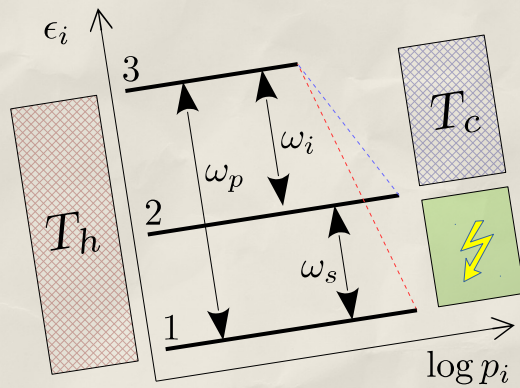


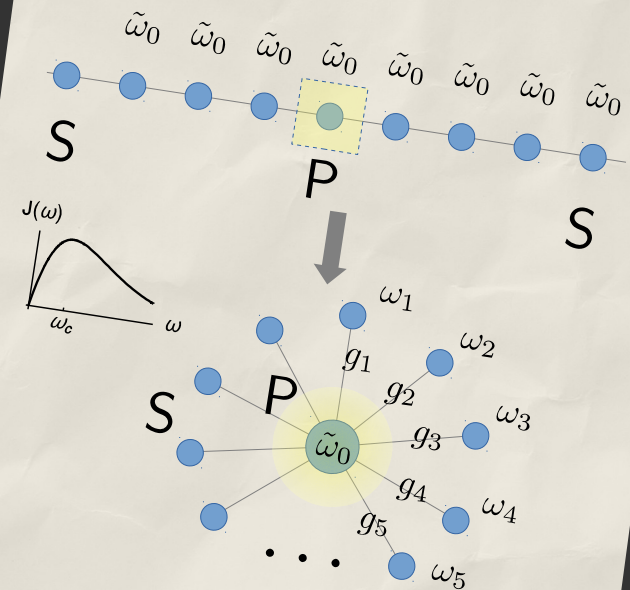
DOUBLE FEATURE!

Quantum thermal engineering



AND

Quantum thermometry



KITP USBC 08/06/18



QUANTUM THERMOMETRY



University of
Nottingham

UK | CHINA | MALAYSIA

LUIS A. CORREA



Mohammad Mehboudi

ICFO: The Institute
of Photonic Sciences



Anna Sanpera

Universitat Autònoma
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University of Nottingham

Christos Charalambous

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Senaida Hernández-Santana

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Karen V. Hovhannisyan

Aahrus University

Aniello Lampo

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Maciej Lewenstein

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Martí Perarnau-Llobet

Max Planck Institut für Quantenoptik

QUANTUM THERMOMETRY

Estimating temperature for **quantum technologies**:

- Nanoscale spatial resolution.
- Minimally disturbing techniques.
- High precision at low temperatures.



