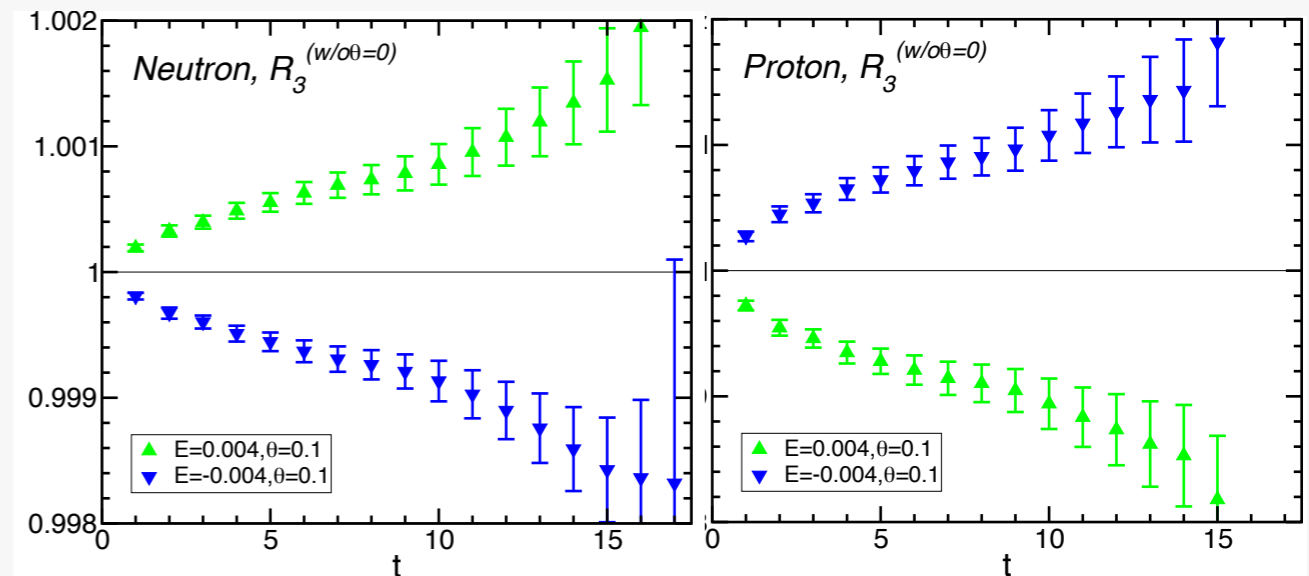
$$\langle N(t)\bar{N}(0) \rangle_{\theta, \vec{E}} \sim e^{-(E \pm \vec{d}_N \cdot \vec{E})t}$$

$$\frac{\langle N_{\uparrow}(t)\bar{N}_{\uparrow}(0) \rangle_{\theta, E_z}}{\langle N_{\downarrow}(t)\bar{N}_{\downarrow}(0) \rangle_{\theta, E_z}} \sim e^{2d_N E_z t} \approx 1 + 2d_N E_z t$$

Wick rotation: $\vec{E} \rightarrow i\vec{E}$

SU(3) g.f. link $U_z \rightarrow U_z e^{iE_z t} \sim e^{-E_z t}$

non-periodic with real(Minkowski) E_z



[E.Shintani et al, PRD75, 034507(2007)]

- P,T-odd Form Factor $d_N = F_3(0)/2m$

[E.Shintani et al '05, '15 ; F.Berruto et al '05 ; A.Shindler et al '15 ; C.Alexandrou et al '15]

$$\langle N | J^\mu | \bar{N} \rangle_{\mathcal{CP}} = \bar{u} \Gamma_{CP-even}^\mu u + \bar{u} \Gamma_{CP-odd}^\mu u$$

$$F_1 \gamma^\mu + F_2 \frac{i\sigma^{\mu\nu} q_\nu}{2m} \quad \hat{F}_3 \frac{\gamma_5 \sigma^{\mu\nu} q_\nu}{2m}$$

Need either extrapolation $F_3(Q^2 \rightarrow 0)$,
or smart tricks [C.Alexandrou's talk]

Nucleon spinors are parity-mixed

$$\langle N(t)\bar{N}(0) \rangle_{\mathcal{CP}} \sim \frac{-i\not{p} + me^{2i\alpha_N \gamma_5}}{2m_N} e^{-E_N t}$$

CP-odd matrix elements require subtraction of $F_{1,2}$ contributions:

$$\langle Q \cdot N J^\mu \bar{N} \rangle \sim \mathcal{K}_3^\mu F_3 + \alpha(\mathcal{K}_1^\mu F_1 + \mathcal{K}_2^\mu F_2)$$

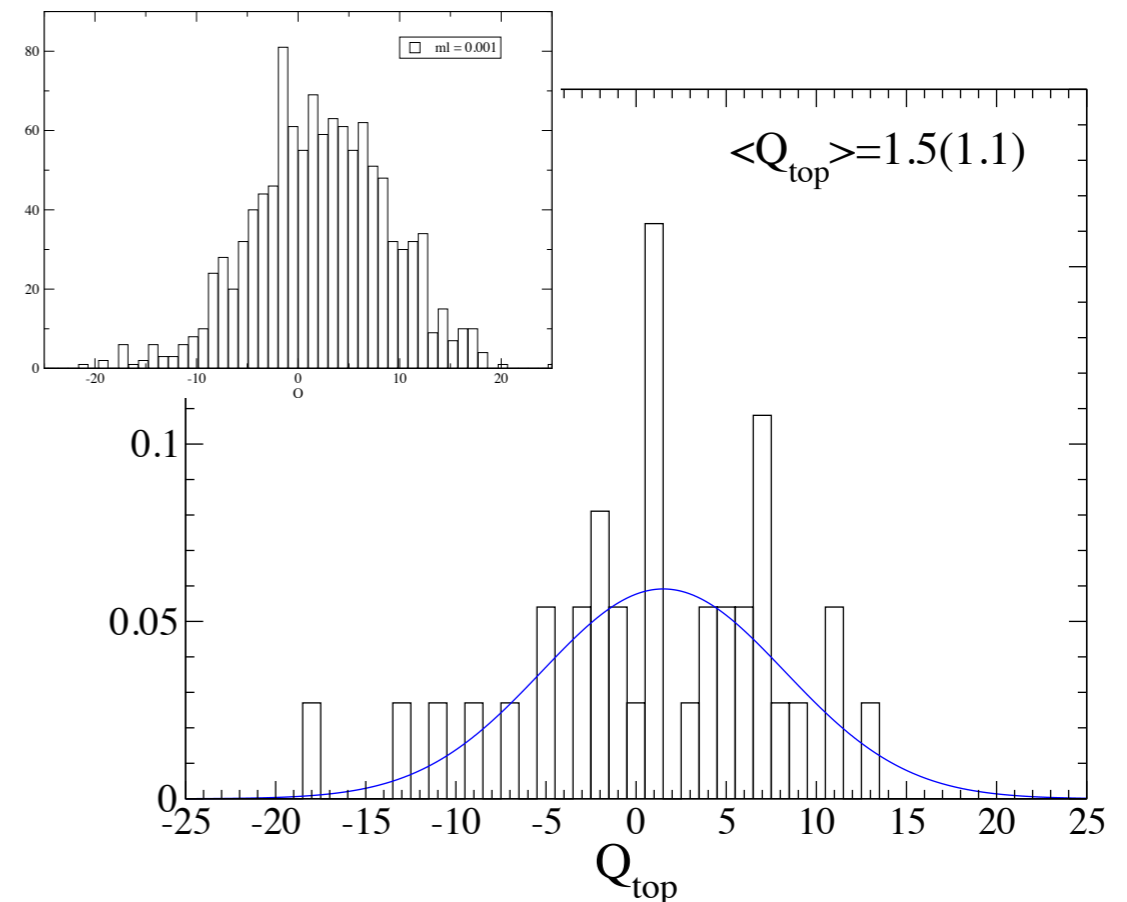
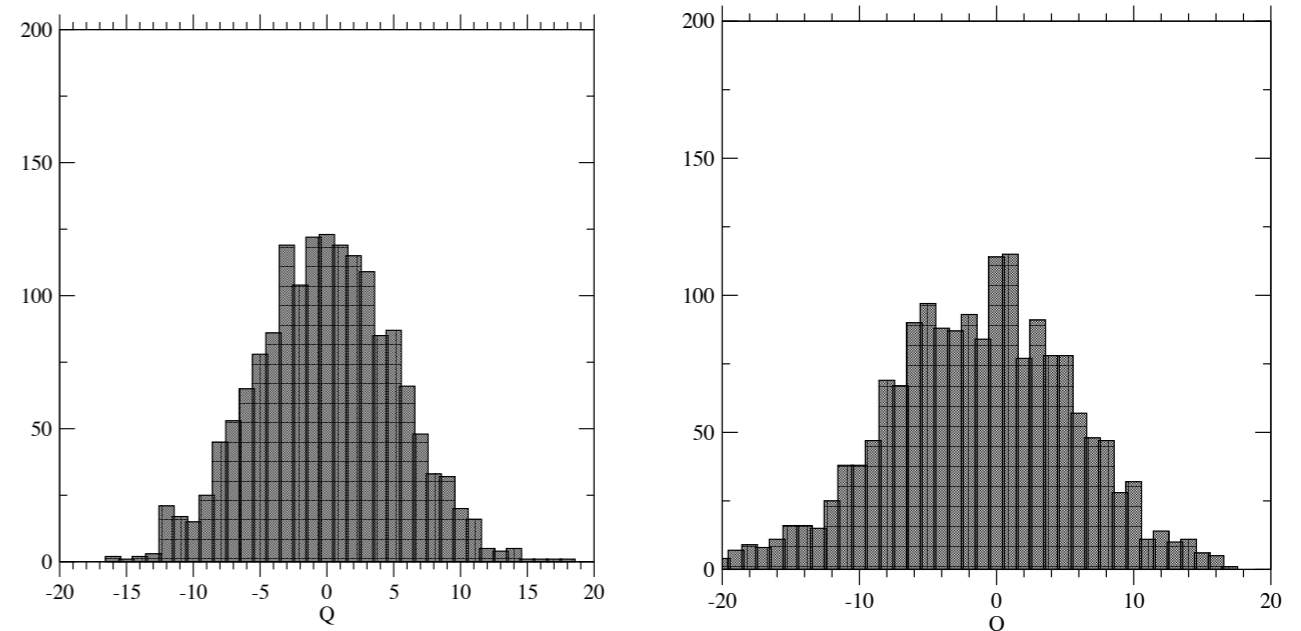
Calculation with Chirally-Symmetric Quarks

[E.Shintani, T.Blum, T.Izubuchi,
A.Soni, PRD93, 094503(2015)]

