

Symmetry protected topological phase for interacting bosons with Rydberg atoms

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Experimental collaboration:

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Research group:

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Kevin Kleinbeck, **Sebastian Weber**



Strongly interacting
Rydberg slow light
polaritons



INTEGRATED QUANTUM
SCIENCE AND TECHNOLOGY



From few to many-body
physics with dipolar quantum
gases

How to characterize states of matter?

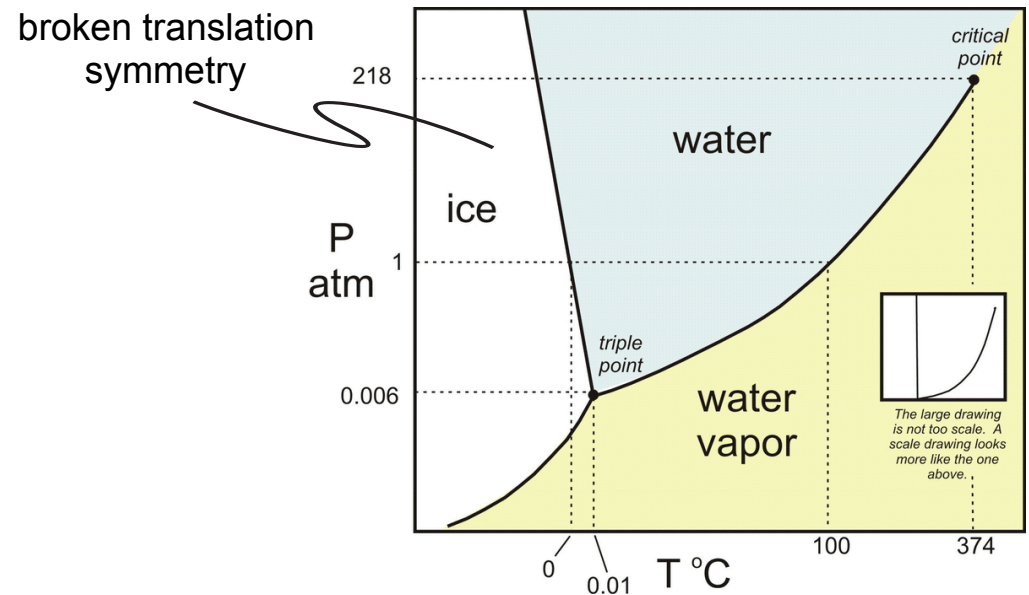
Characterize phases of matter

- based on Landau paradigm
- symmetry breaking
- order parameter and long range order
- thermal and quantum phase transitions

Extremely successful

- band insulator/Fermi liquids
- crystals
- superfluids
 - Bose-Einstein condensate
 - superfluid Helium
- superconductors
- ferromagnets and anti-ferromagnets

phase diagram of water



How to characterize states of matter?

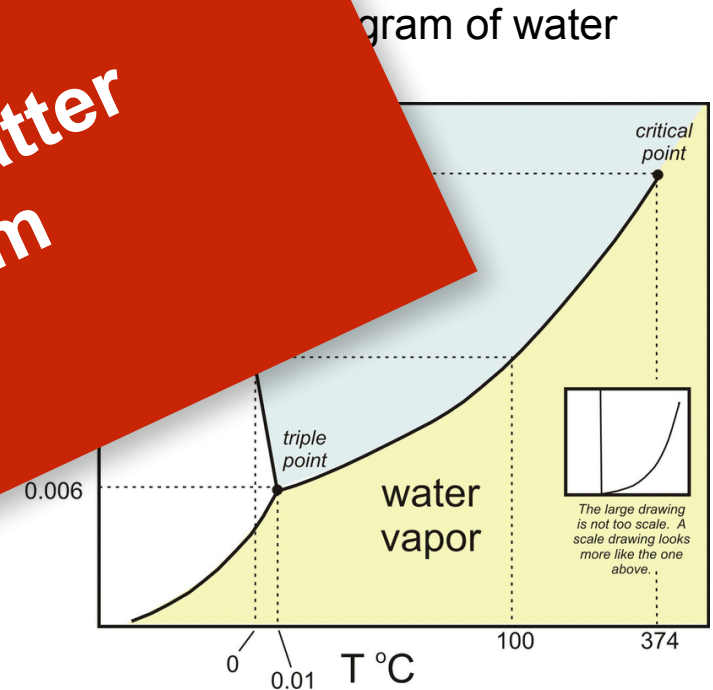
Characterize phases of matter

- based on Landau paradigm
- symmetry breaking
- order parameter and long range order
- thermal and quantum phase transitions

Extremely subtle

- band structure
- crystals
- superfluids
 - Bose-Einstein condensates
 - superfluids
- superconductors
- ferromagnets and anti-ferromagnets

But, not all states of matter follow this paradigm



Topological phases

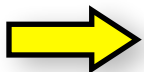
Characterizing ground states of quantum many-body systems at $T=0$

- absence of symmetry breaking
- gapped phases

Definition:

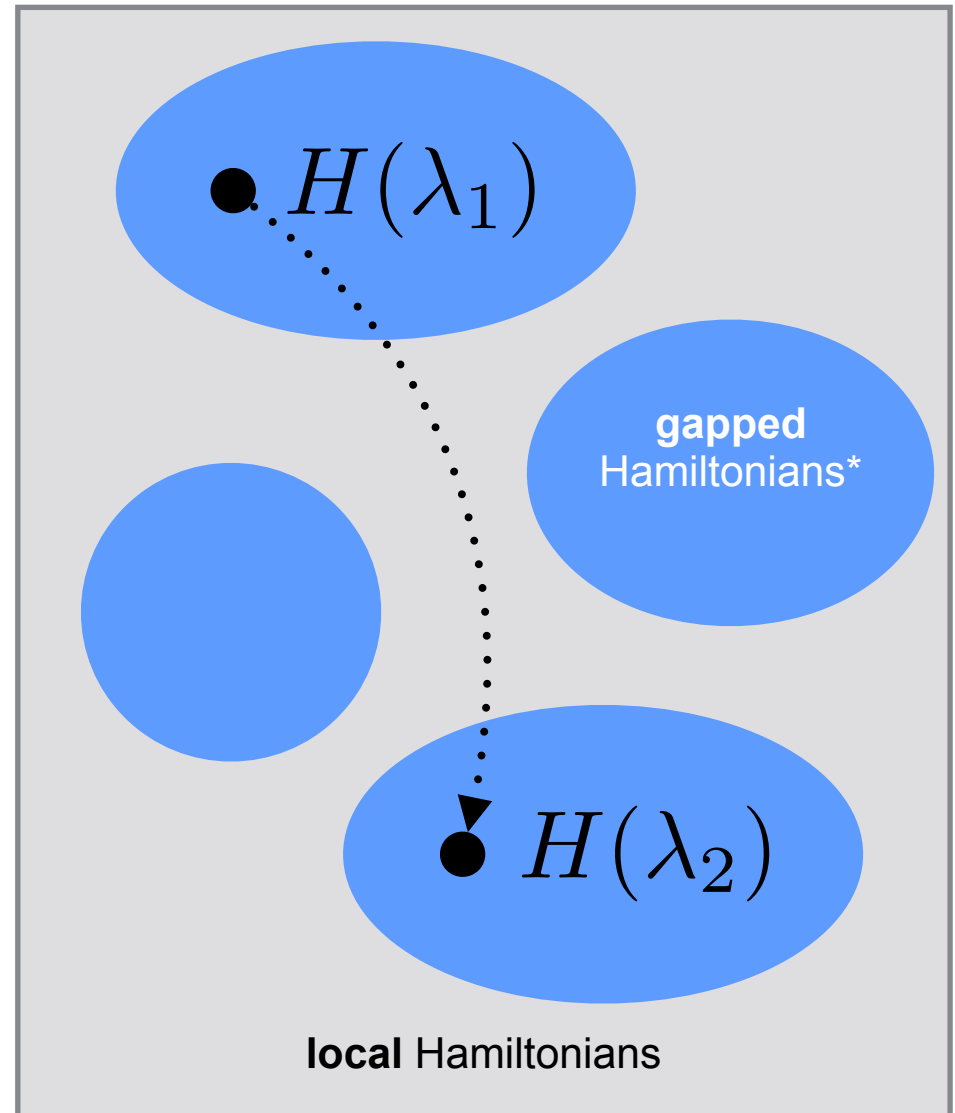
two states are in the same topological phase if they can be smoothly transformed into each other without closing the gap

Example in 2D:
integer quantum Hall states, fractional quantum Hall states, toric code,



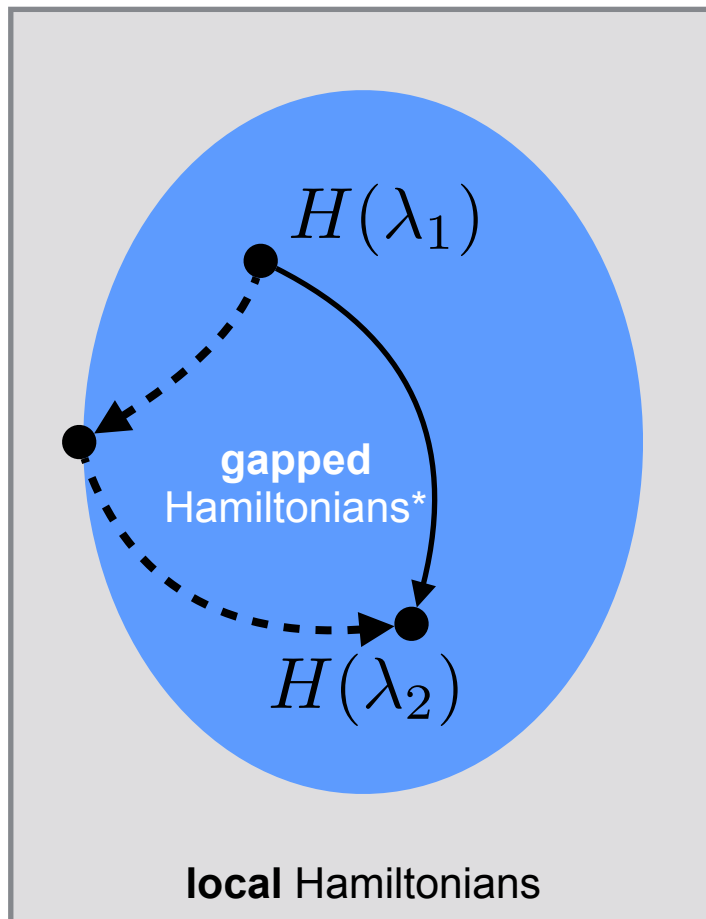
Are there other topological phases also in 1D?