NANOSCALE THERMOMETRY

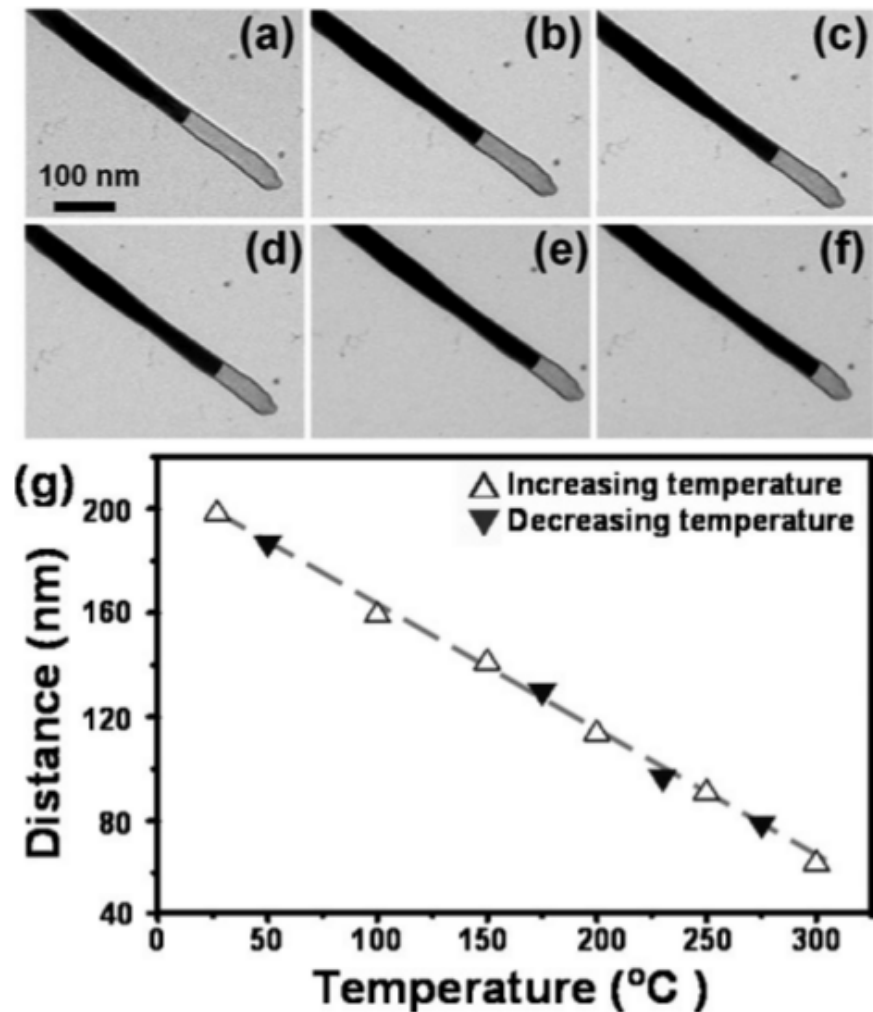


Figure 3. a–f) TEM images of a selected Pb/ZnO nanocable, with Pb filling about 16 μm in length, heated from room temperature to 25, 100, 150, 200, 250, and 300 $^{\circ}\text{C}$, respectively, recorded by video camera in the transmission electron microscope. g) Plot of cavity length versus heating temperature.

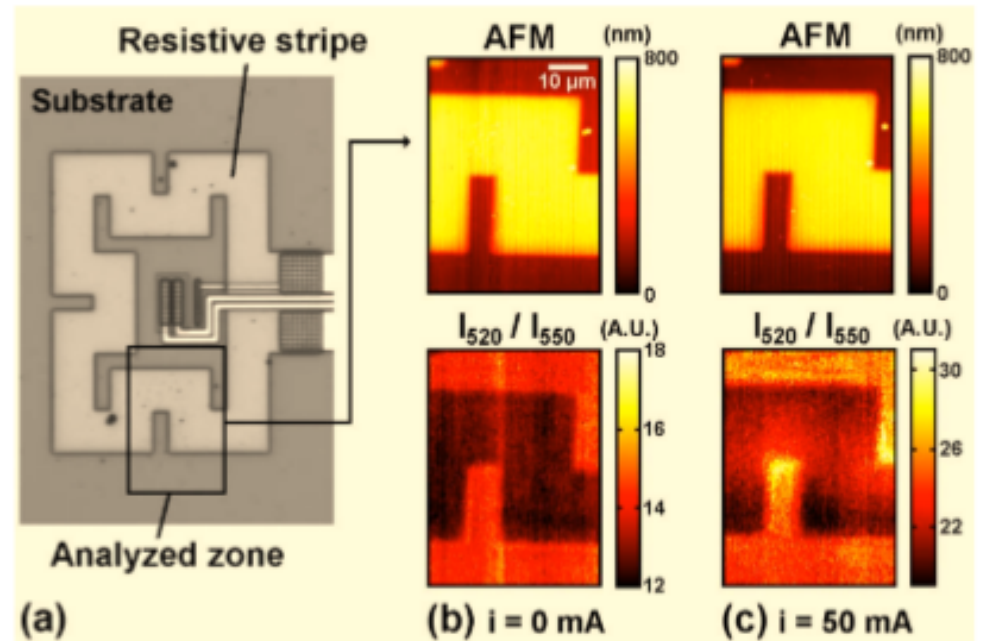


FIG. 3. (Color online) (a) Optical micrograph of the microelectronic circuit; (b) topography and fluorescence ratio images of the structure when no current is passing through the stripe; (c) topography and fluorescence ratio images when a current of ~ 50 mA is passing through the structure. The bright (hot) zones are clearly visible on the stripe. The image size is $45 \times 60 \mu\text{m}^2$.