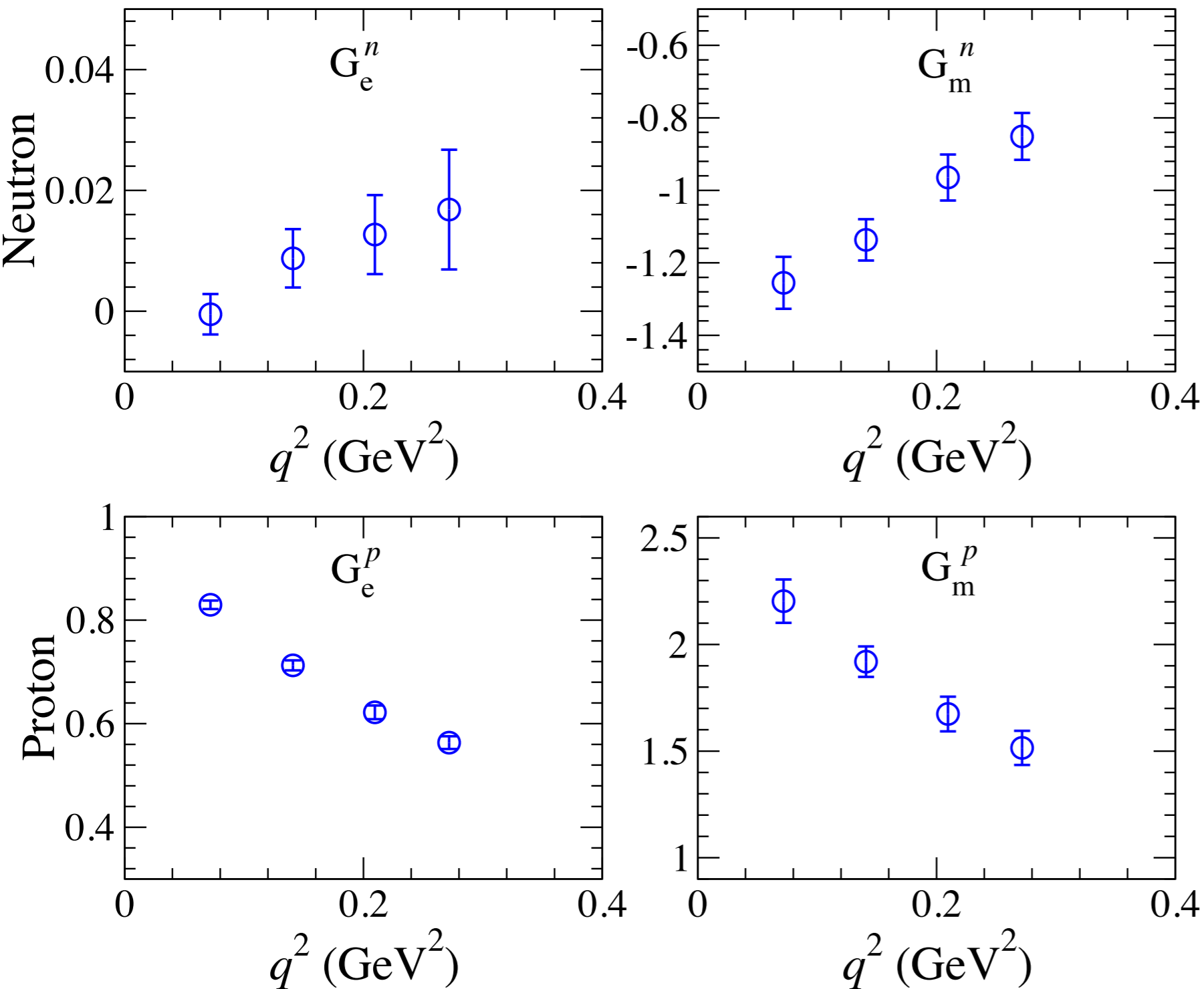
- $1/a = 1.73 \text{ GeV}$
- $V=(2.7 \text{ fm})^3$
- $M_{\text{pi}} = 330, 400 \text{ MeV}$
- 750 configurations

- $1/a = 1.37 \text{ GeV}$
- $V=(4.6 \text{ fm})^3$
- $M_{\text{pi}} = 170 \text{ MeV}$
- 39 configurations

Topological charge



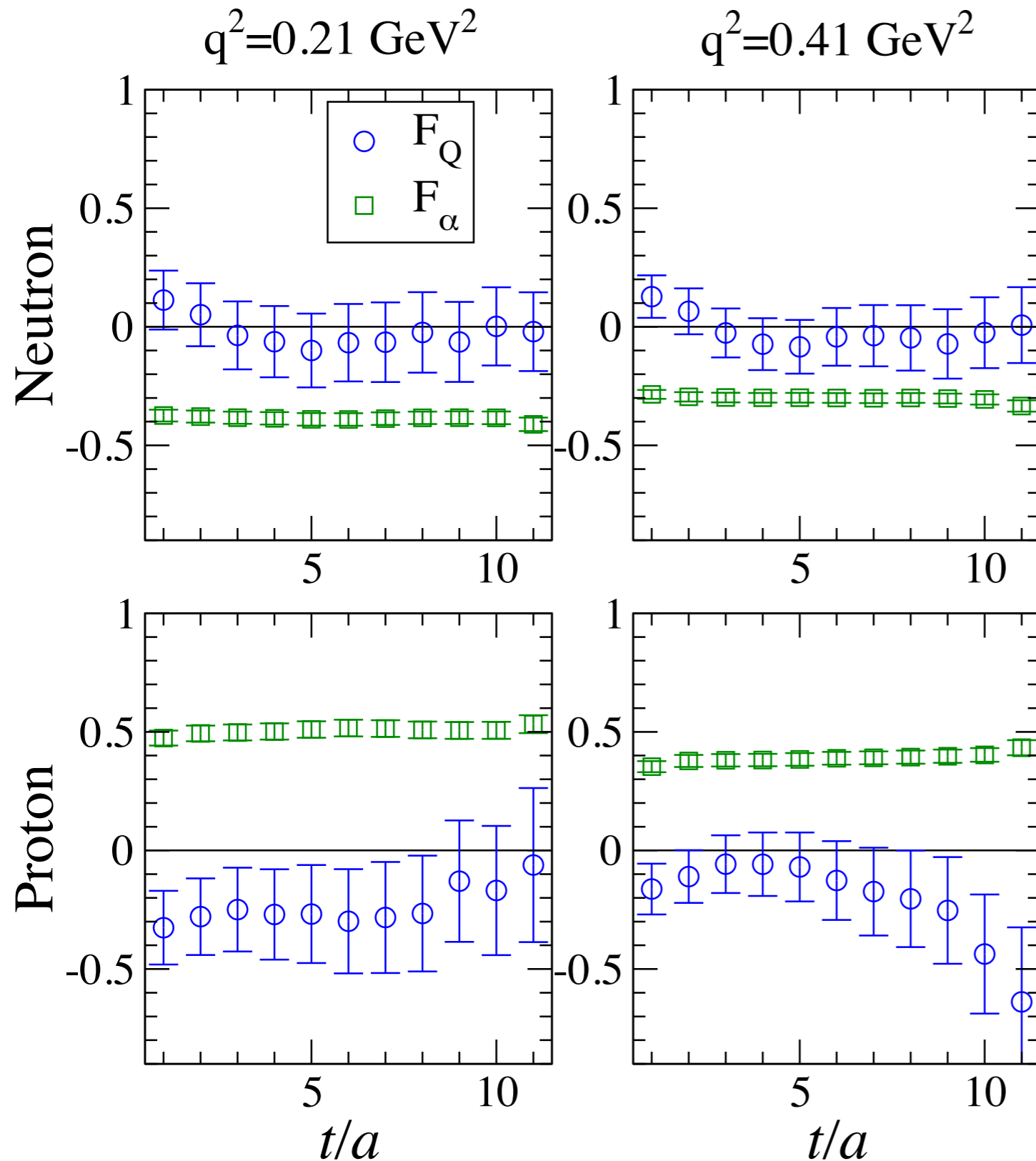
Electric and Magnetic Form Factors



[E.Shintani, T.Blum, T.Izubuchi, A.Soni, PRD93, 094503(2015)]

- $(4.6 \text{ fm})^3 \times (9.2 \text{ fm})$ box
- $m_\pi = 170 \text{ MeV}$

CP-odd Form Factors



[E.Shintani, T.Blum, T.Izubuchi, A.Soni, PRD93, 094503(2015)]

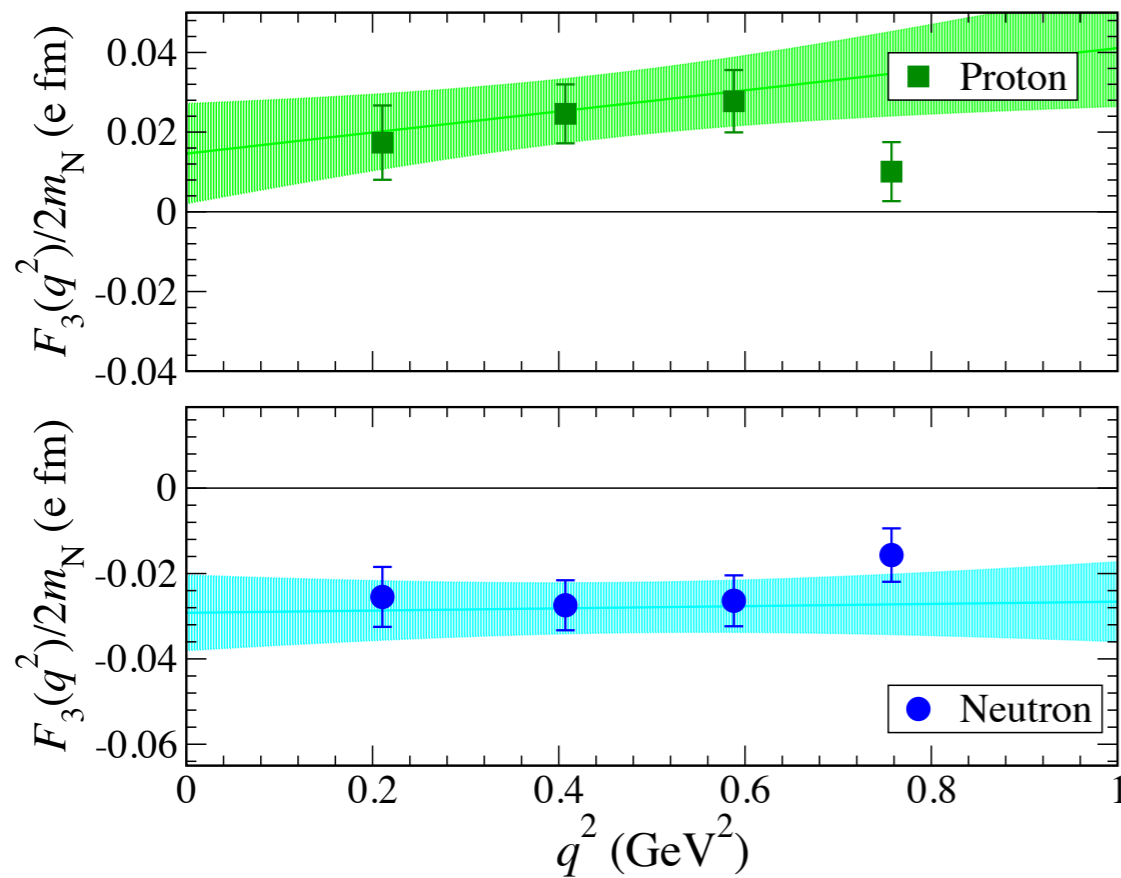
- $(2.7 \text{ fm})^3 \times (7.3 \text{ fm})$ box
- $m_\pi = 330 \text{ MeV}$