Two dimensions




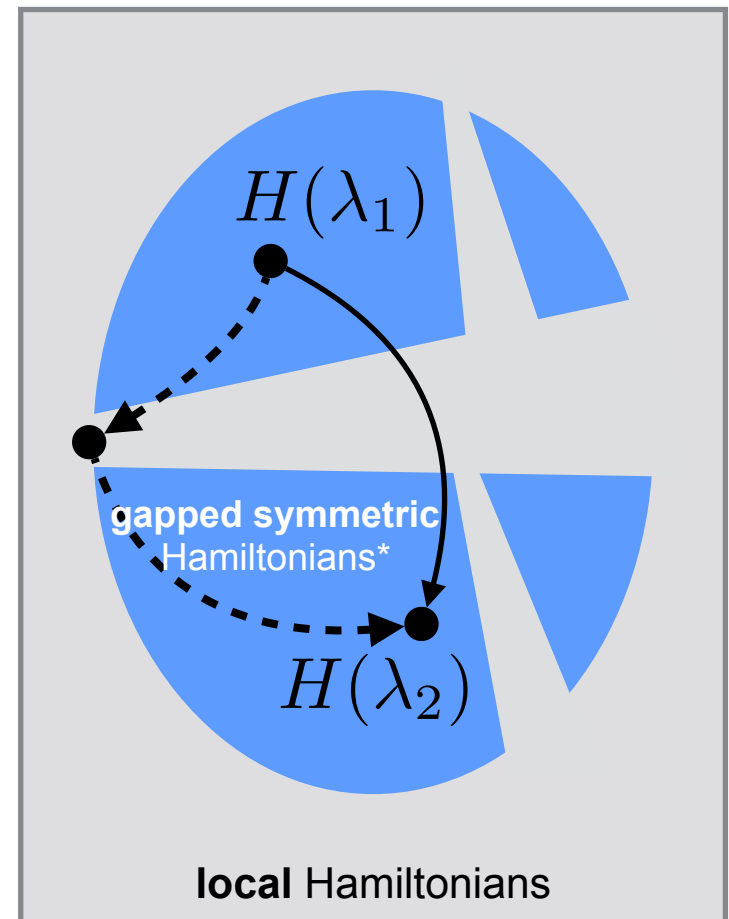
Symmetry protected topological phases

- restrict to systems, which satisfy a certain symmetry
- symmetry group \mathcal{S} with $[H, \mathcal{S}] = 0$



One dimension

require a
symmetry

 $[H, \mathcal{S}] = 0$



Topological insulators: fermionic SSH model

- non-interacting Fermions

- Hamiltonian

$$H = -J \sum_{i \in \text{even}} c_{i+1}^\dagger c_i - J' \sum_{i \in \text{odd}} c_{i+1}^\dagger c_i + h.c.$$

$$= \sum_{ij} c_i^\dagger \hat{H}_{ij} c$$

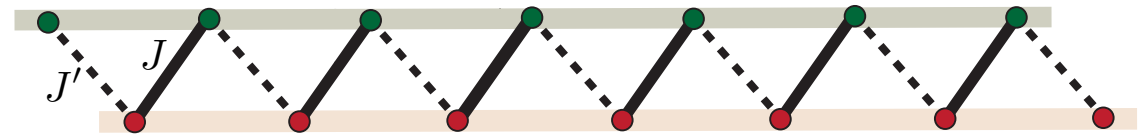
- gapped ground state at half-filling

- 10-fold way classifications of topological insulators

(Chiu, et al., RMP 2016; Ludwig, 2016)



AIII,
Z-index



- sub-lattice (chiral) symmetry
(anti-unitary operator)

$$\mathcal{S}_S = \prod_i \left[c_i + (-1)^i c_i^\dagger \right] K$$

complex conjugation

- on single particle Hamiltonian

$$U_S \hat{H} U_S^\dagger = -\hat{H} \quad (U_S)_{ij} = (-1)^j \delta_{ij}$$

Trivial phase $J' > J$

- conventional band insulator

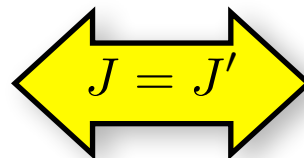
- excitation gap

- filling $n = 1/2$

- decoupled limit $J = 0$

$$|\psi\rangle = \prod_{i \in \text{odd}} (c_i^\dagger + c_{i+1}^\dagger) |0\rangle$$

topological phase
transition



gapless state

Topological phase $J > J'$

- bulk excitation gap

- half filling $n = 1/2$

- four-fold degenerate ground state for open chain

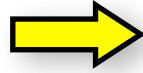


zero energy modes
localized at the edges

Topological insulators: fermionic SSH model

SPT phase

- properties of the SSH model are robust under any perturbation commuting with the symmetry



- four-fold degenerate ground state for open chain
- zero energy modes localized at the edge

Perturbations respecting the symmetry

- random variations of the hopping
- longer range hoppings between the chains

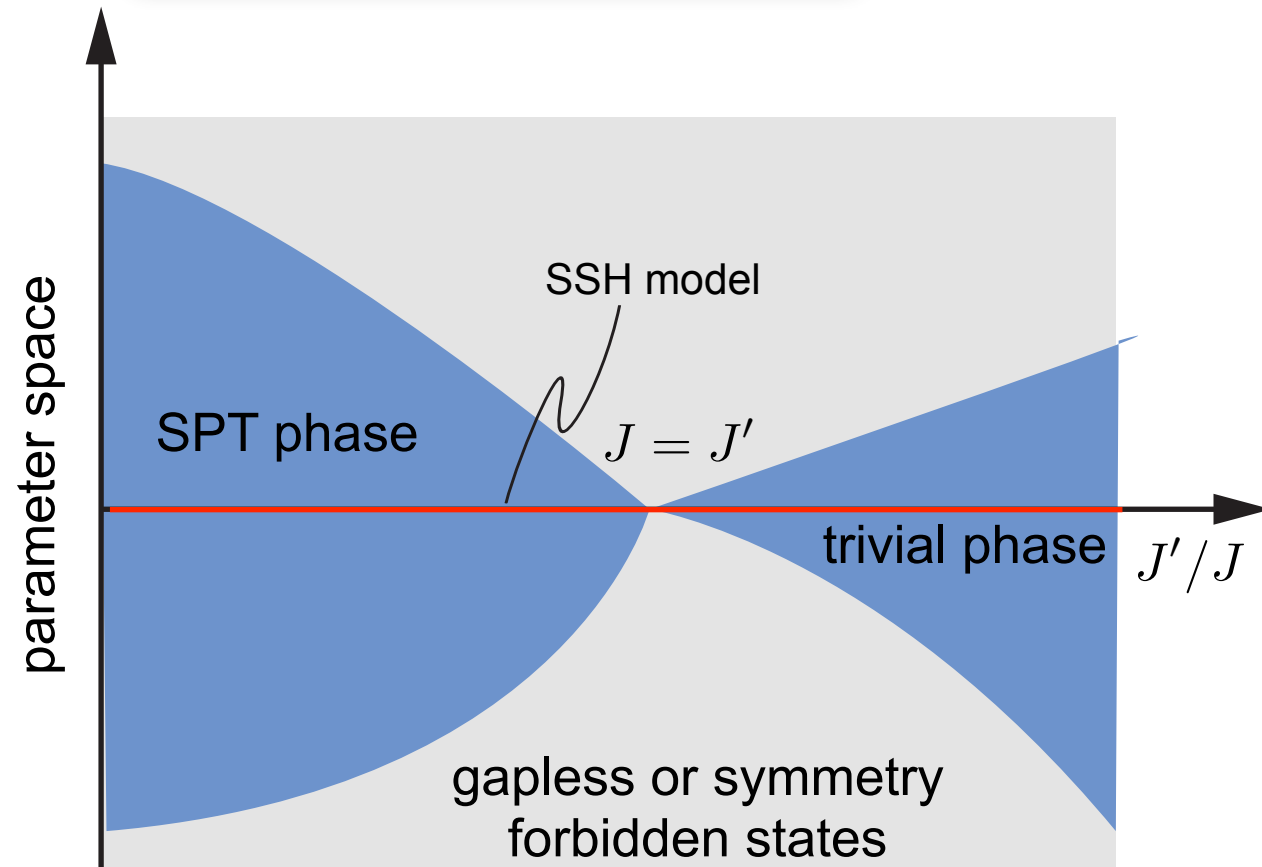
$$c_{i+3}^\dagger c_i + c_i^\dagger c_{i+3}$$

Perturbations breaking the symmetry

- next-nearest neighbor hopping

$$c_{i+2}^\dagger c_i + c_i^\dagger c_{i+2}$$

- on-site shifts $c_i^\dagger c_i$



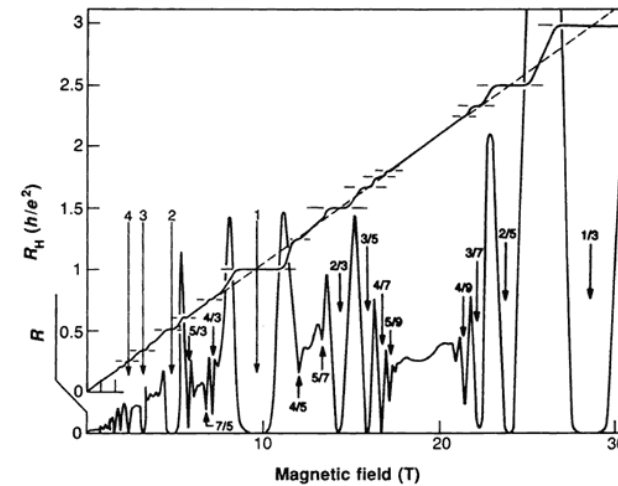
Topological phases: condensed matter

Integer quantum Hall effect (1980):

- 2D electron gas in a strong magnetic field
- present for non-interacting fermions

Fractional quantum Hall effect (1982):

- 2D electron gas in a strong magnetic field
- strong interactions between the particles

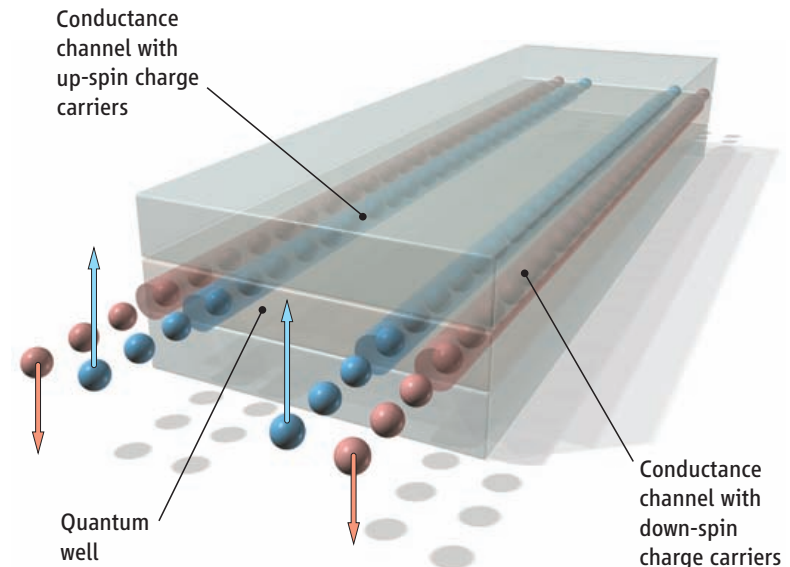


Spin-1 anti-ferromagnets in 1D (1990):

- realization of a bosonic SPT phase
- protecting symmetry $SO(3)$
- NENP materials $[\text{Ni}(\text{C}_2\text{H}_8\text{N}_2)_2(\text{NO}_2)]\text{ClN}_4$

Topological insulators (2007):

- spin-Hall effect in HgTe quantum wells
- time-reversal symmetry



Topological phases: artificial matter

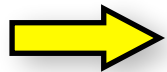
Artificial matter

- cold atomic gases in optical lattices
- Ion traps
- Photonic circuits
-

Topological phases: artificial matter

Topological band structures and edge modes

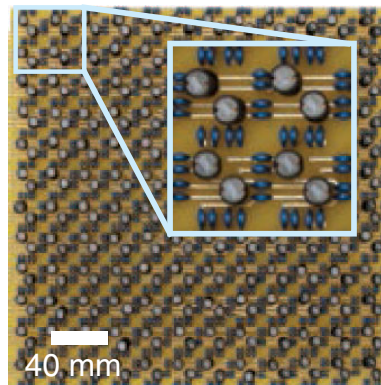
- motivated by topological phases for non-interacting fermions
- probing spectrum of a single particle Hamiltonian
- property of the coupling matrix



accessible in classical and quantum systems

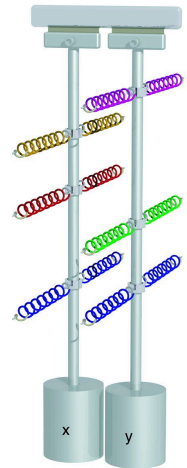
Edge modes in photonic systems:

- optical regime
Hafezi, et al, Nat Phot. 2013
Rechtsman et al, Nature 2013
- radio-frequencies
Ningyuan et al. PRX 2015



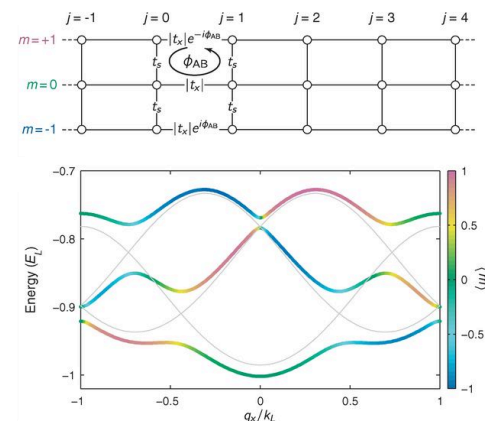
Classical coupled pendulums:

- time reversal invariant
Süsstrunk Huber, Science 2015
- driven rotating pendulums
Nash et al., PNAS 2015



Cold atomic gases

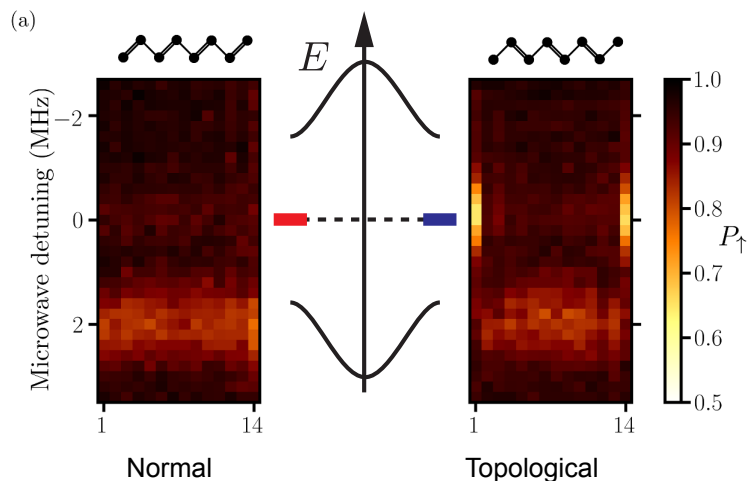
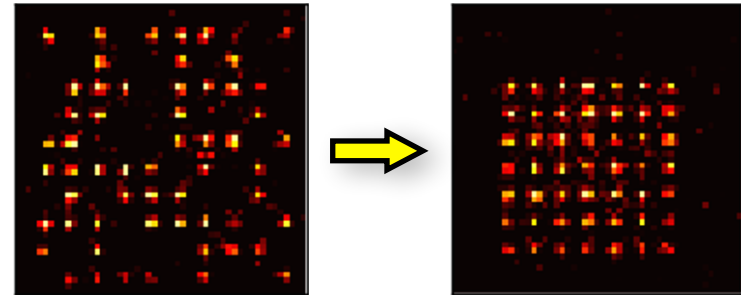
- artificial gauge fields
Stuhl et al, Science, 2015
- artificial dimensions
Mancini et al, Science, 2015
- optical lattices and lattice shaking
Jotzu et al. Nature 2014, Aidelsburger et al. Nat. Phys. 2014
Lohse et al. Nat. Phys. 2015, Flaschner et al. Science 2016



Outline

Experimental setup with Rydberg atoms

- single atoms in optical tweezers
- assembly of arbitrary structures



Edge modes in the SSH model

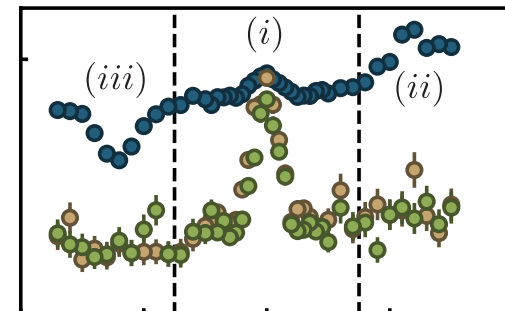
(S. de Léséleuc, *et al*, arXiv:1810.13286)

- single particle physics (independent on statistics)
- observation of localized edge modes

Symmetry protected topological phase

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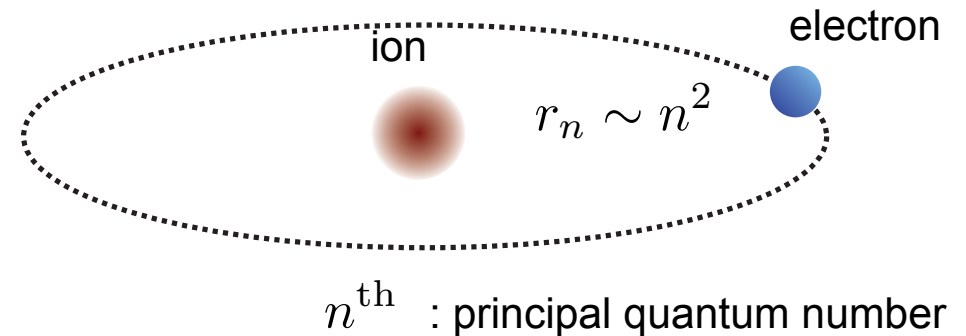
- ground state of the interacting many-body system at half filling
- Spectroscopic detection of zero energy edge states



Experimental setup with Rydberg atoms

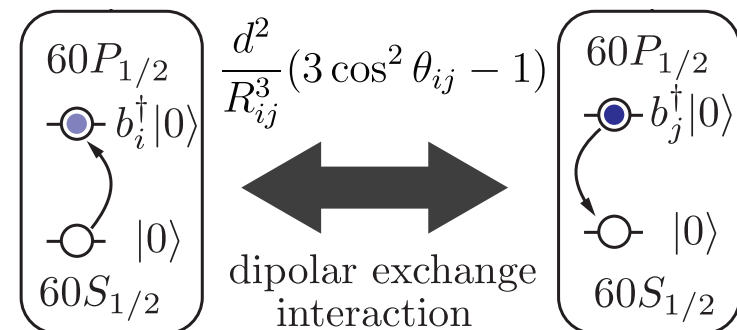
Rydberg atoms

- one electron excited into a state with high principal quantum number n
- here, Rubidium atoms $n \sim 40 - 100$, excited into s-states and p-states



Rydberg-Rydberg interaction

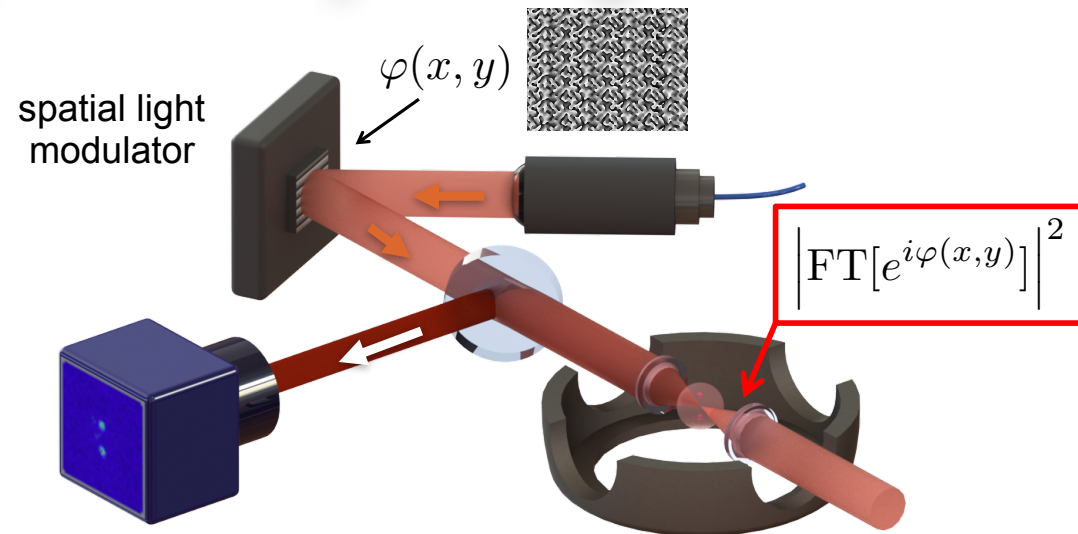
- strong van der Waals interactions between Rydberg states
 - attractive or repulsive
 - $C_6 \sim n^{11}$
- dipolar exchange interactions
 - exchange of excitation between two different Rydberg states
 - $d \sim n^2$



Experimental setup with Rydberg atoms

Single atoms trapped in optical tweezers

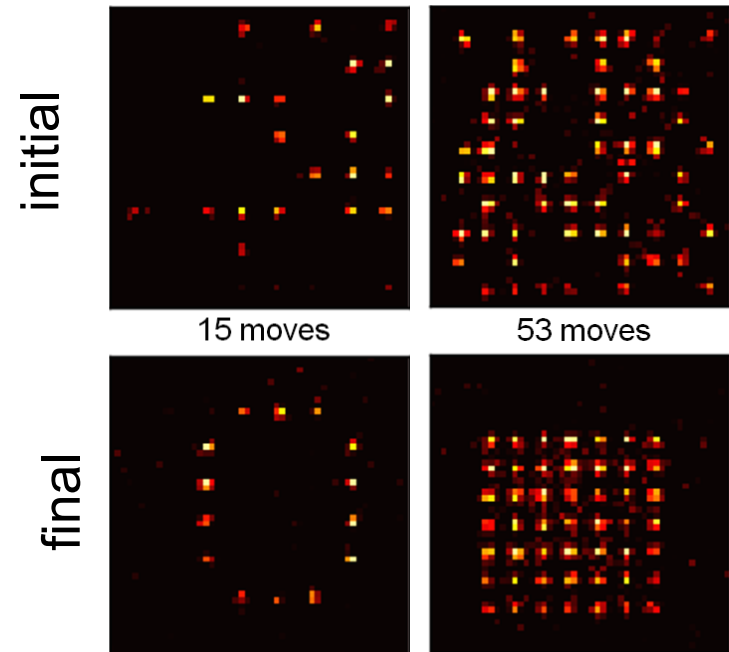
- individual traps for a single atom
- not in ground state of the trapping potential
- single site resolution



Deterministic assembly in arbitrary structures and lattices

- loading from a cold thermal cloud
 - stochastic loading
- prepare lattice structure by moving the filled traps
- prepare arbitrary 2D as well as 3D structures
- achieved by different groups:
 - Paris (2D) Science **354**, 1021 (2016)
 - Harvard (1D), Science **354**, 1024 (2016)
 - Korea (2D), Nat. Comm. **7**, 13317 (2016)