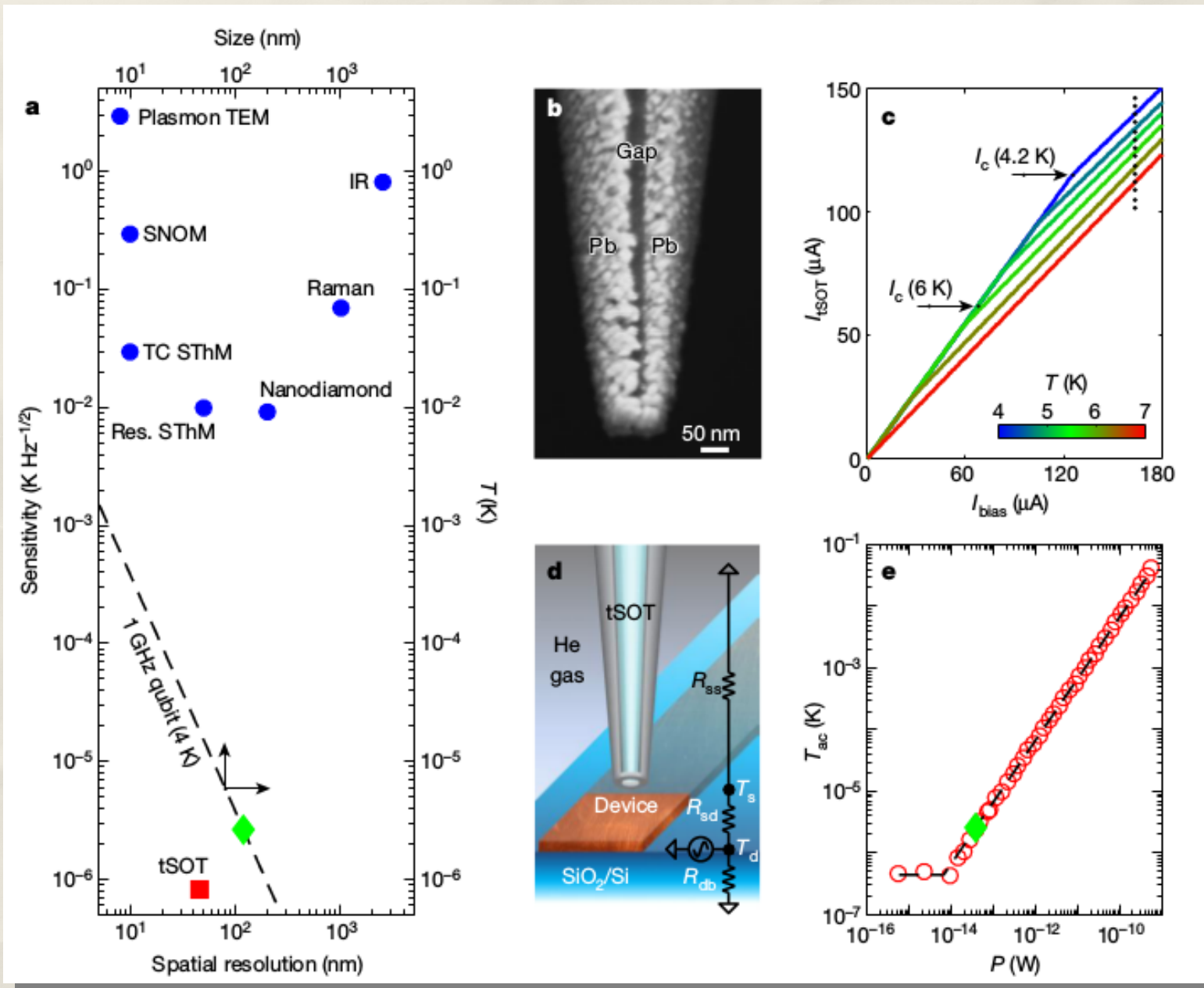
L. Aigouy *et al.*, *Appl. PhLeadys. Lett.* **87**, 184105 (2005)

QUANTUM THERMOMETRY

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TOWARDS QUANTUM THERMOMETRY...