$$F_3 = F_Q + F_\alpha$$

\swarrow \searrow
 Lattice CP-mixing
 CP-odd f.f. correction

Q²-Dependence of F₃

- (2.7 fm)³x(7.3 fm) box
- m_π=330 MeV

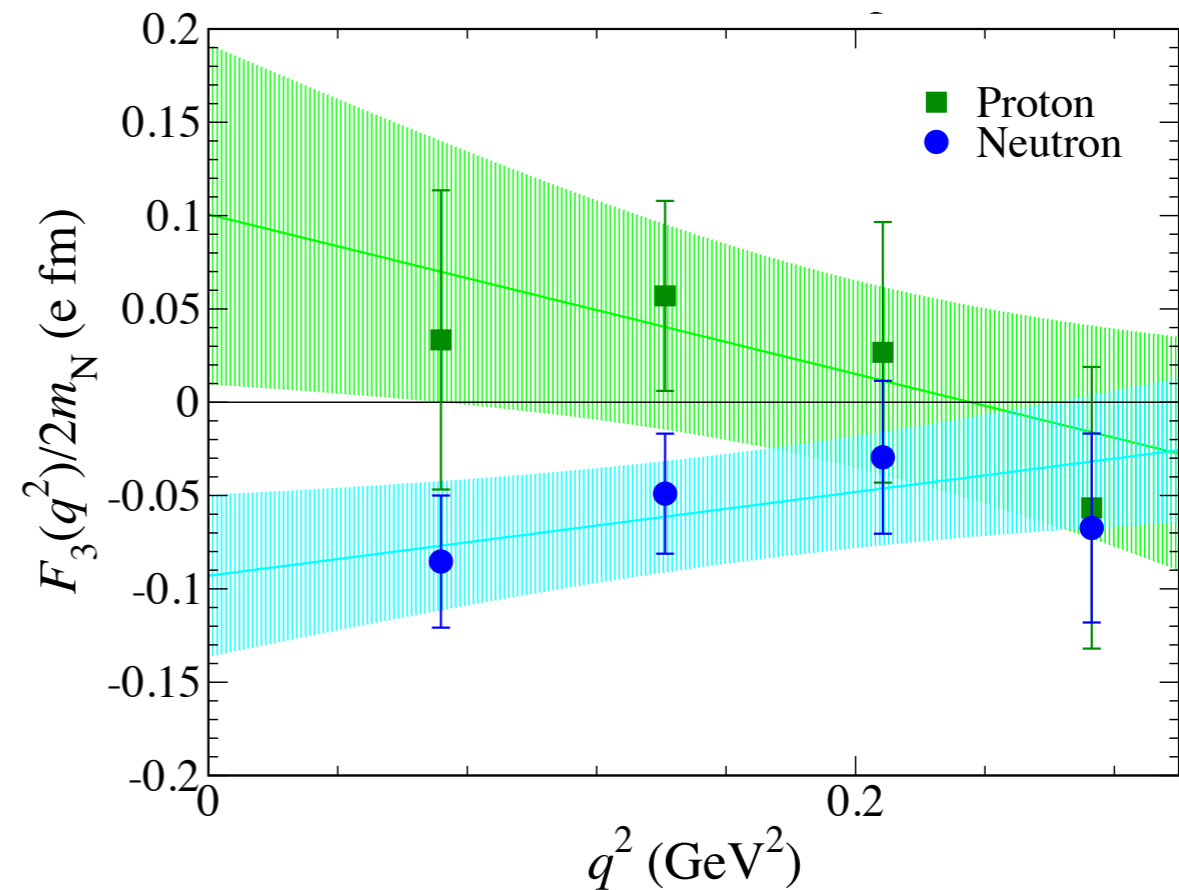


$$S'_p = -11(21) \cdot 10^{-4} e \cdot \text{fm}^3$$

$$S'_n = 24(14) \cdot 10^{-4} e \cdot \text{fm}^3$$

- Schiff moments from linear fit

- (4.6 fm)³x(9.2 fm) box
- m_π=170 MeV



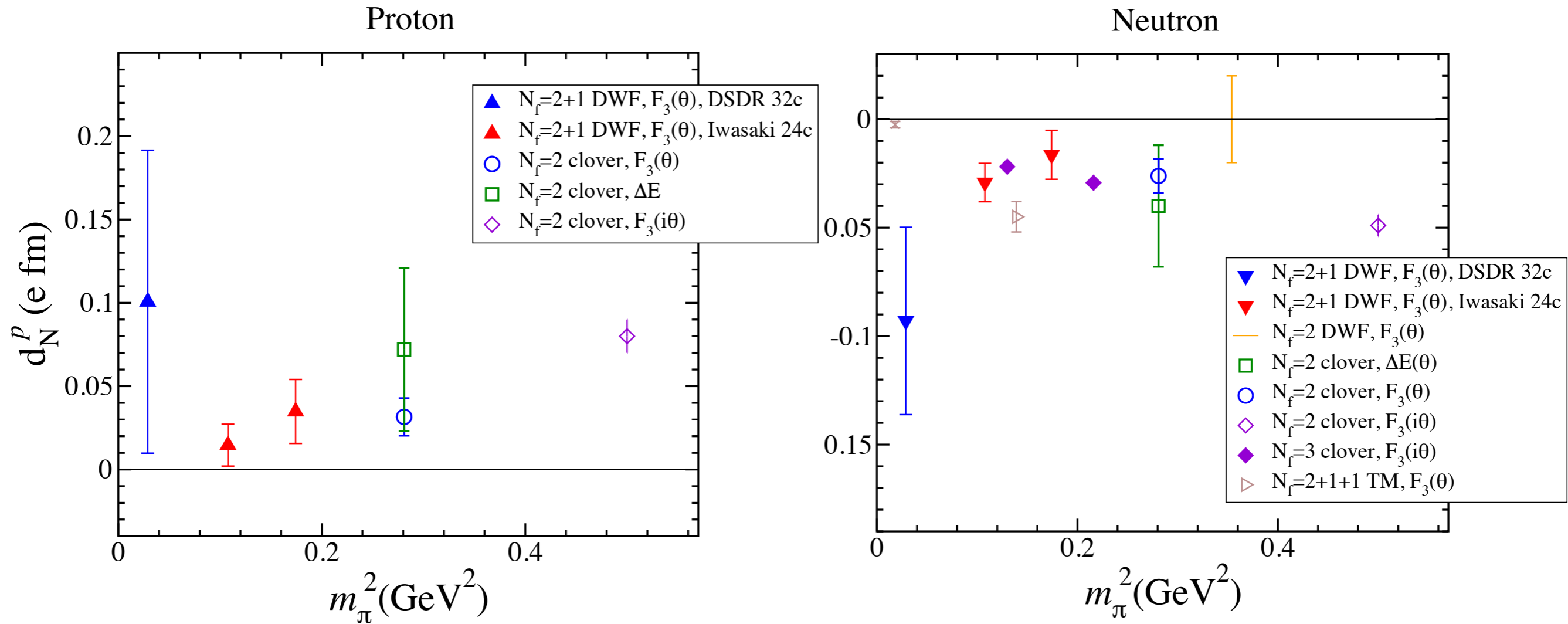
$$S'_p = -170(150) \cdot 10^{-4} e \cdot \text{fm}^3$$

$$S'_n = 87(94) \cdot 10^{-4} e \cdot \text{fm}^3$$

$$\frac{1}{2m_N} F_3(Q^2) = d_N + S'Q^2 + O(Q^4)$$

[E.Shintani, T.Blum, T.Izubuchi,
A.Soni, PRD93, 094503(2015)]

EDM vs. Pion Mass



- Substantial MC noise due to extensive nature of top.charge

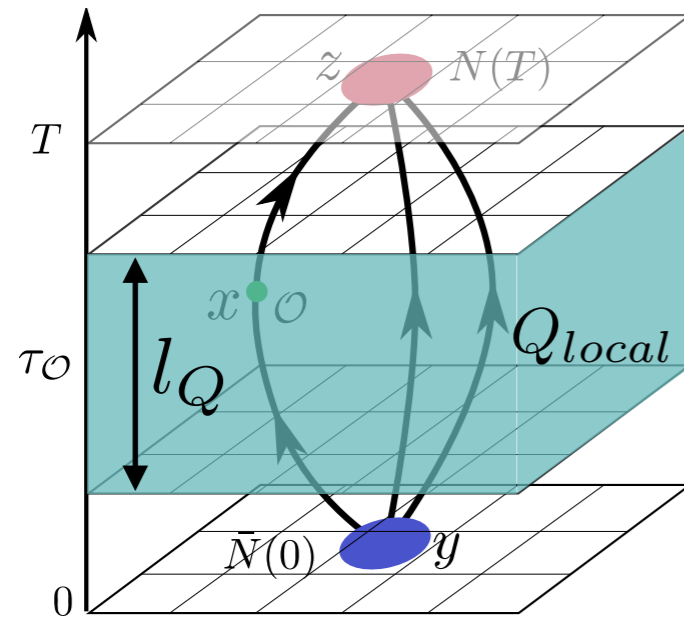
[E.Shintani, T.Blum, T.Izubuchi, A.Soni, PRD93, 094503(2015)]

Localized Sampling of $Q=\tilde{F}F$

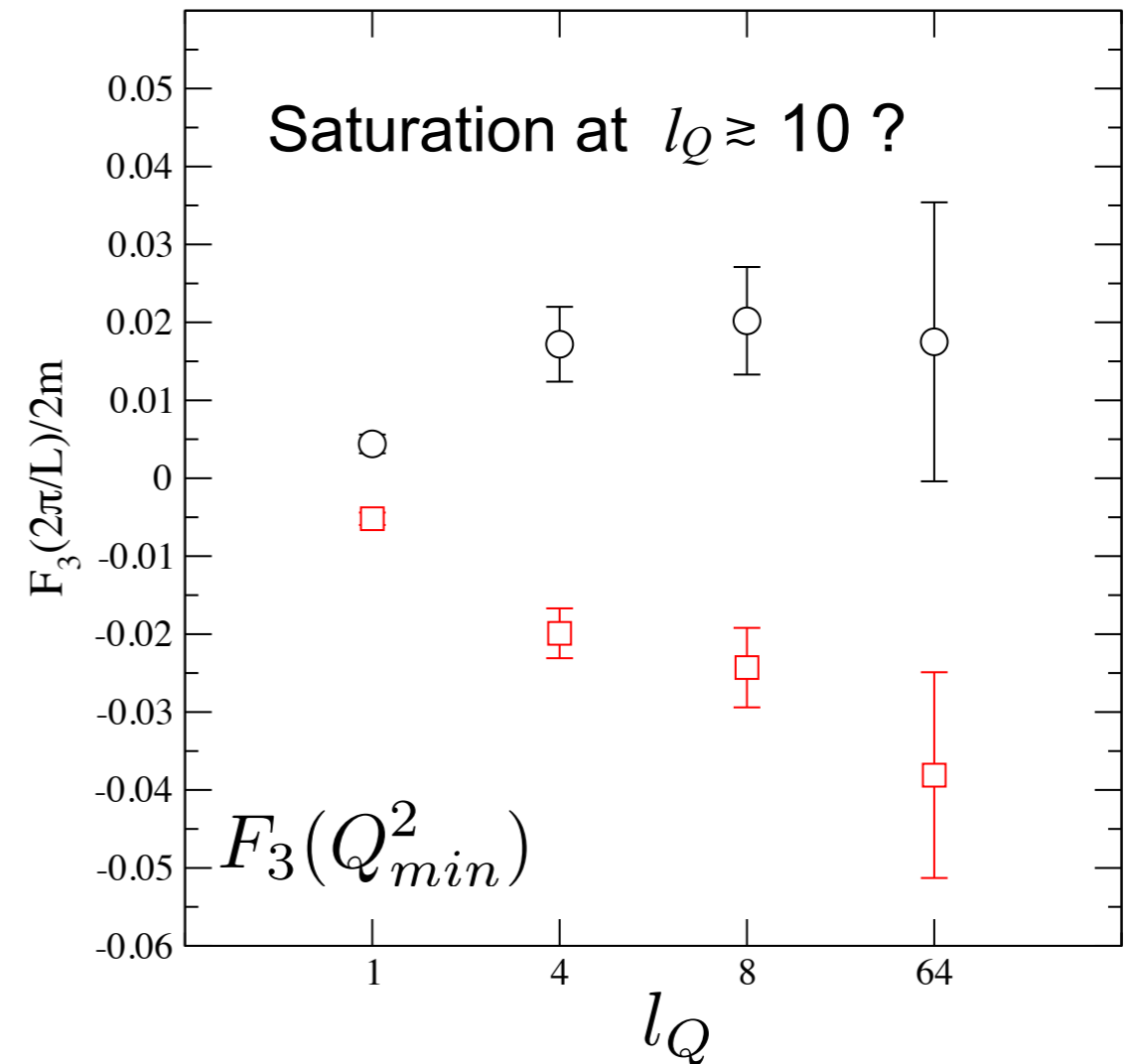
$$\langle \tilde{F}F(x)\tilde{F}F(0) \rangle \sim e^{-m_{\eta'}|x|}$$