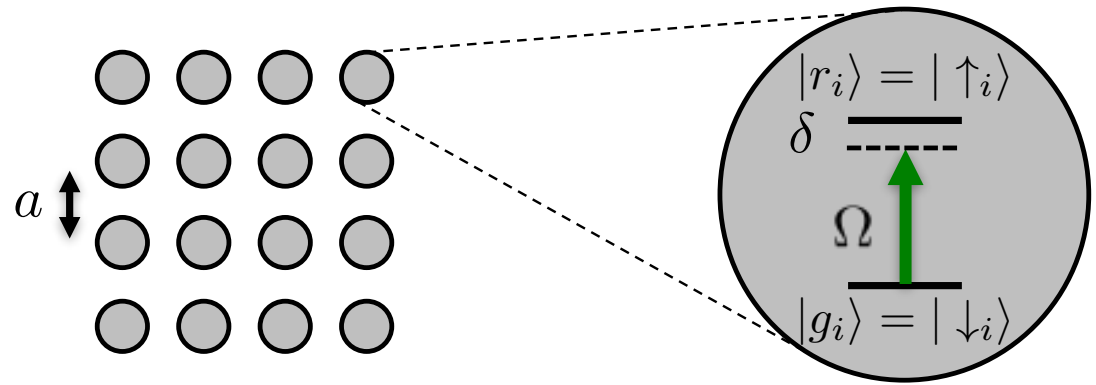
Barredo, *et al.*, Science **354**, 1021 (2016)



Experimental setup with Rydberg atoms

Quantum Ising like models

- all atoms coupled to a Rydberg S-state
- van der Waals interaction between Rydberg states



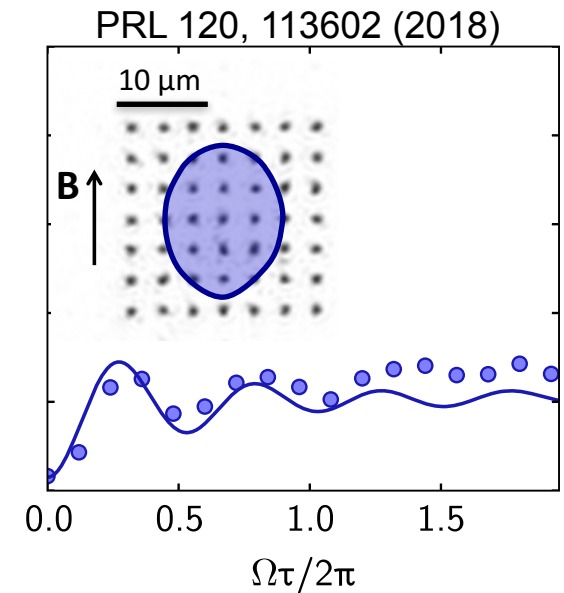
$$H = \Omega \sum_i \sigma_i^x + \sum_i \Delta_i \sigma_i^z + \sum_{i \neq j} \frac{C_6}{|\mathbf{r}_i - \mathbf{r}_j|^6} n_i n_j$$

transverse field
longitudinal field
Ising type interaction

$$n_j = \frac{1 + \sigma_j^z}{2}$$

Quantum simulation of spin models

- non-equilibrium quench dynamics
- time dependent driven and disordered systems
- dissipative systems by including spontaneous decay
- Labuhn, *et al.*, Nature 534, 667 (2016)
- Bernien, *et al.*, Nature 551, 579 (2017)
- ...



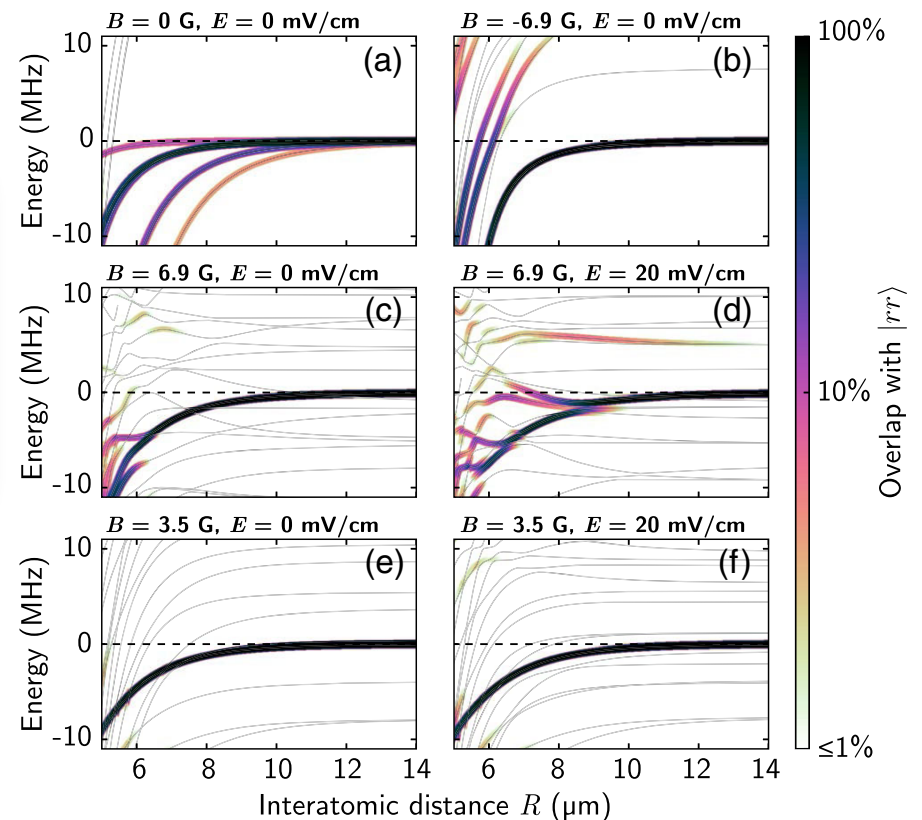
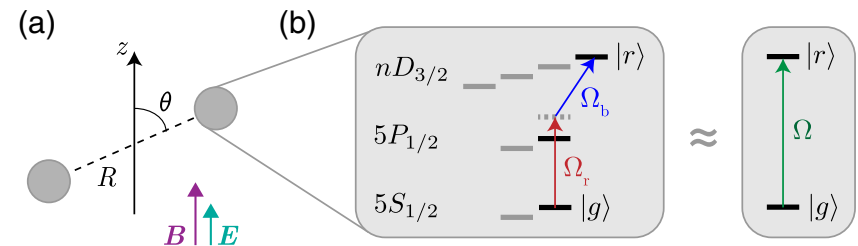
Experimental setup with Rydberg atoms

Importance of microscopic details of interaction potentials

- pair-interactions are in general much more complex than pure van der Waals
- many crossings with additional Rydberg pair states

A Rydberg interaction calculator

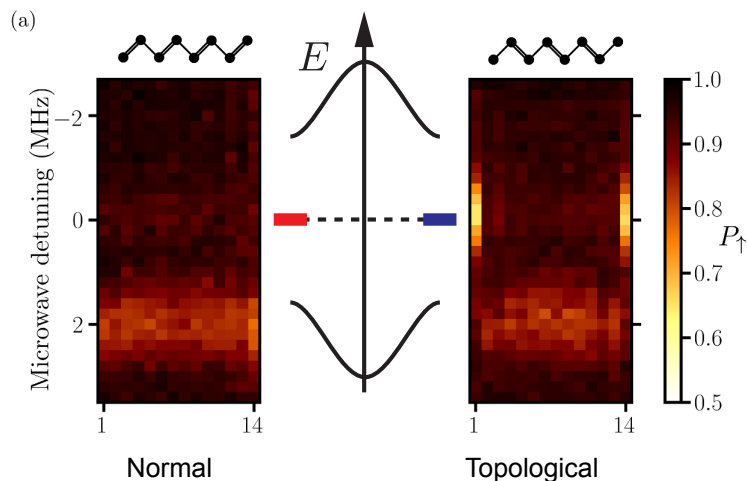
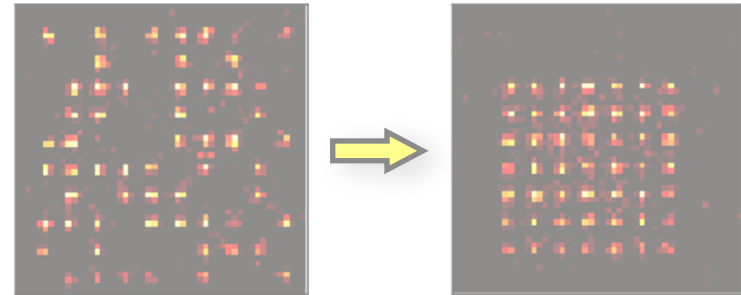
- open source software
- S. Weber, *et al.*, J. Phys. B **50**, 133001 (2017)



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Edge modes in the SSH model

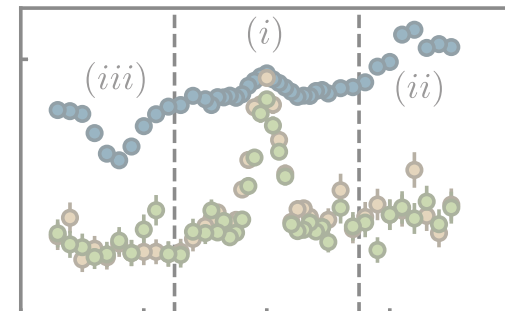
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- single particle physics (independent on statistics)
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Symmetry protected topological phase

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- ground state of the interacting many-body system at half filling
- Spectroscopic detection of zero energy edge states

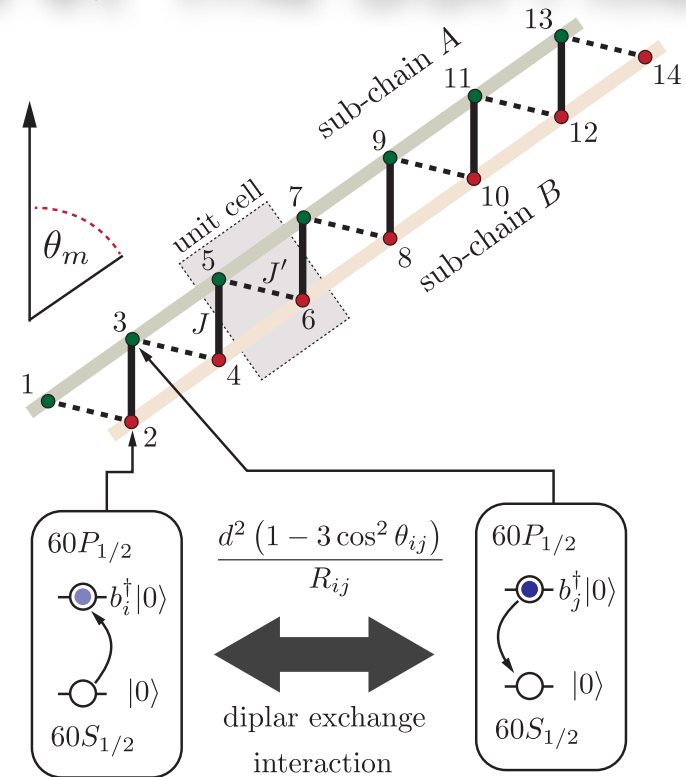


Single particle properties of the SSH model

Hamiltonian of the single particle SSH model for bosons

- all atoms in a Rydberg S-state
- bosonic excitation: Rydberg P-state
- hopping by dipolar exchange on a dimerized chain

$$H = \sum_{ij} b_i^\dagger \hat{H}_{ij} b_j = \sum_{i \in A, j \in B} J_{ij} [b_i^\dagger b_j + b_j^\dagger b_i]$$