OUTLINE

~~1. Motivation~~

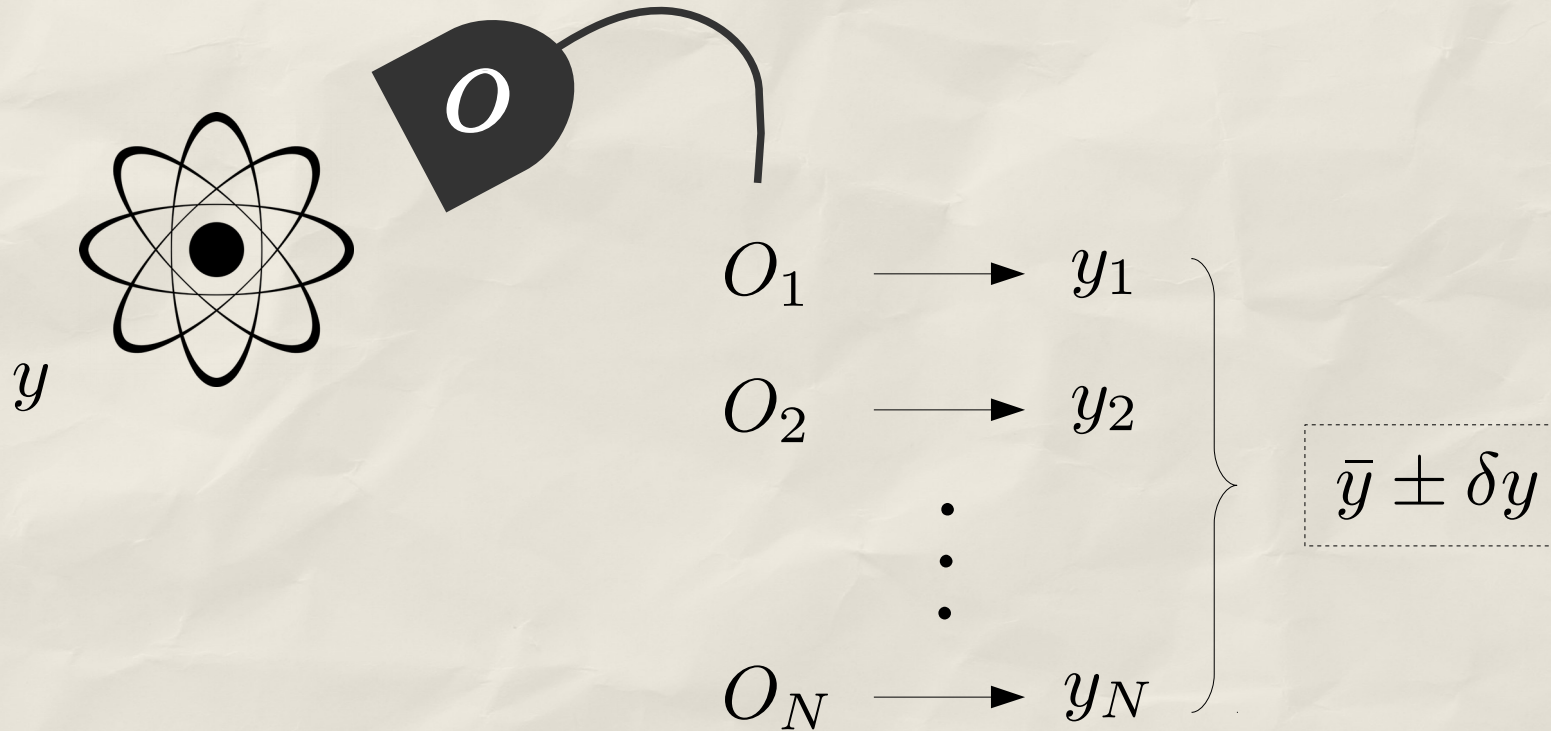
2. Thermal sensitivity

3. Equilibrium quantum thermometry

4. Quantum thermometry out of equilibrium

5. Low-temperature scaling of thermal sensitivity

QUANTUM ESTIMATION THEORY



$$y = \bar{y}$$

$$N \rightarrow \infty$$

$$\delta y \geq \frac{1}{\sqrt{N \mathcal{F}_y(\mathbf{O})}}$$

Cramér-Rao
bound

QUANTUM ESTIMATION THEORY

$$\frac{\partial_y \langle \mathbf{O} \rangle}{\Delta \mathbf{O}} \leq \mathcal{F}_y(\mathbf{O}) \leq F_y \equiv \sup_{\mathbf{O}} \mathcal{F}_y(\mathbf{O})$$

Error propagation formula Classical Fisher information Quantum Fisher information

$$y = \bar{y}$$
$$N \rightarrow \infty$$

$$\delta y \geq \frac{1}{\sqrt{N \mathcal{F}_y(\mathbf{O})}}$$

Cramér-Rao bound

OUTLINE

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SENSITIVITY AND HEAT CAPACITY

$$\rho_T = \mathcal{Z}^{-1} e^{-\mathbf{H}/T}$$

$$F_T(\rho_T) = \frac{(\Delta \mathbf{H})^2}{T^4}$$

$$F_T(\rho_T) = \frac{C(T)}{T^2}$$

$$C(T) \equiv \frac{d\langle \mathbf{H} \rangle}{dT} = \frac{(\Delta \mathbf{H})^2}{T^2}$$

T. Jahnke *et al.*, *PRE* **83**, 011109 (2011)

LAC *et al.*, *PRL* **114**, 220405 (2015)