[E.Shintani, T.Blum, T.Izubuchi, A.Soni, PRD93, 094503(2015)]

- Overcome noise from top.charge fluctuations: sample $\tilde{F}F$ locally



$$Q_{local}(\tau, l_Q) \sim \int_{\tau-l_Q/2}^{\tau+l_Q/2} dt dV \tilde{F}F$$



- $(2.7 \text{ fm})^3 \times (7.3 \text{ fm})$ box
- $m_\pi = 330 \text{ MeV}$

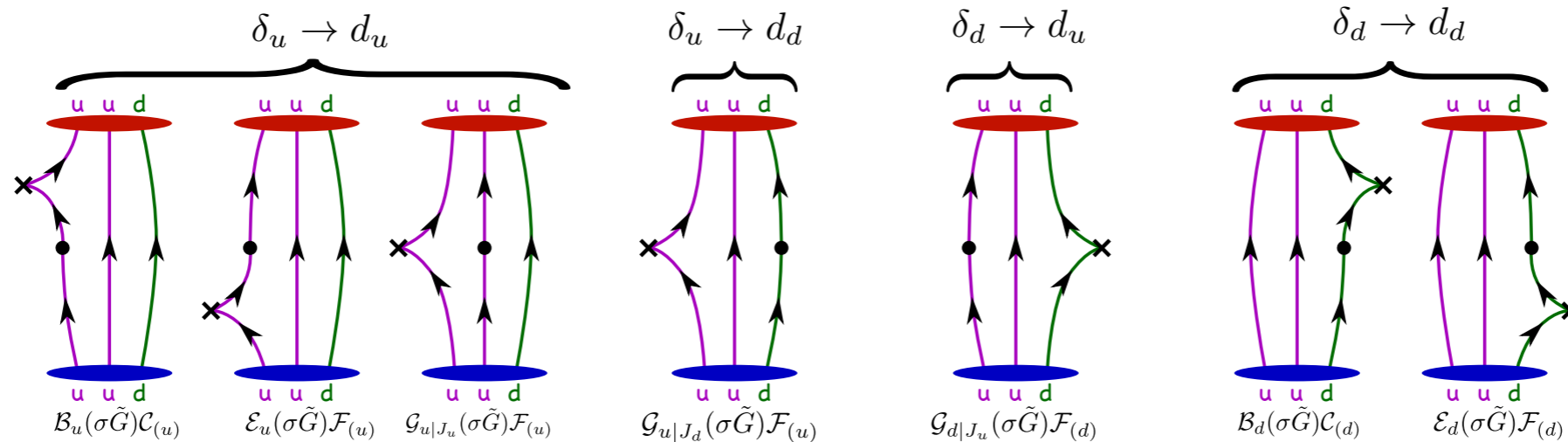
Quark Chromo-EDM

$$\mathcal{L}^{(5)} = \sum_q \tilde{d}_q \bar{q}(G \cdot \sigma) \gamma_5 q \quad \begin{array}{l} \nearrow \\ \searrow \end{array} \begin{array}{l} \langle N(y) \bar{N}(0) \int d^4x (\tilde{G} \cdot \sigma) \rangle \\ \langle N(y) [\bar{q} \gamma^\mu q](z) \bar{N}(0) \int d^4x (\tilde{G} \cdot \sigma) \rangle \end{array}$$

Quark-Gluon EDM: Insertions of dim-5 Operators

$$\mathcal{L}^{(5)} = \sum_q \tilde{d}_q \bar{q} (G \cdot \sigma) \gamma_5 q \quad \begin{array}{l} \nearrow \\ \searrow \end{array} \begin{array}{l} \langle N(y) \bar{N}(0) \int d^4x (\tilde{G} \cdot \sigma) \rangle \\ \langle N(y) [\bar{q} \gamma^\mu q](z) \bar{N}(0) \int d^4x (\tilde{G} \cdot \sigma) \rangle \end{array}$$

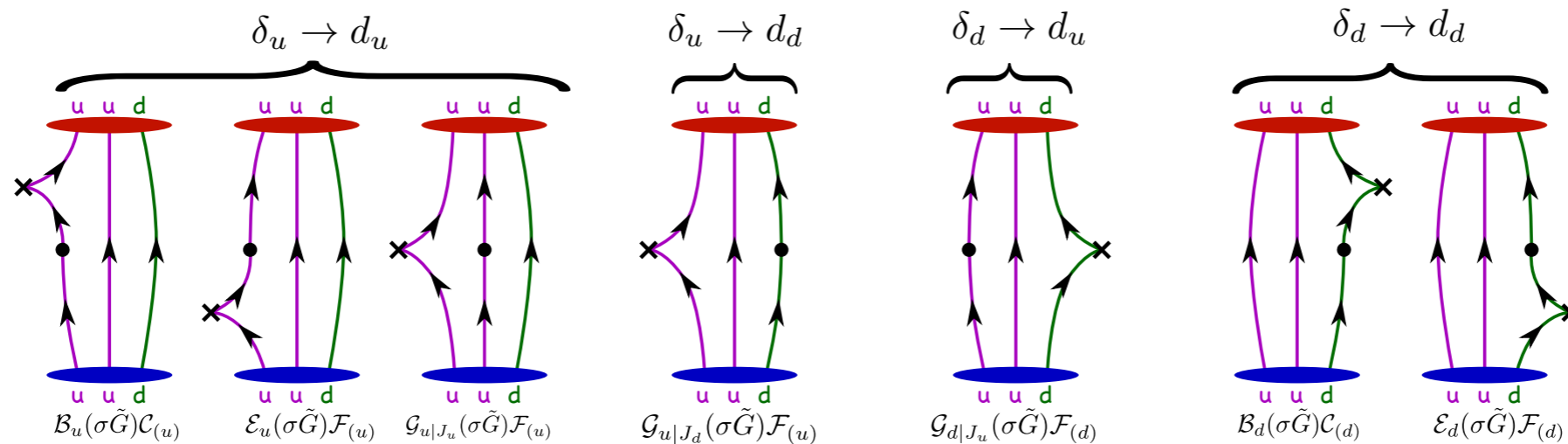
● Now: Only quark-connected insertions



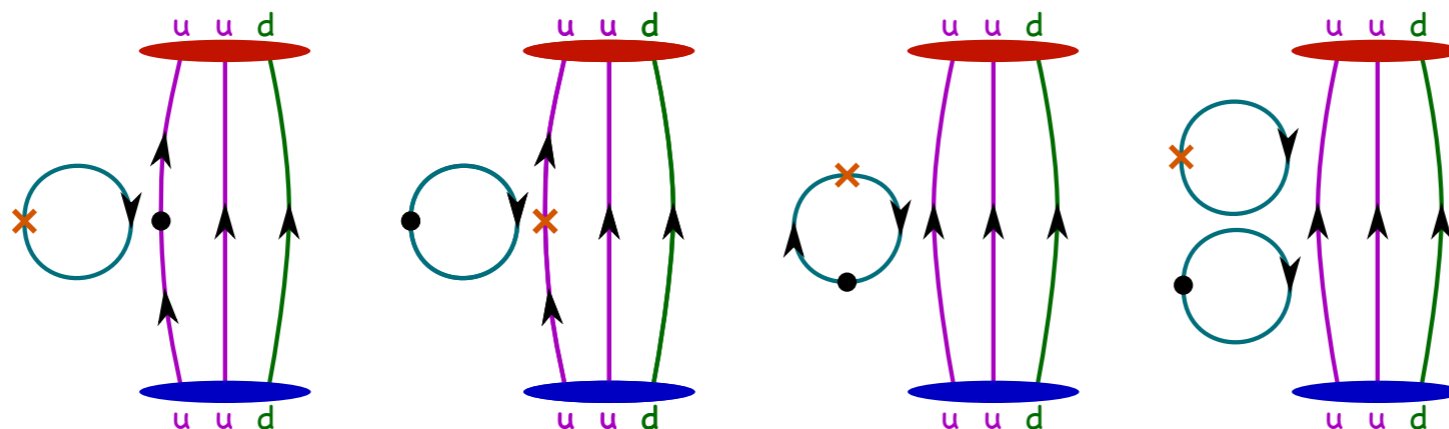
Quark-Gluon EDM: Insertions of dim-5 Operators

$$\mathcal{L}^{(5)} = \sum_q \tilde{d}_q \bar{q}(G \cdot \sigma)\gamma_5 q \quad \begin{matrix} \nearrow \\ \searrow \end{matrix} \begin{matrix} \langle N(y) \bar{N}(0) \int d^4x (\tilde{G} \cdot \sigma) \rangle \\ \langle N(y) [\bar{q}\gamma^\mu q](z) \bar{N}(0) \int d^4x (\tilde{G} \cdot \sigma) \rangle \end{matrix}$$