Chiral symmetry

- unitary matrix with

$$U_S \hat{H} U_S^\dagger = -\hat{H}$$

- here, we obtain

$$(U_S)_{ij} = (-1)^j \delta_{ij}$$



- chiral symmetry does not allow for hopping within the same sub-chain

- dipolar hopping naturally gives rise to longer range hopping



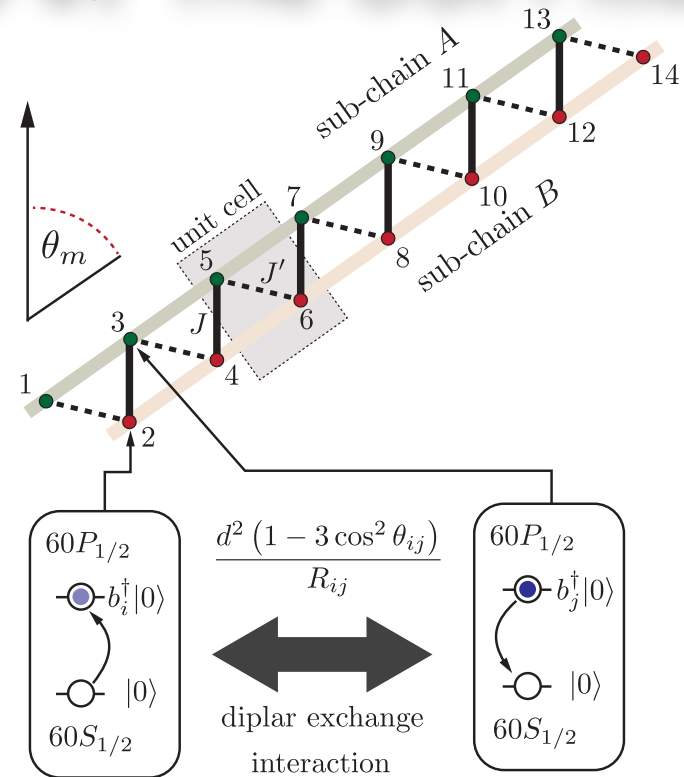
Here, we satisfy the symmetry by employing the anisotropy of the dipole-dipole interaction

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Chiral symmetry

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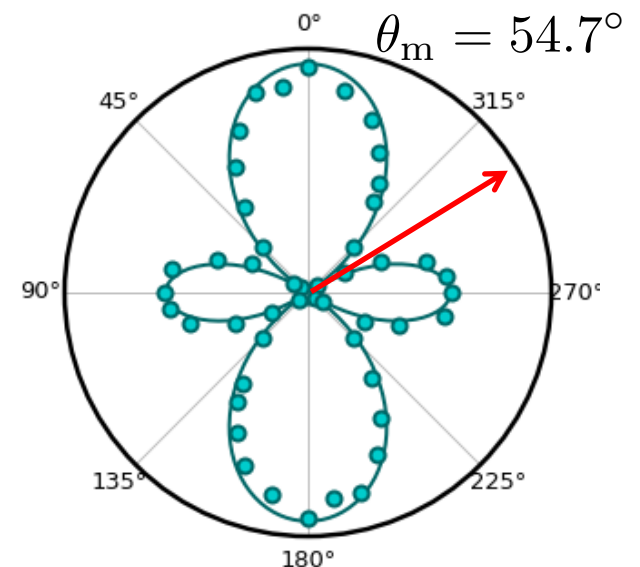
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$$(U_S)_{ij} = (-1)^j \delta_{ij}$$

experimentally measured



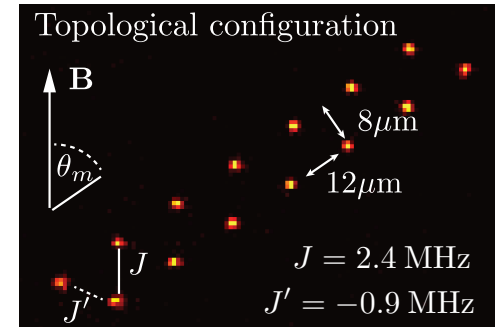
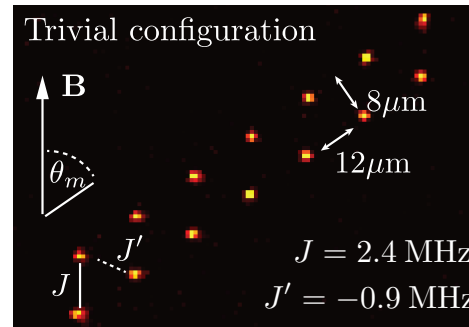
angular dependence



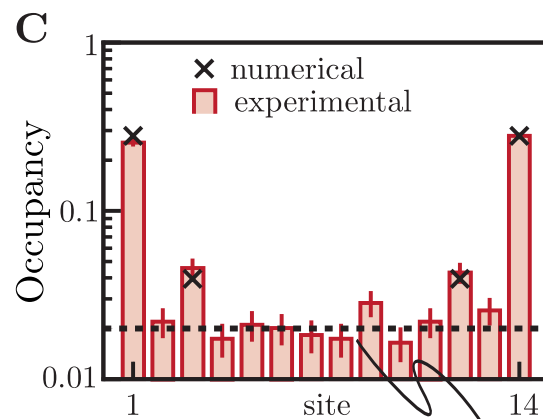
Edge modes of the SSH model

Spectroscopy of the spectrum

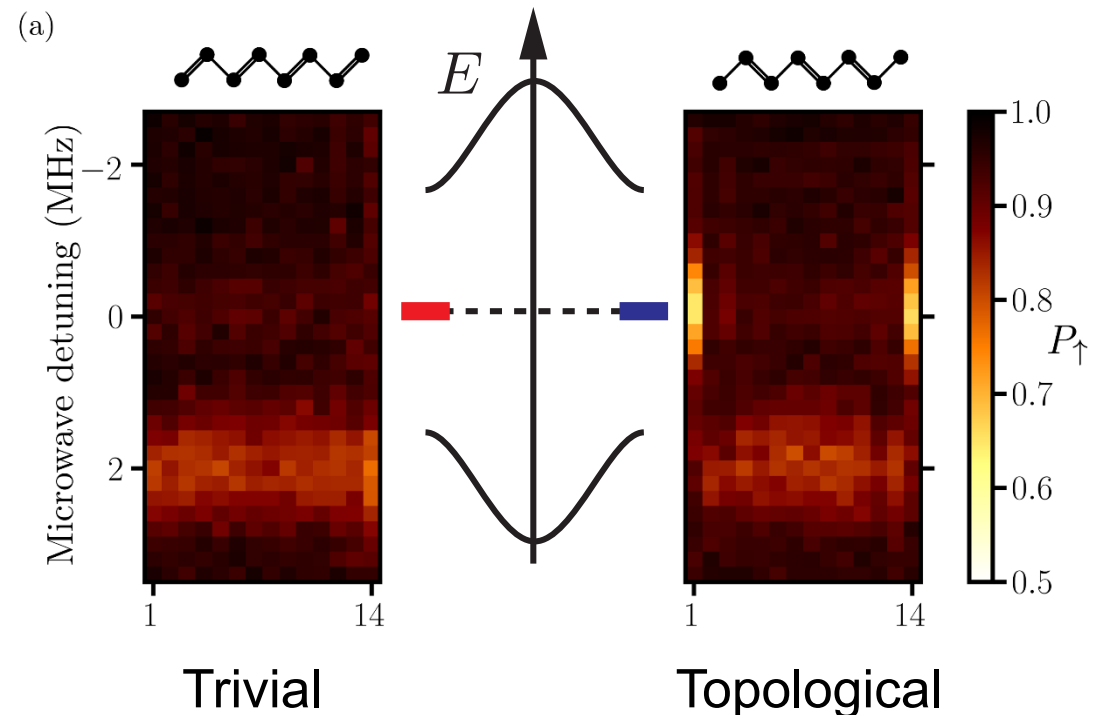
- microwave field coupling the S-state to the Rydberg P-state
- no-momentum transfer: only couple to half of the states
- observation of zero energy states localized at the edge of the system



Edge state localization



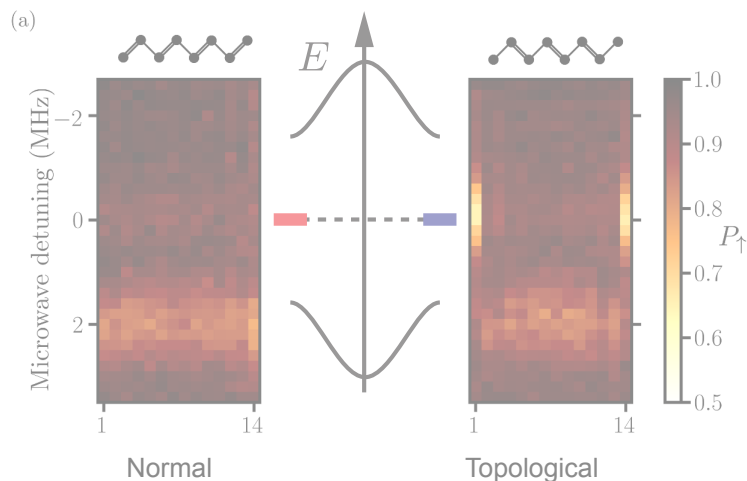
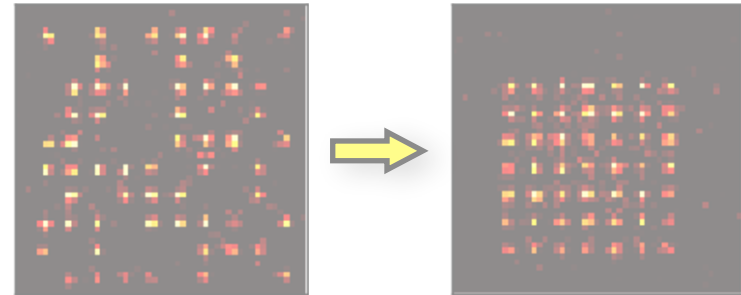
imperfections of the experiments



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Edge modes in the SSH model

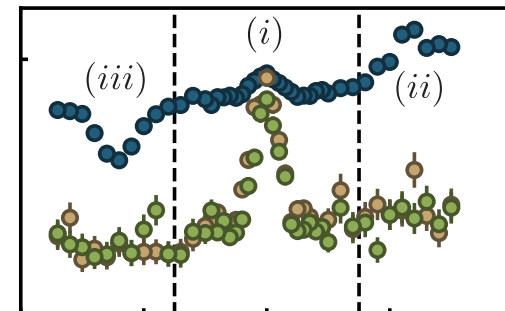
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Symmetry protected topological phase