SENSITIVITY AND HEAT CAPACITY



$$\Delta T \leq \frac{1}{\sqrt{F_T(\rho_T)}}$$

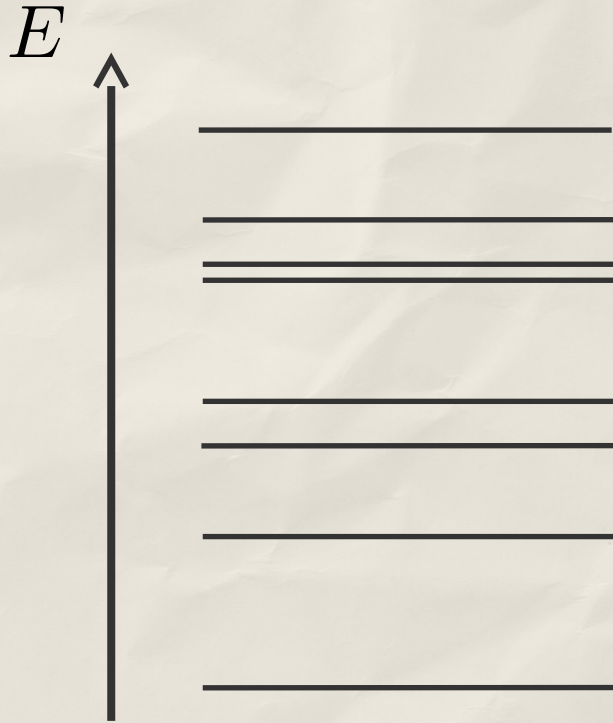
$$F_T(\rho_T) = \frac{C(T)}{T^2}$$

$$\left(\frac{T}{\Delta T}\right)^2 \leq C(T)$$

T. Jahnke *et al.*, *PRE* **83**, 011109 (2011)

LAC *et al.*, *PRL* **114**, 220405 (2015)

THE BEST EQUILIBRIUM THERMOMETER

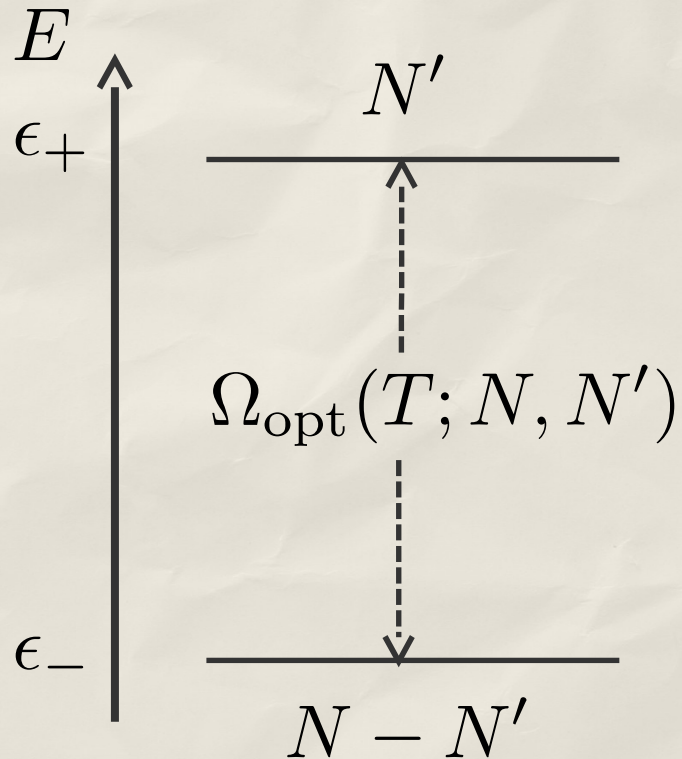


$$\partial_{\epsilon_i} \Delta H = 0$$

Reeb & Wolf, *IEEE Trans. Inf. Theor.*, 1458 (2015)

LAC et al., *PRL* **114**, 220405 (2015)