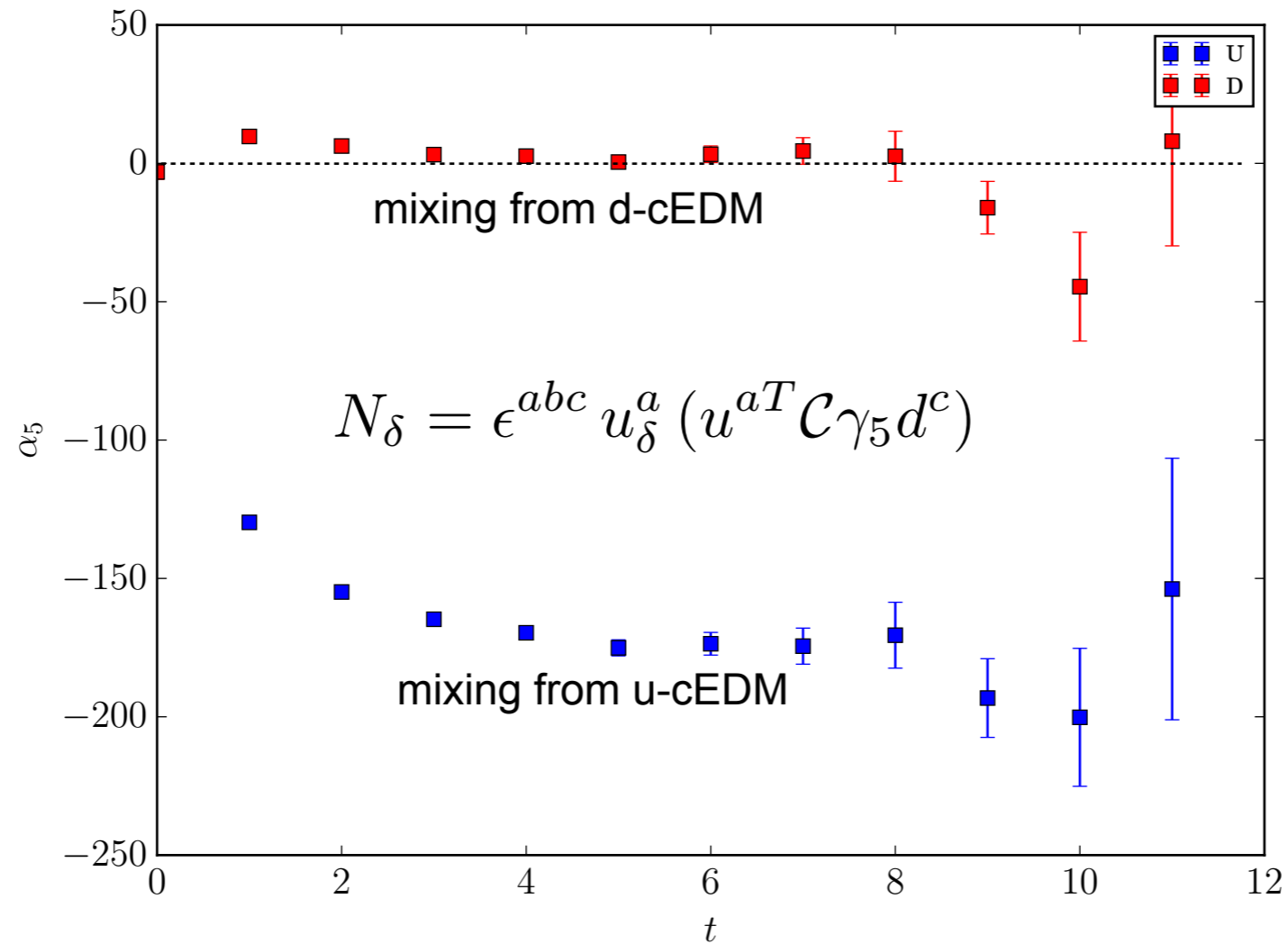
- Now: Only quark-connected insertions



- Some day: Single- and double-disconnected diagrams (contribute to isosinglet cEDM, mix with θ -term)



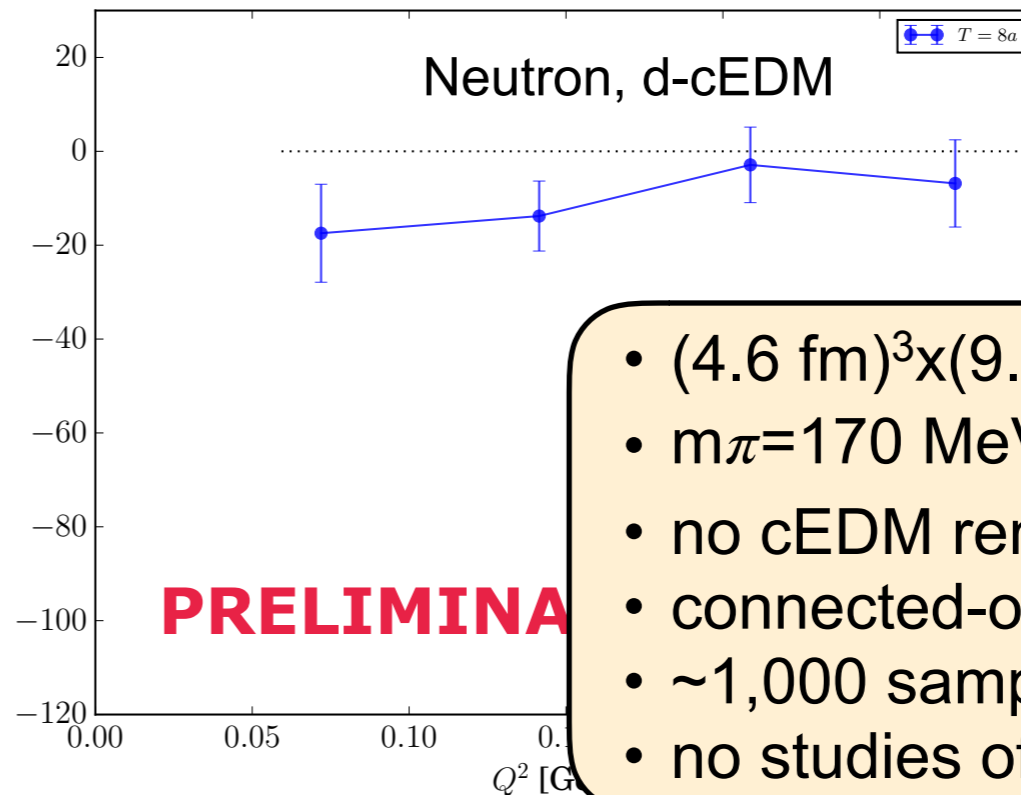
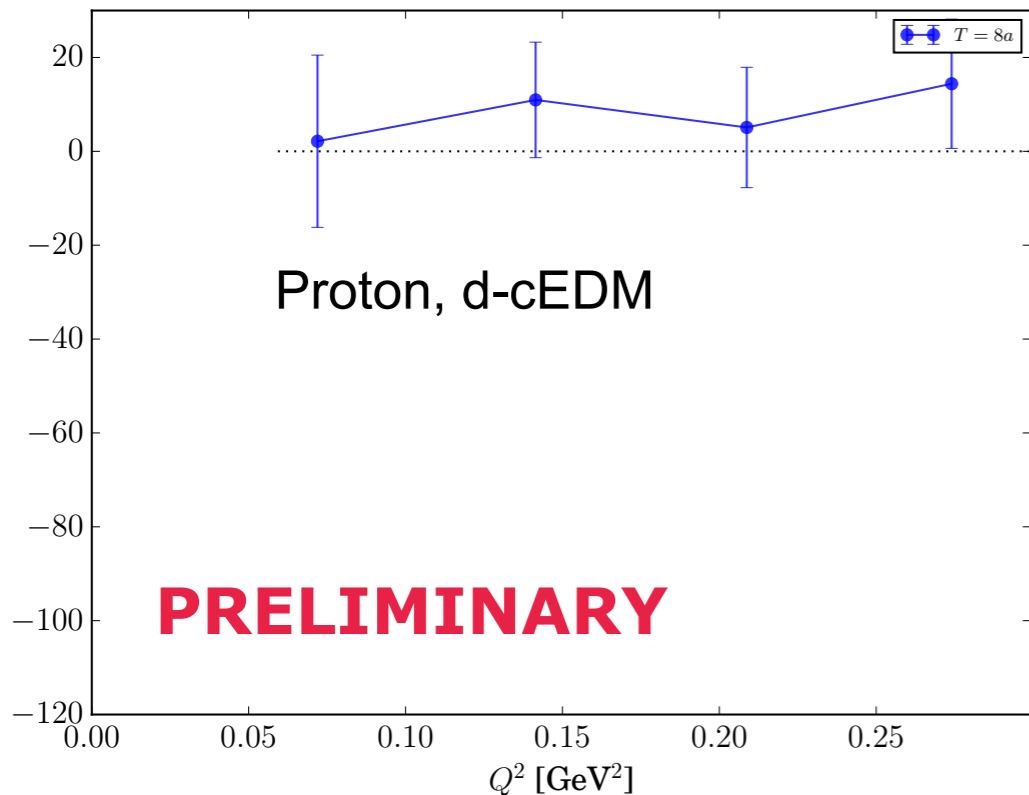
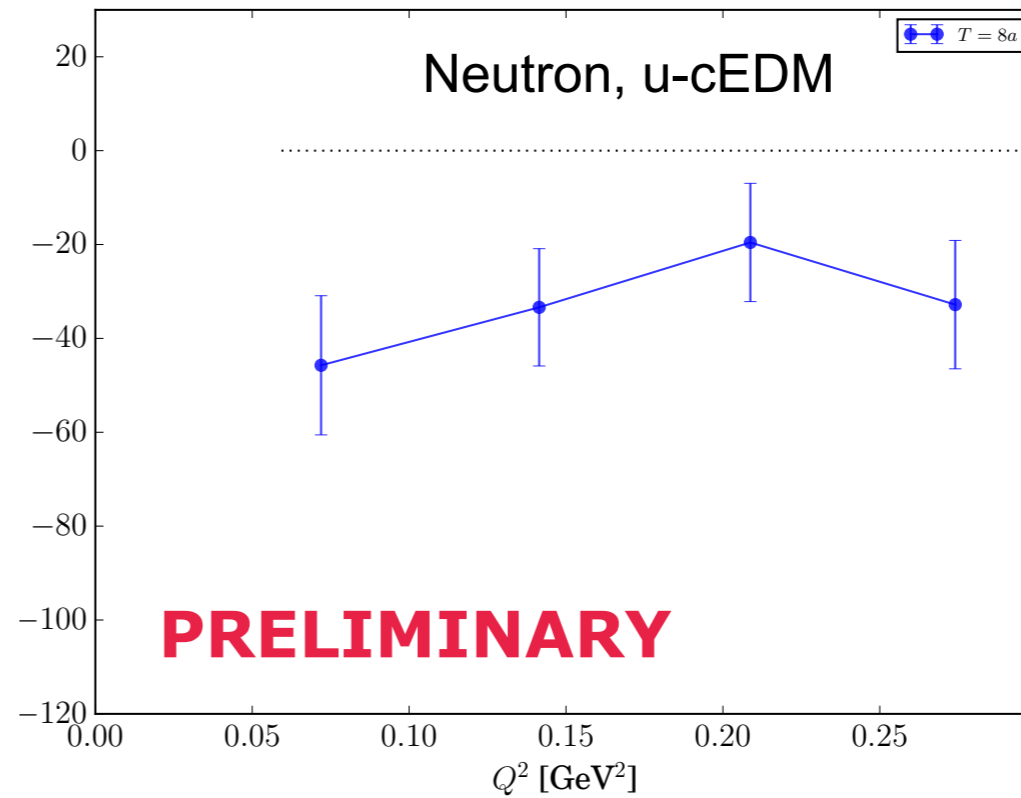
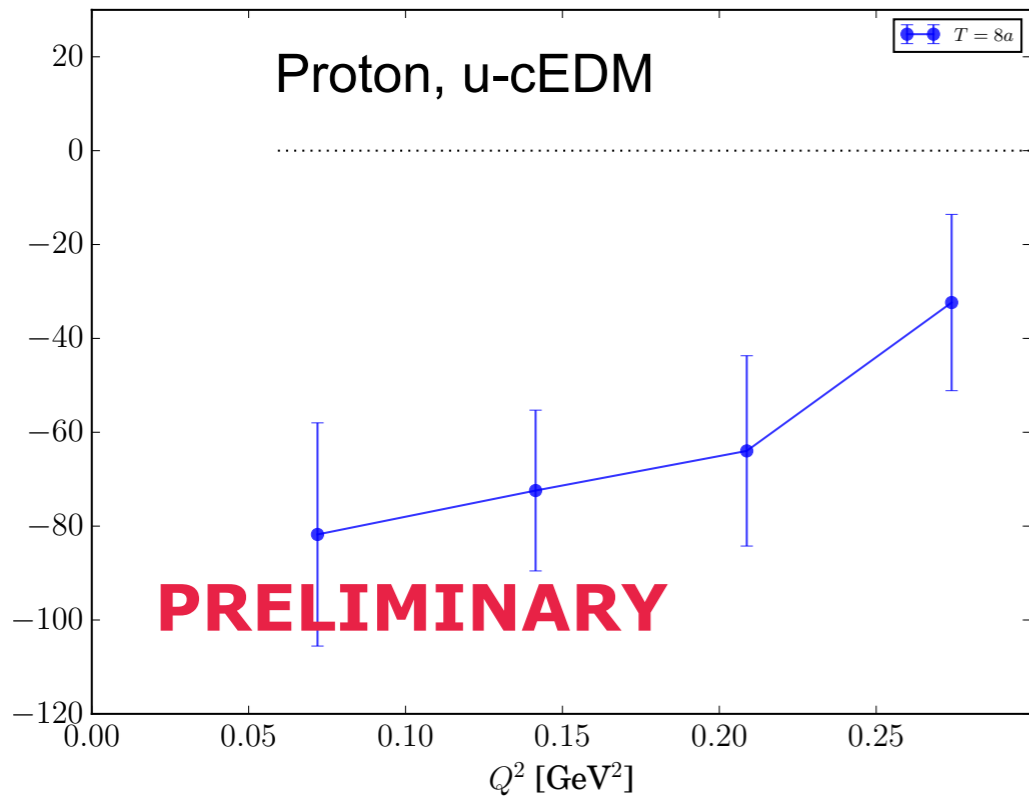
Parity Mixing (Proton)



$$\langle N(t) \bar{N}(0) \rangle_{\mathcal{CP}} = e^{i\alpha_5 \gamma_5} \frac{-i\not{p} + m_N}{2m_N} e^{2i\alpha_5 \gamma_5} e^{-E_N t}$$