Topological phases with Bosons

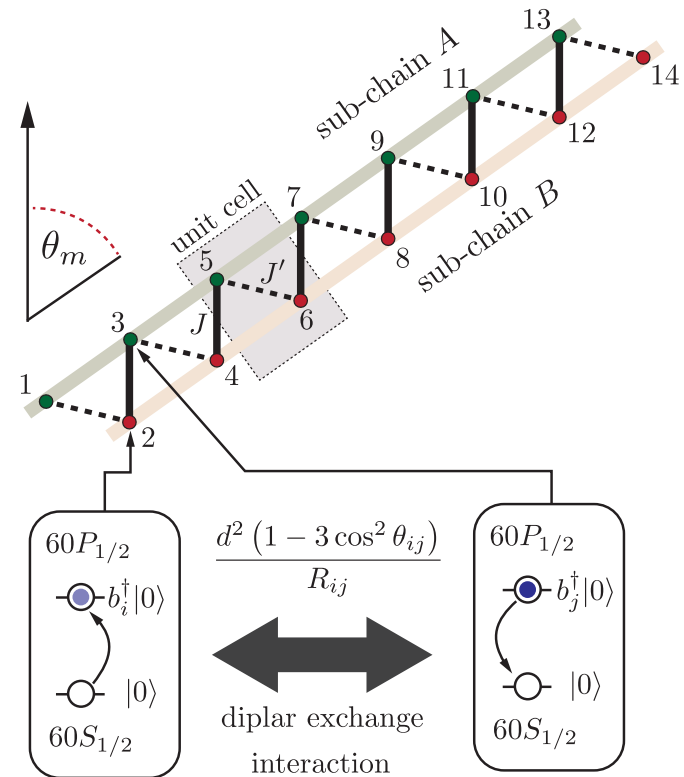
- non-interacting Bosons always from a BEC



topological phases with bosons require strong interactions

Here: hard-core bosons

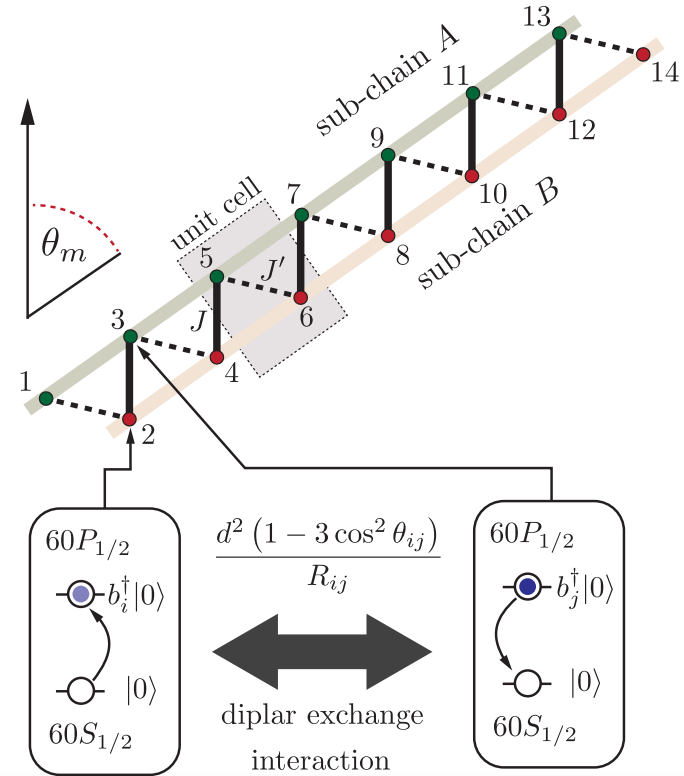
- full classification of SPT phases for interacting Bosons in one-dimension
F. Pollman, *et al.*, PRB 81, 064439 (2010)
X.-G. Wen, *et al.*, Science (2012).
- different SPT phases carry different projective symmetry on the edge
- classification in terms of twisted cohomology group



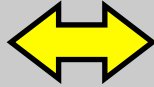
Symmetry protected topological phase

Bosonic version of the SSH model

- hardcore bosons in one dimension
- Jordan Wigner transformation: $b_j = e^{i\pi \sum_{k < j} c_k^\dagger c_k} c_j$
 - non-local transformation
 - longer range hopping give rise to interactions



non-interacting fermions



hard-core bosons

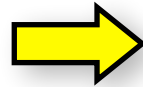
Fermions

- Hamiltonian

$$H = -J \sum_{i \in \text{even}} c_{i+1}^\dagger c_i - J' \sum_{i \in \text{odd}} c_{i+1}^\dagger c_i + h.c.$$

- chiral symmetry

$$\mathcal{S}_S = \prod_i \left[c_i + (-1)^i c_i^\dagger \right] K$$



Bosons

- Hamiltonian

$$H = -J \sum_{i \in \text{even}} b_{i+1}^\dagger b_i - J' \sum_{i \in \text{odd}} b_{i+1}^\dagger b_i + h.c.$$

- symmetry

$$\mathcal{S}_B = \prod_i \left[b_i^\dagger + b_i \right] K \quad \text{complex conjugation}$$

Symmetry protected topological phase

Protecting symmetries

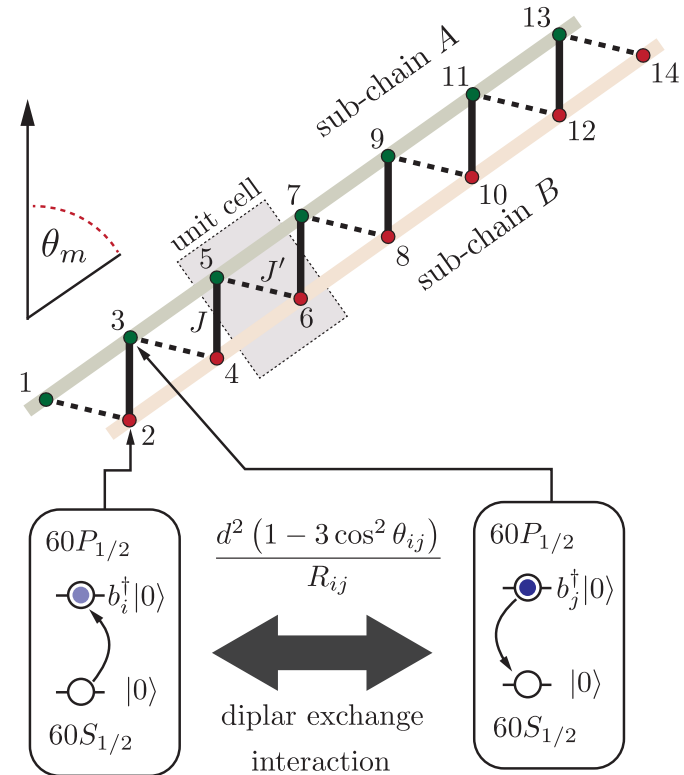
- particle conservation

- discrete operation $\mathcal{S}_B = \prod_i [b_i^\dagger + b_i] K$

- symmetry group $U(1) \times Z_2^T$

- Hamiltonian $[H, \mathcal{S}_B] = 0$

- allows for 4 different SPT phases
X.-G. Wen, *et al*, Science (2012).



SPT phase

- gapped ground state at half-filling

- four-fold ground state degeneracy

- zero energy edge states

Special point: $J' = 0$

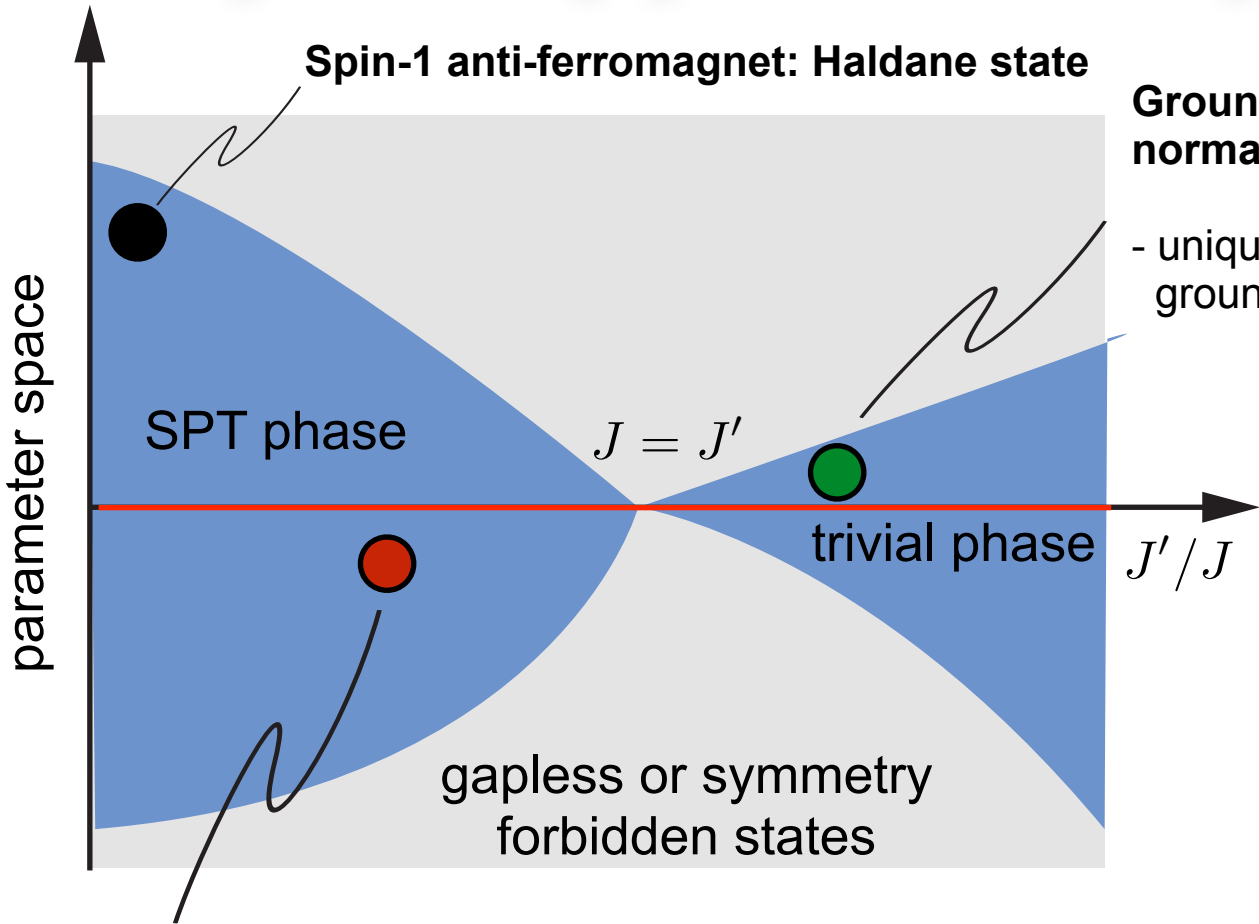
$$|m, m'\rangle = (b_1^\dagger)^m (b_L^\dagger)^{m'} \prod_{i \in \text{even}} \frac{1}{\sqrt{2}} (b_i^\dagger + b_{i+1}^\dagger) |0\rangle$$

Perturbations respecting the symmetry:

- arbitrary hoppings (also complex) $b_i^\dagger b_j + b_j^\dagger b_i$

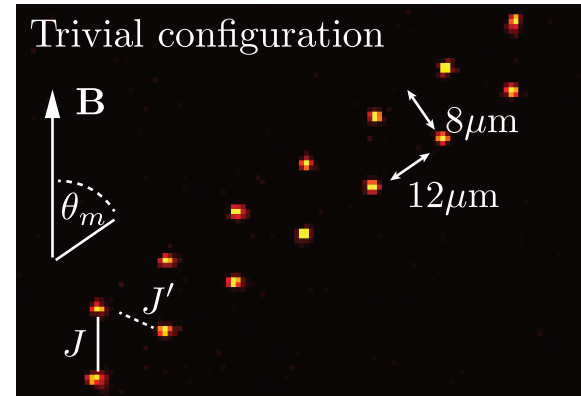
- interactions $(b_i^\dagger b_i - 1/2) (b_j^\dagger b_j - 1/2)$