THE BEST EQUILIBRIUM THERMOMETER

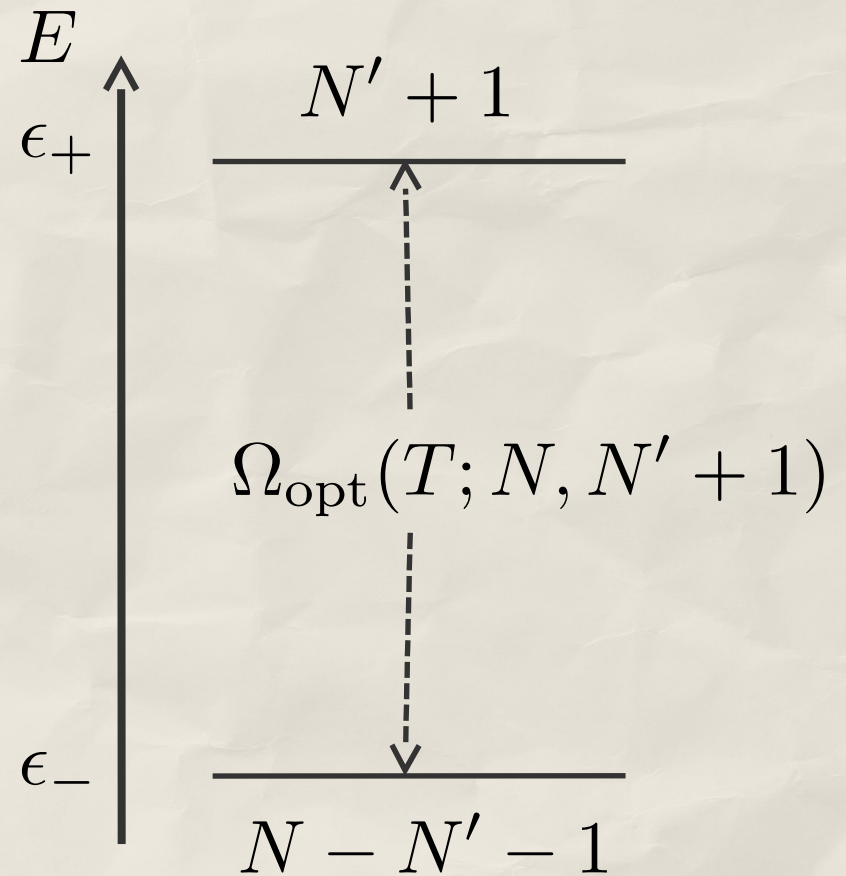
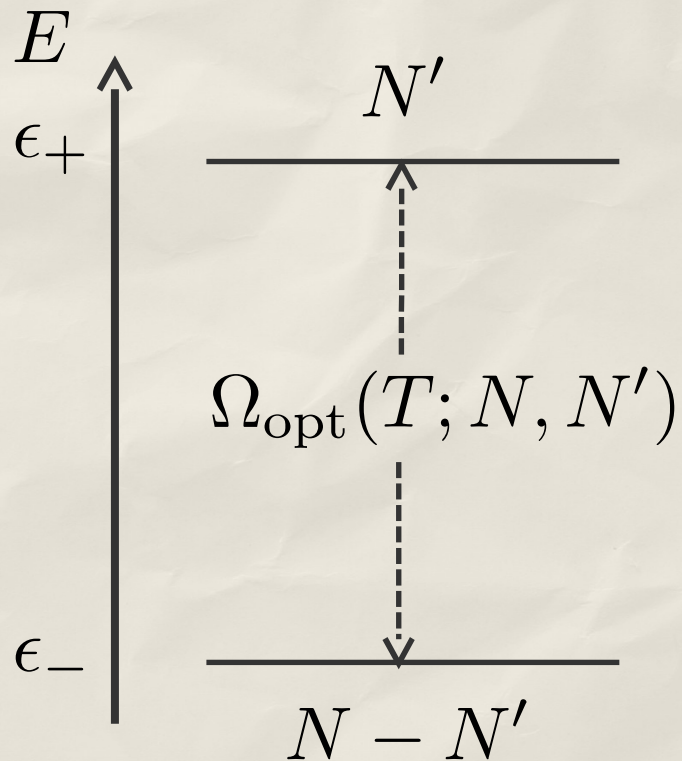


$$\partial_{\epsilon_i} \Delta H = 0$$

Reeb & Wolf, *IEEE Trans. Inf. Theor.*, 1458 (2015)

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THE BEST EQUILIBRIUM THERMOMETER