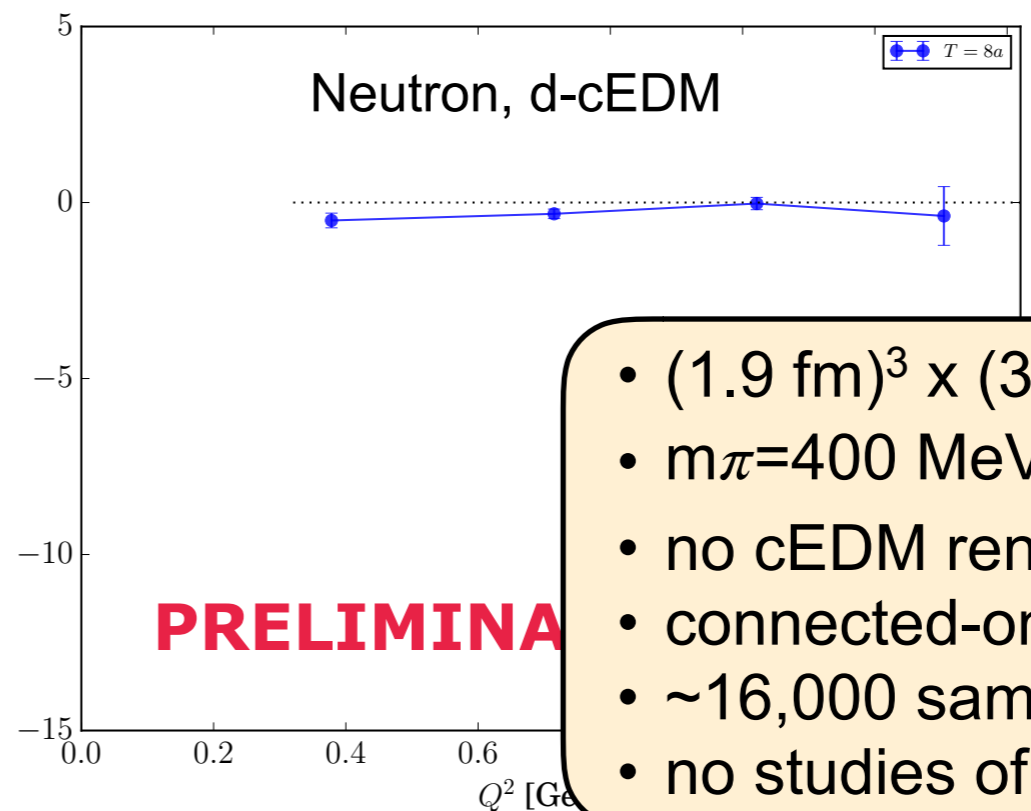
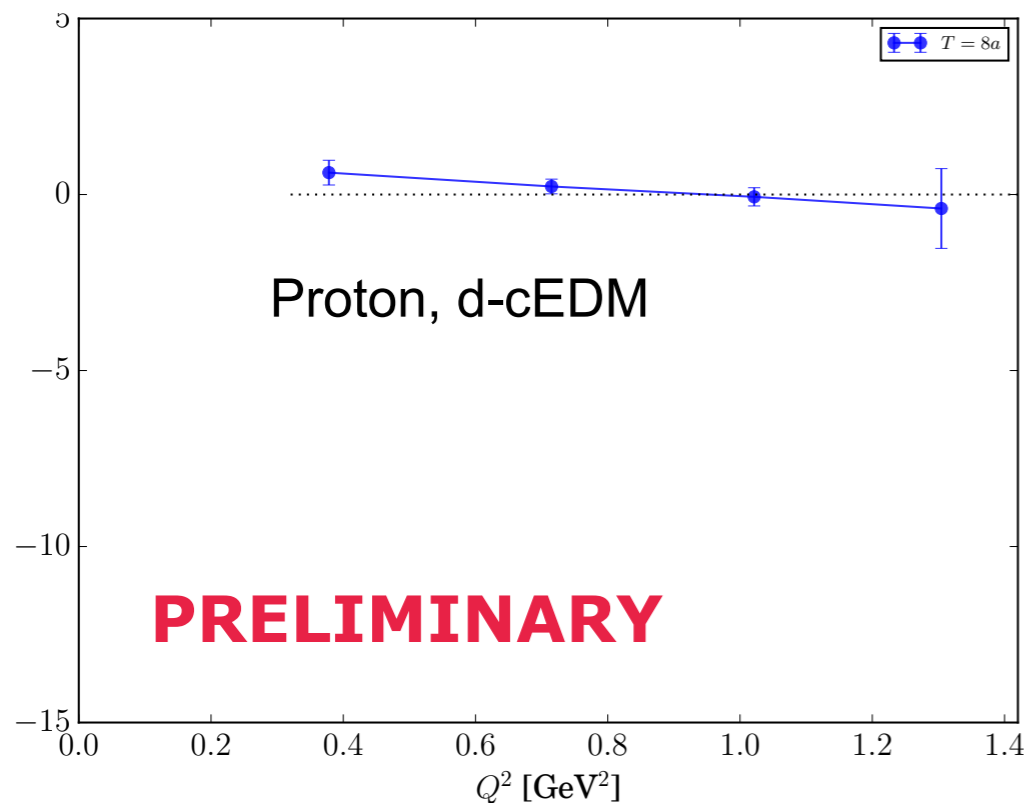
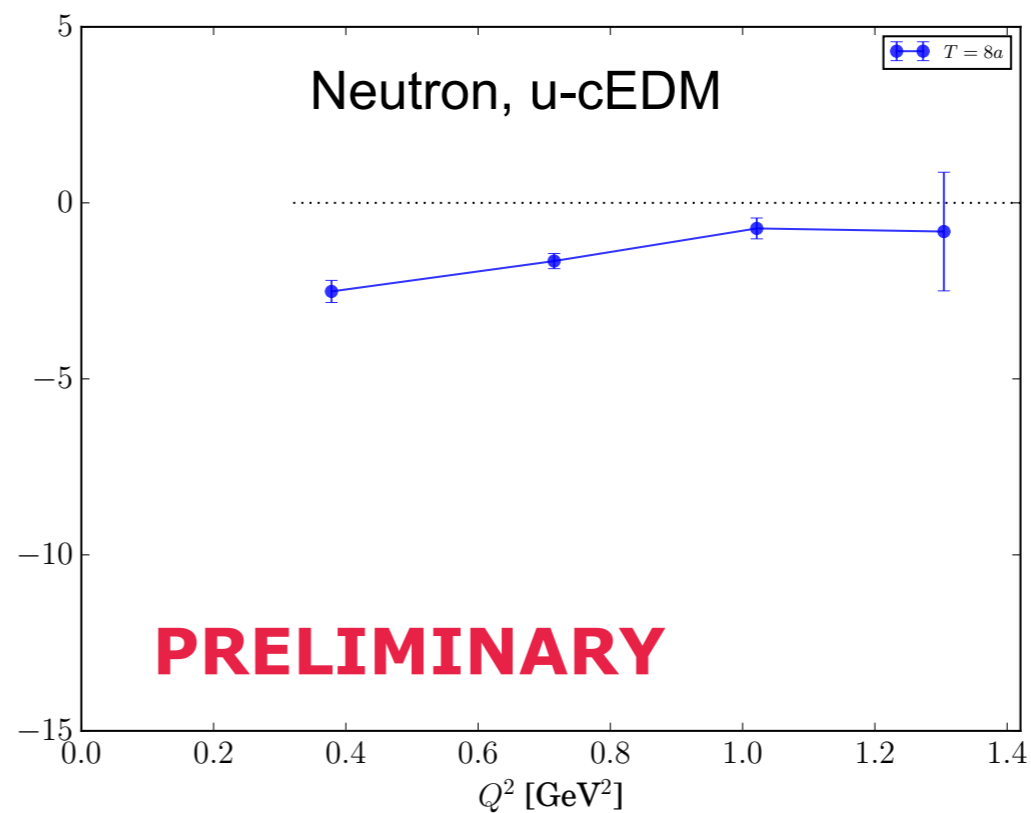
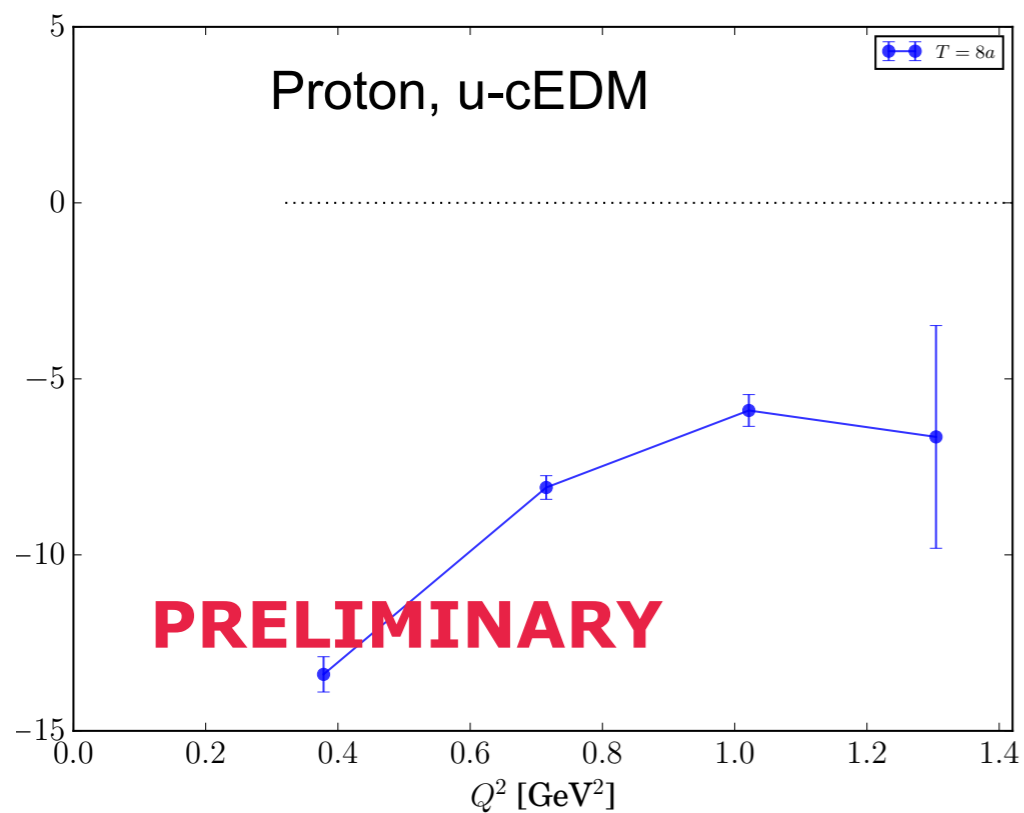
$$\hat{\alpha}_5 = \frac{\alpha_5}{\tilde{d}} = -\frac{\text{ReTr} [T^+ \gamma_5 \cdot C_{2pt}^{CP}(t)]}{\text{ReTr} [T^+ \cdot C_{2pt}^{CP}(t)]}, \quad t \rightarrow \infty$$