Symmetry protected topological phase



Ground state in normal phase

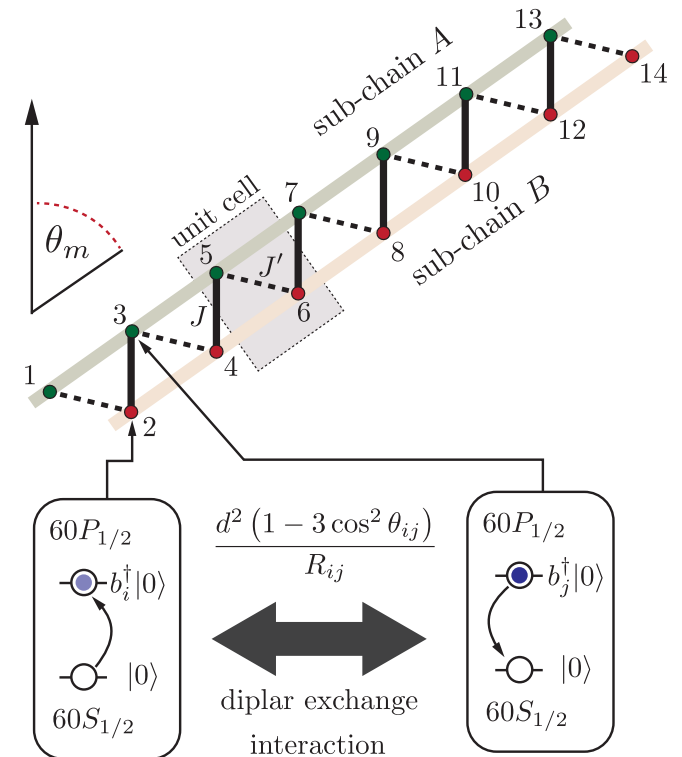
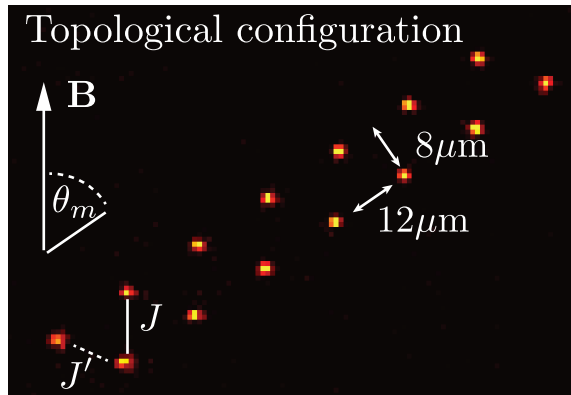
- unique gapped ground state



Ground state in SPT phase

- ground state degeneracy

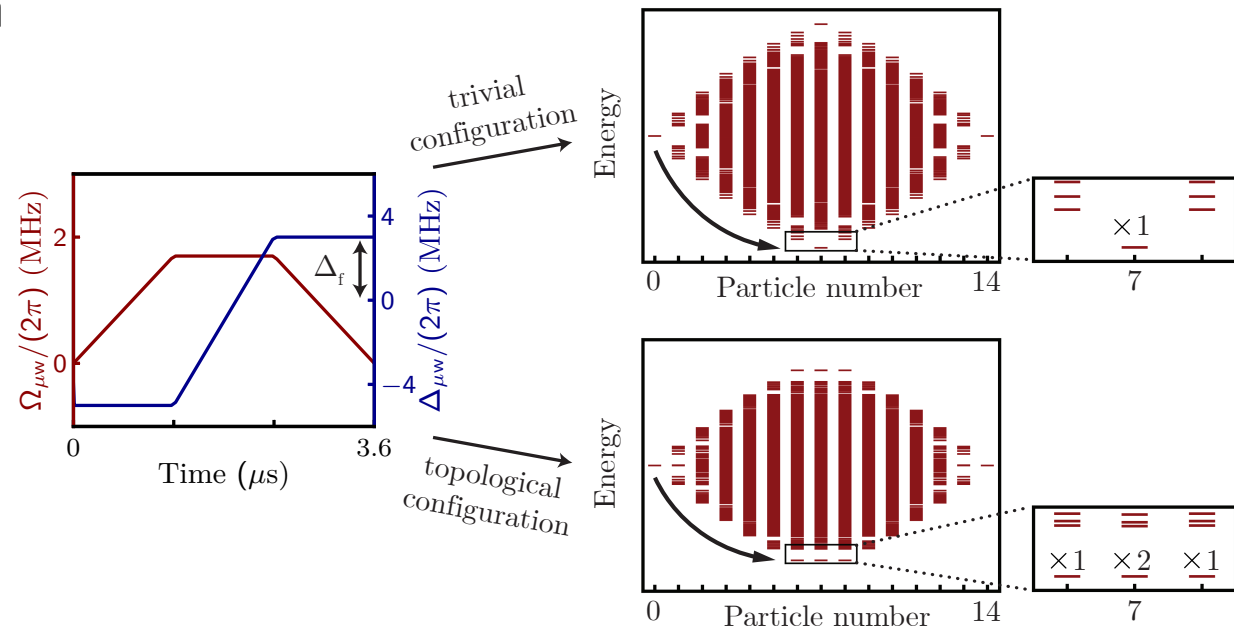
- edge states



Preparation of ground state at half-filling

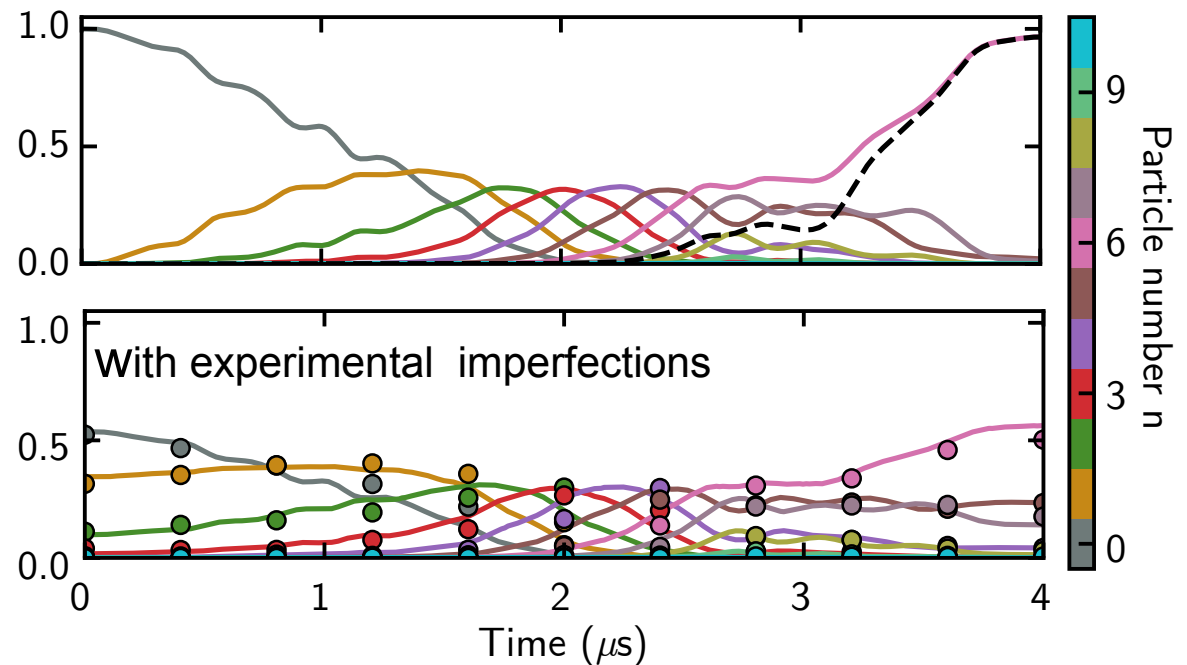
Adiabatic preparation of quantum many-body ground state

- prepare the ground state at half-filling
- highly efficient due to large gap
- ramping scheme motivated for perfect dimerization



Full numerical simulation

- efficient preparation of ground state
- high fidelity to prepare exactly ground state with $N/2$ excitations



Symmetry protected topological phase

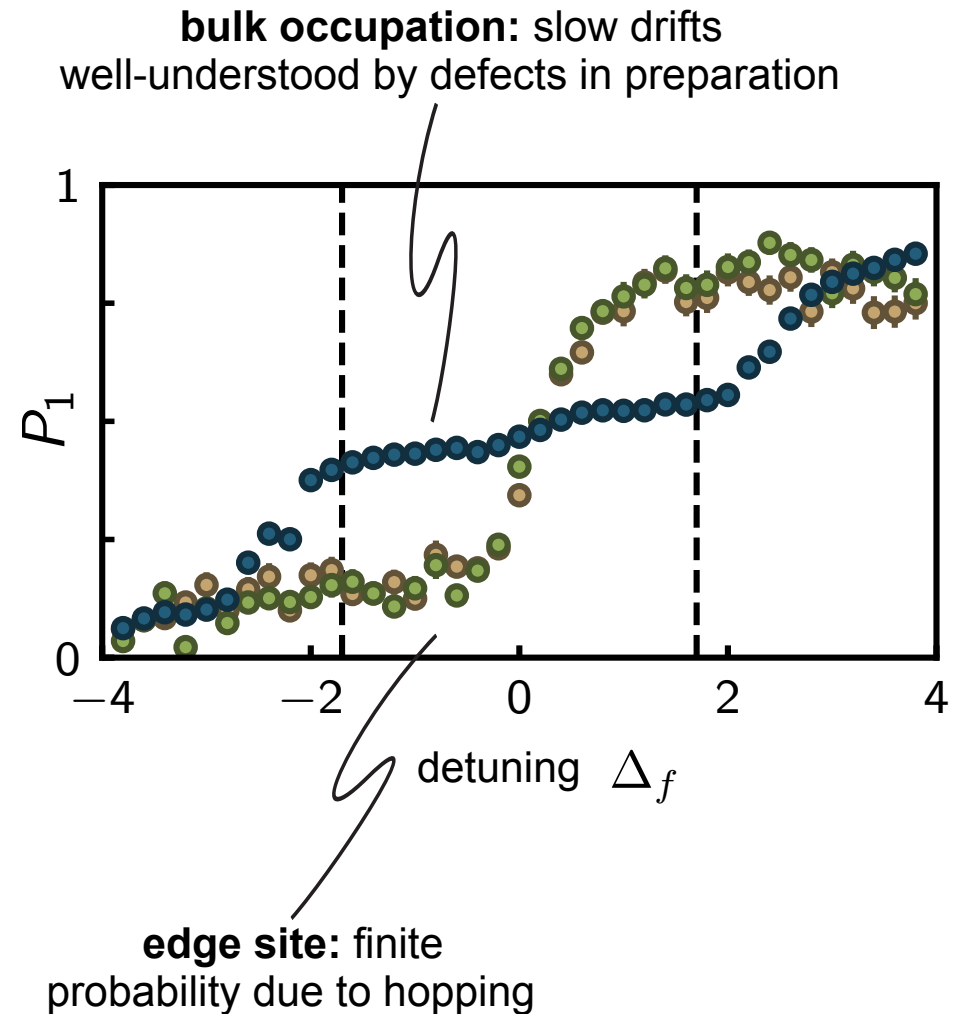
Characterization of topological phase

- half-filling in the bulk
- occupation of edge states depending on ramping procedure



ground state degeneracy

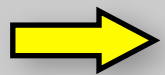
- zero-energy modes in spectroscopy
- robust to perturbation
- absence of any spontaneous symmetry breaking
- experimental detection of string order parameter