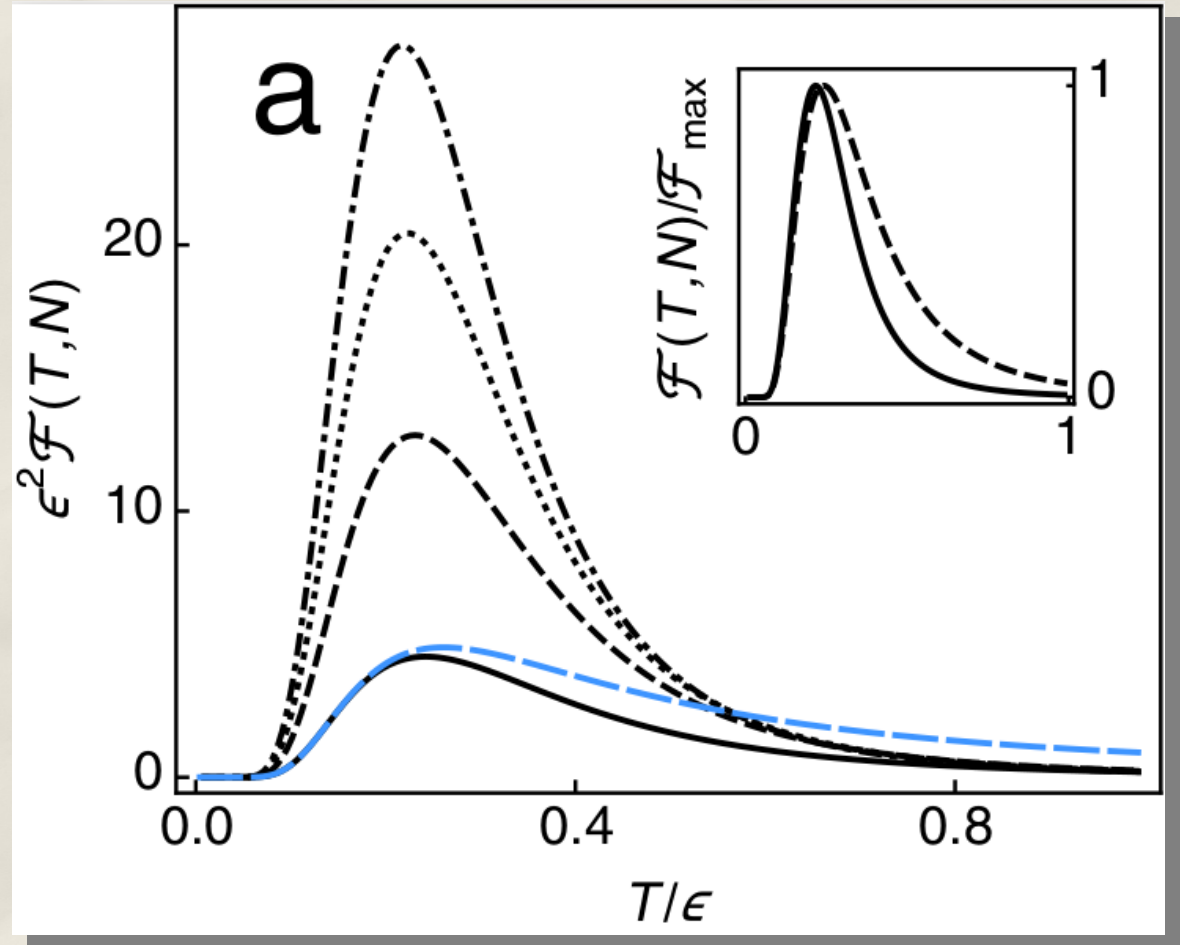
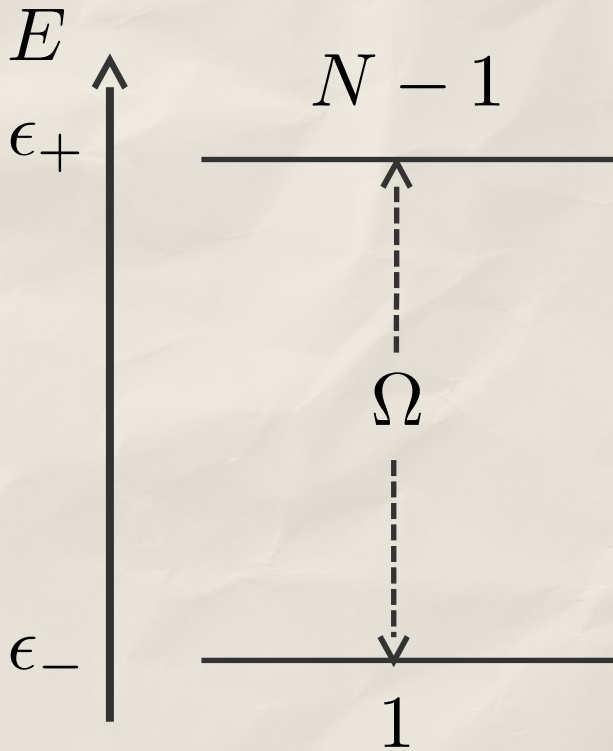


$$C(T; N, N') \leq C(T; N, N' + 1)$$

Reeb & Wolf, *IEEE Trans. Inf. Theor.*, 1458 (2015)

LAC et al., *PRL* **114**, 220405 (2015)