Proton & Neutron EDFF Form Factors



- $(4.6 \text{ fm})^3 \times (9.2 \text{ fm})$ box
- $m_\pi = 170 \text{ MeV}$
- no cEDM renormalization
- connected-only
- $\sim 1,000$ samples
- no studies of sys.errors yet

Proton & Neutron EDFF Form Factors



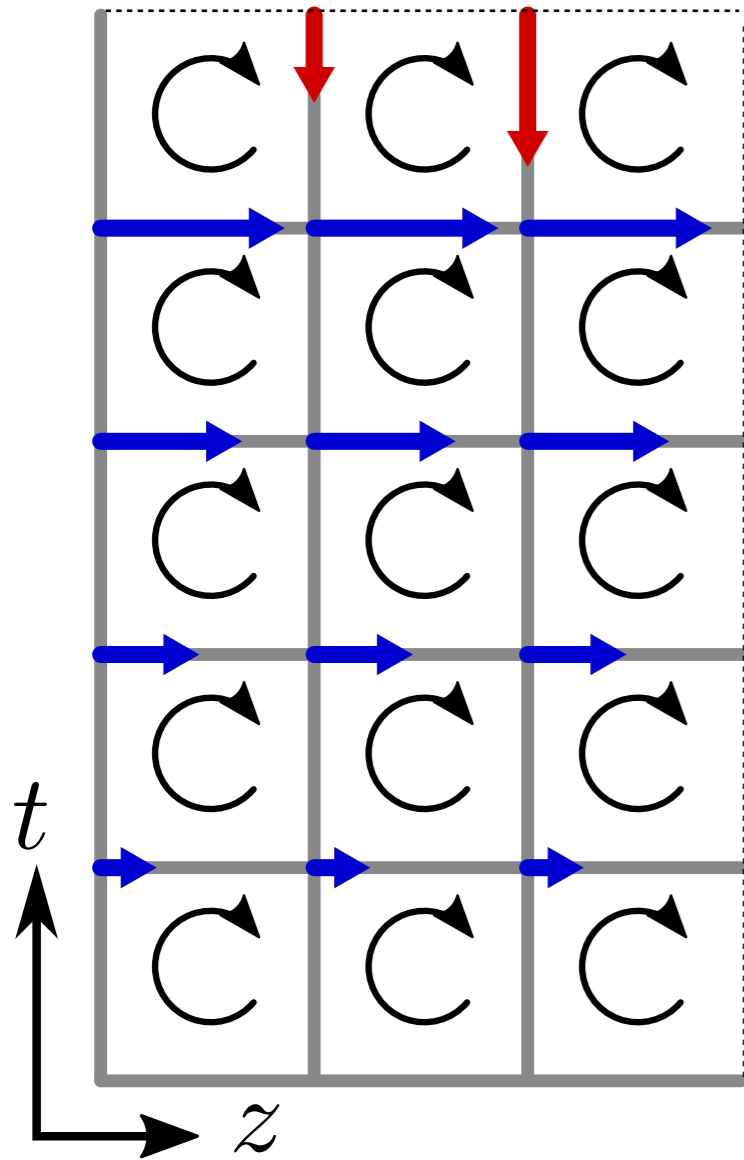
- $(1.9 \text{ fm})^3 \times (3.8 \text{ fm})$
- $m_\pi = 400 \text{ MeV}$
- no cEDM renormalization
- connected-only
- $\sim 16,000$ samples
- no studies of sys.errors yet

Background Electric Field

Accessing magnetic and electric moments at $Q^2=0$

Imag.Minkowski/Real Euc. electric field on a lattice

[W.Detmold et al (2009)] : calculation of hadron polarizabilities



Full flux through the "side" of the periodic box $= q\Phi = 2\pi$

Constant Electric field has to be quantized, $\mathcal{E}_{\min} = \frac{1}{|q_d|} \frac{2\pi}{L_x L_t}$

$$U_\mu \rightarrow e^{iqA_\mu} U_\mu$$

$$A_z(z, t) = n \mathcal{E}_{\min} \cdot t$$

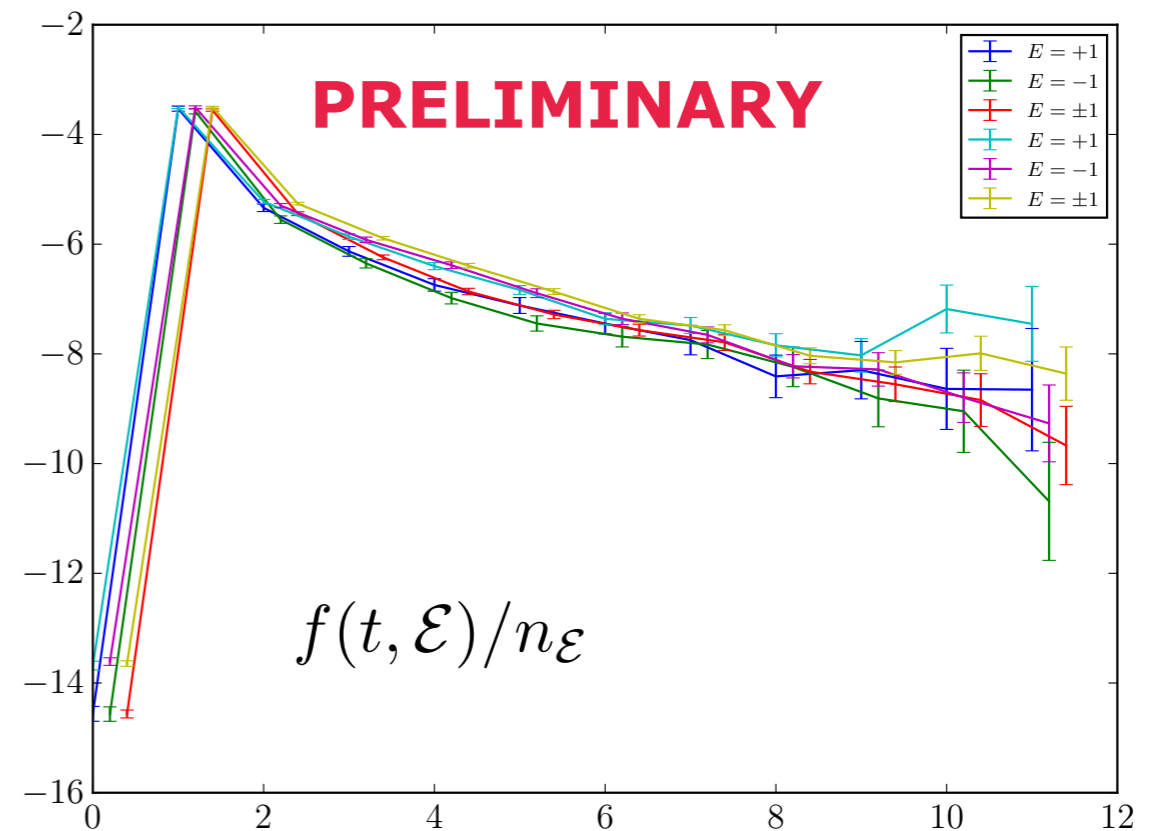
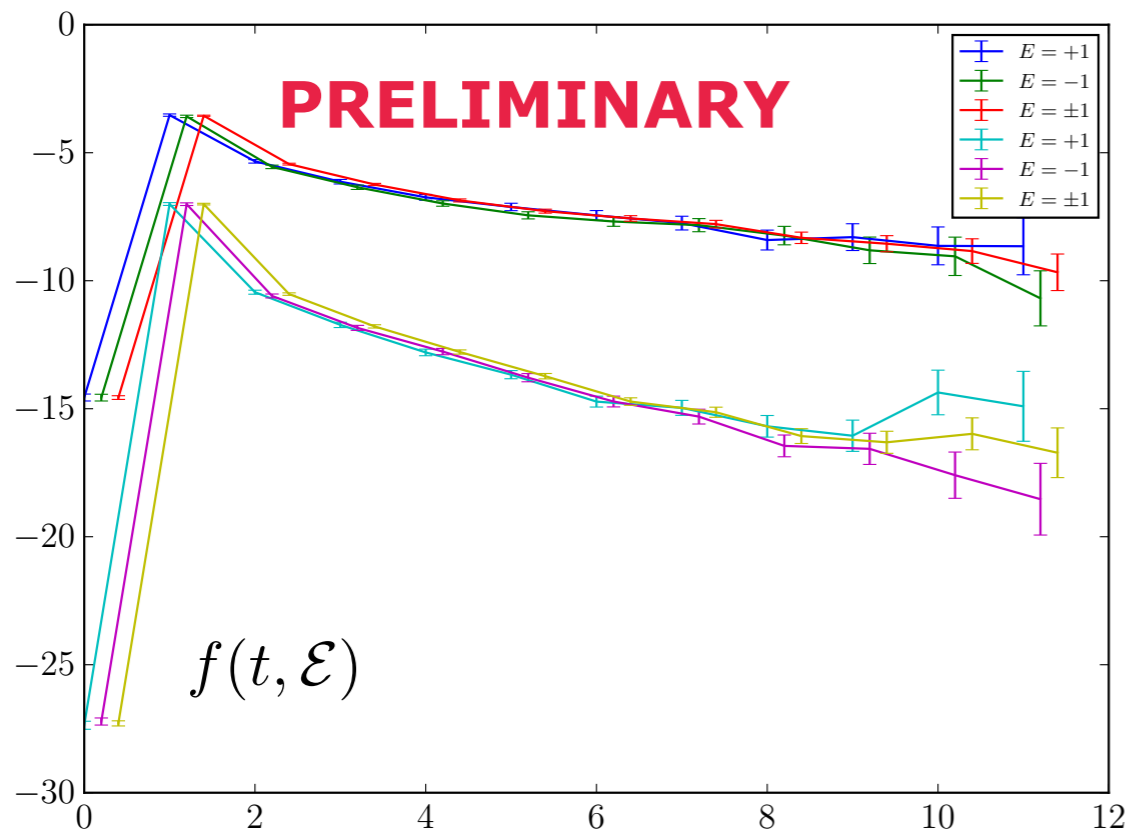
$$A_t(z, t = L_t - 1) = -n \mathcal{E}_{\min} \cdot L_t z$$

CP-odd Neutron Energy Shift

$$\langle N(t)\bar{N}(0) \mathcal{O}_{\overline{CP}} \rangle_{\mathcal{E}} \sim e^{-E_N t} [A - d_N \mathcal{E}_z \Sigma_z t]$$

$$f(t, \mathcal{E}) = \frac{\text{ReTr}[\Sigma_z \cdot \langle N(t)\bar{N}(0) \mathcal{O}_{\overline{CP}} \rangle_{\mathcal{E}}]}{\text{ReTr}[\langle N(t)\bar{N}(0) \rangle_{\mathcal{E}}]} \sim A + d_N \mathcal{E} t$$