Symmetry protected topological phase

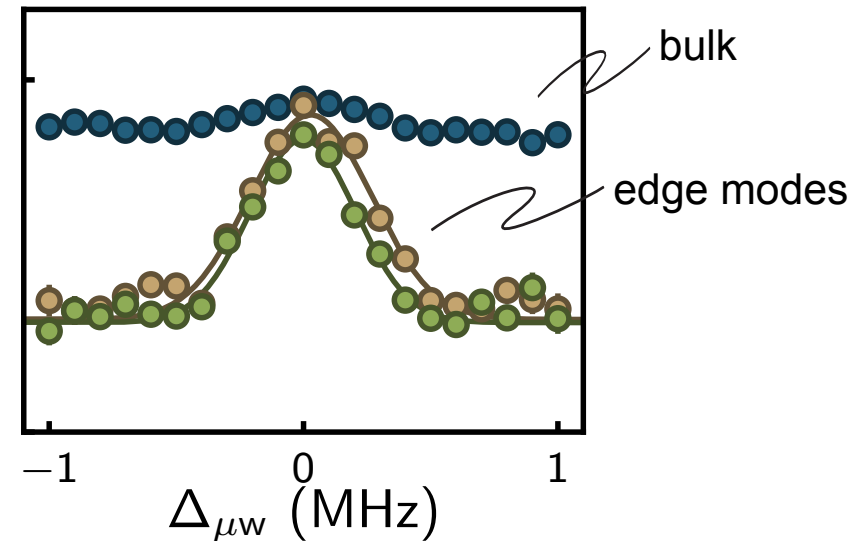
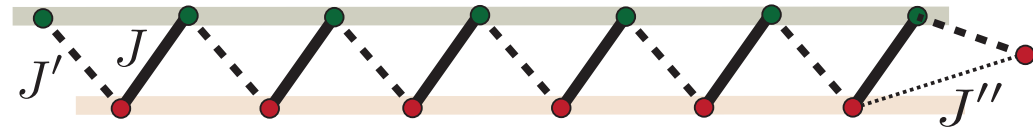
Characterization of topological phase

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ground state degeneracy

- zero-energy modes in spectroscopy
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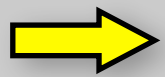


- edge modes excitation is at zero detuning
- in perturbed and unperturbed setup
- line-width well accounted for by finite pulse shape

Symmetry protected topological phase

Characterization of topological phase

- half-filling in the bulk
- occupation of edge states depending on ramping procedure



ground state degeneracy

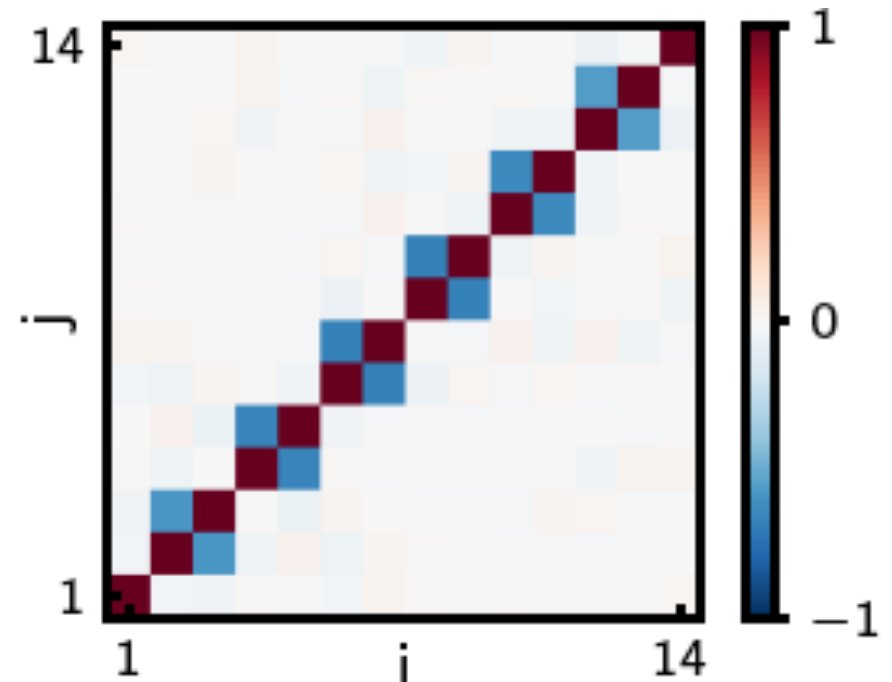
- zero-energy modes in spectroscopy
- robust to perturbation

- absence of any spontaneous symmetry breaking

- experimental detection of string order parameter

- correlation function

$$C_{ij} = \langle n_i n_j \rangle - \langle n_i \rangle \langle n_j \rangle$$

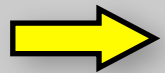


- one excitation shared between two dimers

Symmetry protected topological phase

Characterization of topological phase

- half-filling in the bulk
- occupation of edge states depending on ramping procedure



ground state degeneracy

- zero-energy modes in spectroscopy
- robust to perturbation

- absence of any spontaneous symmetry breaking

- experimental detection of string order parameter

- string order parameter

$$Z_i = 1 - b_i^\dagger b_i$$

$$C_{\text{string}}^z = - \left\langle Z_2 e^{i\frac{\pi}{2} \sum_{k=3}^{N-2} Z_k} Z_{N-1} \right\rangle$$

	C^z	C^x	C_{string}^z	C_{string}^x
Th. (no errors)	-0.96	0.98	0.78	0.88
Full simulation	-0.69(1)	0.68(2)	0.11(2)	0.10(2)
Experiments	-0.67(1)	0.48(2)	0.11(2)	0.05(2)

Symmetry protected topological phase

Does the different symmetry between bosons and fermions play a role?

Non-interacting Fermions

- SSH chain requires chiral symmetry

$$\mathcal{S}_S = \prod_i \left[c_i + (-1)^i c_i^\dagger \right] K$$

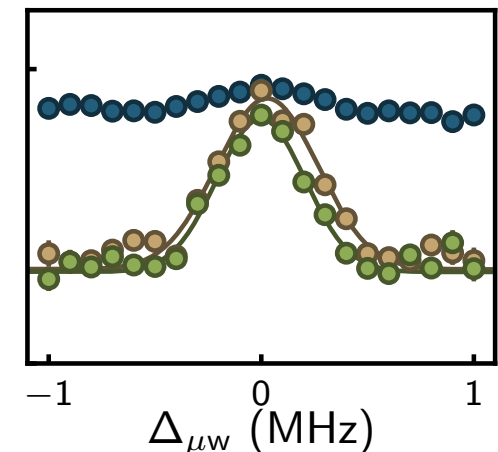
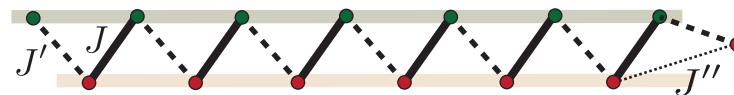
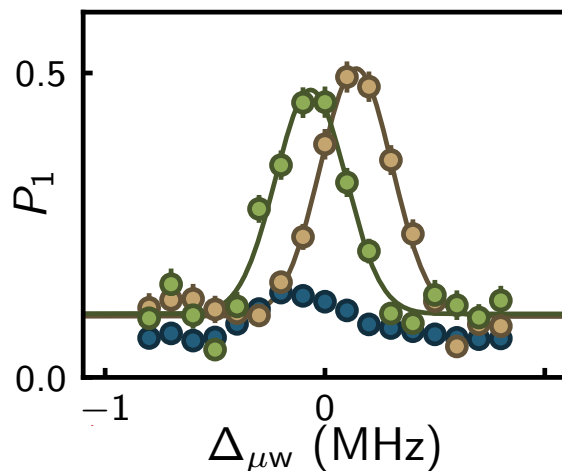
- only hopping between the sub-chains
- also required for single particle Hamiltonian

Interacting bosons

- protecting symmetry

$$\mathcal{S}_B = \prod_i \left[b_i^\dagger + b_i \right] K$$

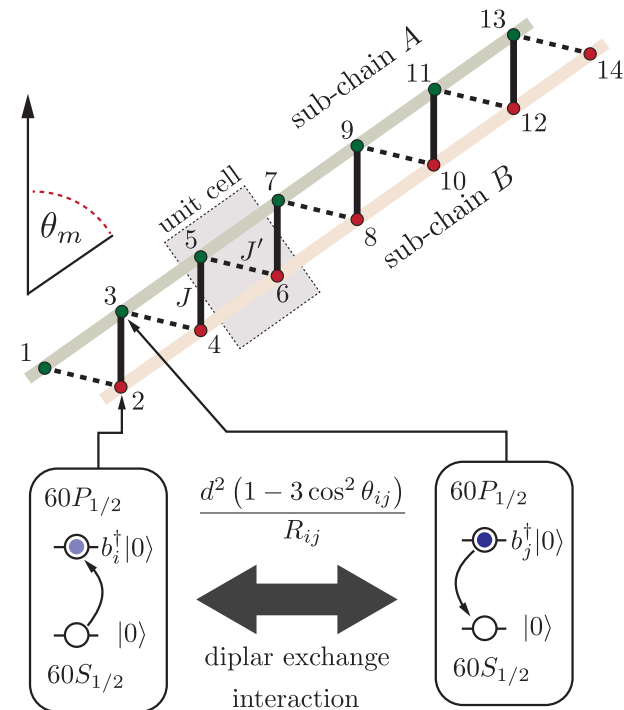
- allows arbitrary hopping
- robust ground state degeneracy by shifting the edge atom



Summary

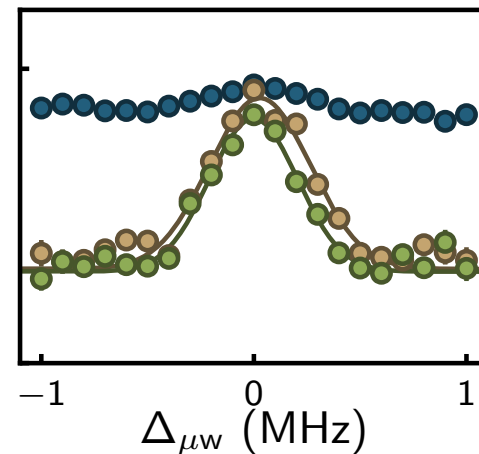
Edge modes in the SSH model

- implementation of the chiral symmetry
- observation of localized edge modes
- verification of exponential splitting of energy with system size
- benchmarking the experimental results with theoretical predictions



Symmetry protected topological phase

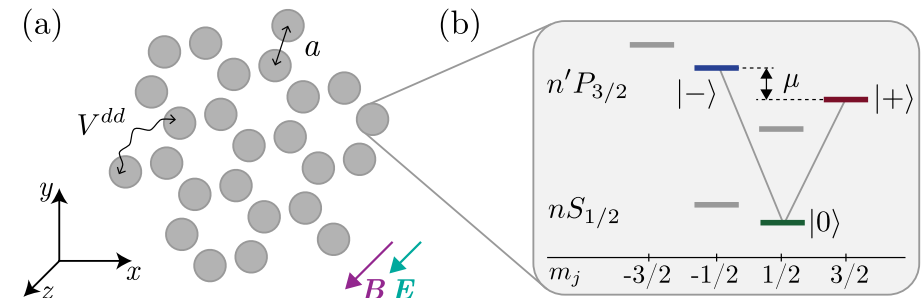
- ground state of the interacting many-body system at half filling
- first realization of a symmetry protected topological phase in artificial
- quantum simulation of novel states of matter



Outlook

Topological band structure

- dipolar exchange interactions
- spin-orbit coupling due to anisotropy of dipole-dipole interaction
- Chern number $C=1$ or $C=2$
- probing dynamics of edge states



quantum many-body states?

fractional bosonic Chern insulators?

dynamical preparation?

competition with losses?