- $(1.9 \text{ fm})^3 \times (3.8 \text{ fm})$
- $m_\pi = 400 \text{ MeV}$
- $\mathcal{E}_{\text{min}} = 0.0966 \text{ GeV}^2 = 490 \text{ MeV/fm}$



- Linearity in \tilde{d}_q/d_N , t , and \mathcal{E}
- No renormalization yet

Electric field
on $16^3 \times 32$ lattice

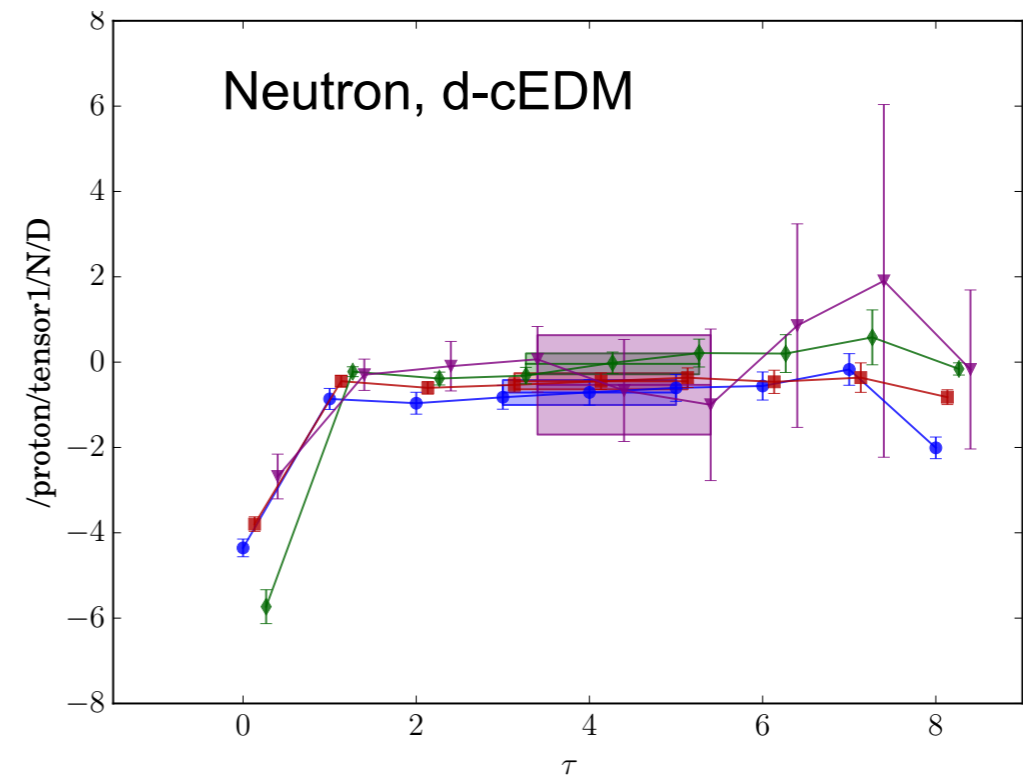
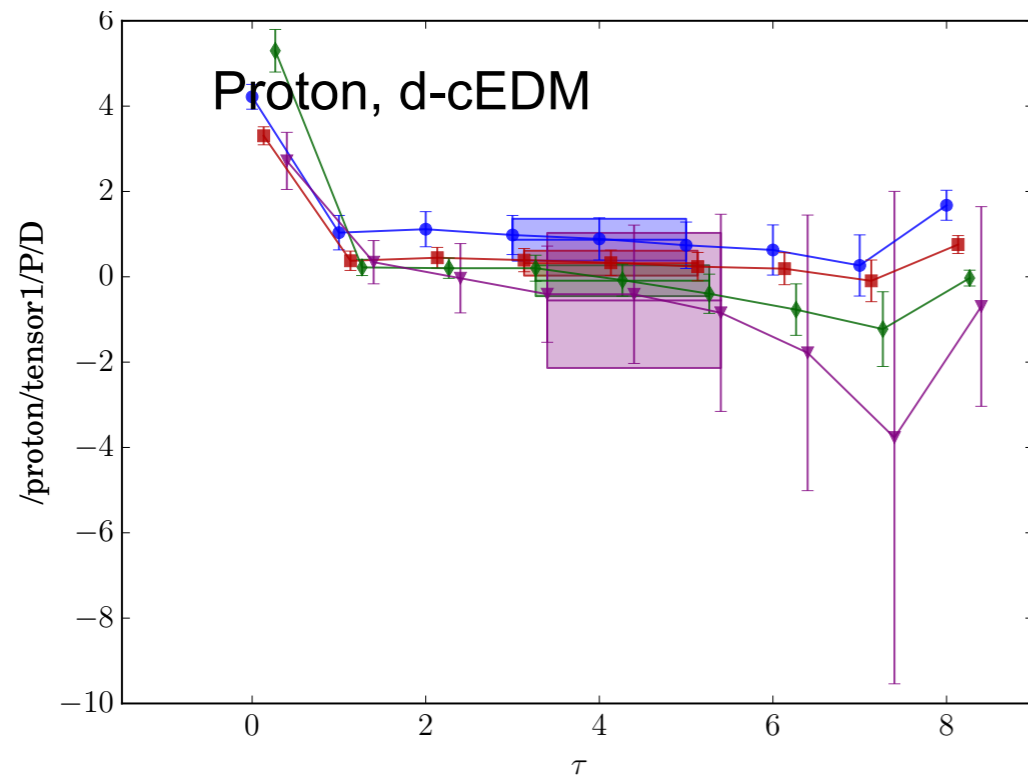
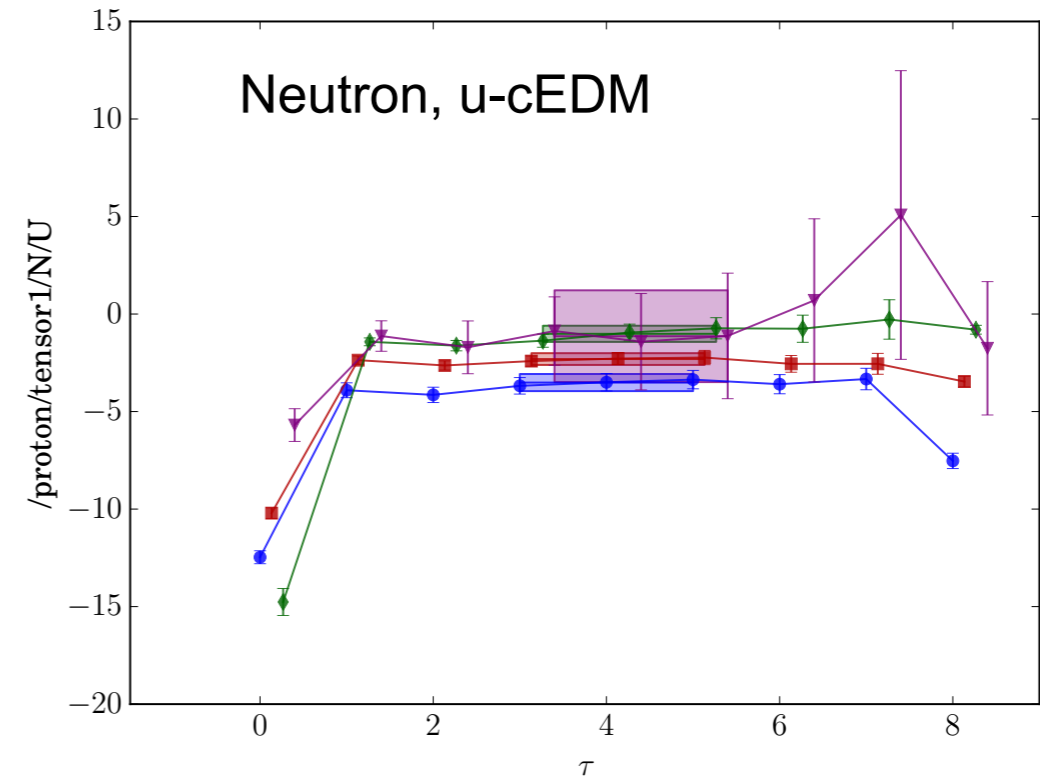
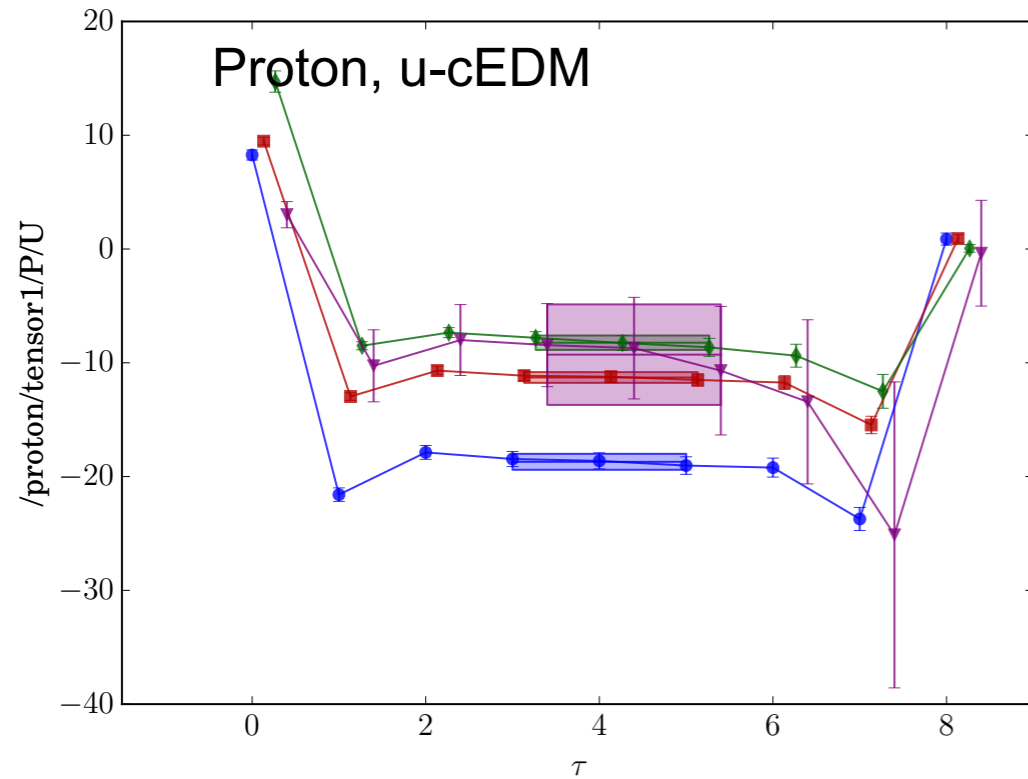
$$\mathcal{E} = \frac{6\pi}{L_x L_t} \approx 0.1 \text{ GeV}^2$$

$$\approx 500 \text{ M(e)V/fm}$$

Summary

- Calculations of θ -induced NEDM are very noisy close to the physical pion mass
- Additional techniques may be necessary
local sampling of topology?
- Initial results for *quark-connected* cEDM-induced EDFF look promising
- Preliminary study with background field methods shows expected qualitative behavior
- Challenges in computing NEDM from cEDM
subtraction of lower-dimension operators
disconnected diagrams
mixing with θ -term in the isoscalar channel

F3 Plateaus (16c32, 400 MeV)



F3 Plateaus (32c64, 170 MeV)