THE BEST EQUILIBRIUM THERMOMETER



OUTLINE

~~1. Motivation~~

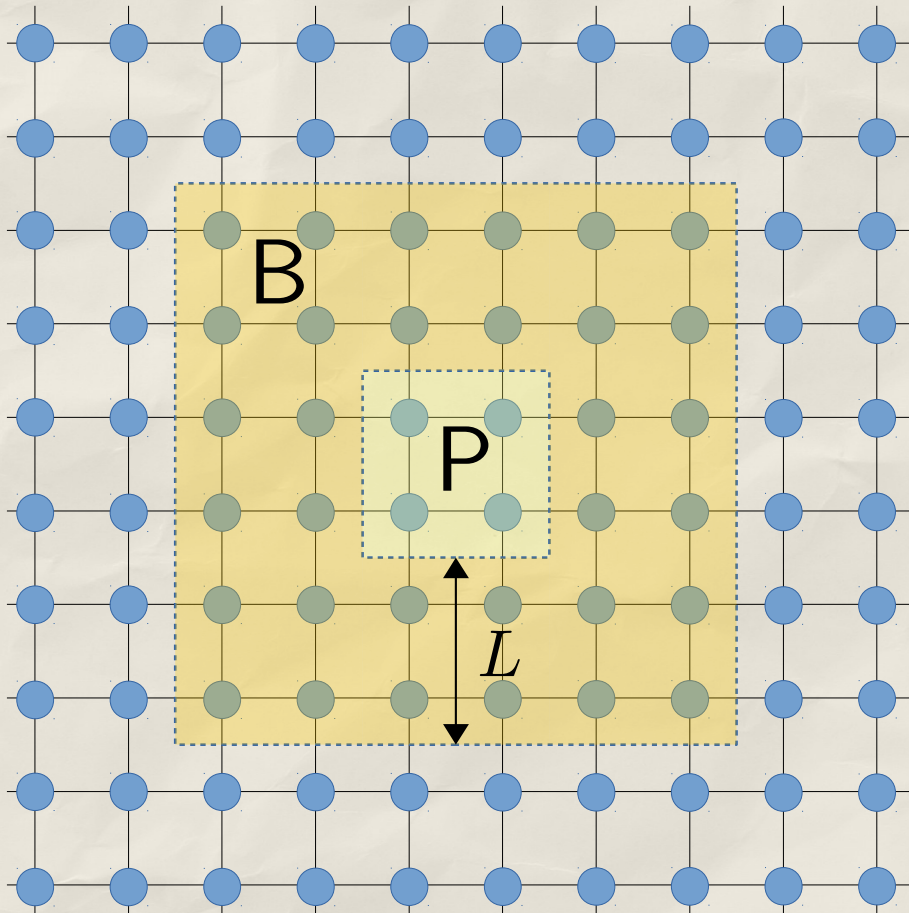
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SENSITIVITY AND LOCAL HEAT CAPACITY



- * $|B(L)| \ll |P|$

- * $L > 2\xi(T)$

- * $T > T^*$



$$F_T(\text{trs } \varrho_T) \simeq \frac{(\Delta H_P)^2}{T^4}$$

A. De Pasquale *et al.*, *Nat. Commun.* **7**, 12782 (2016)

G. De Palma *et al.*, *PRA* **95**, 052115 (2017)

SENSITIVITY AND DRESSED HAMILTONIAN

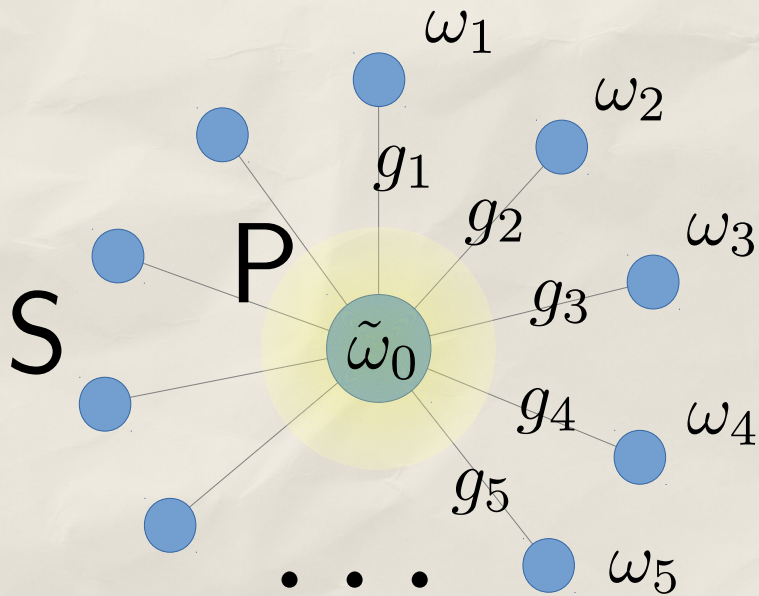
$$\mathcal{Z}_P^* \equiv \frac{\text{tr} e^{-\beta(\mathbf{H}_P + \mathbf{H}_S + \mathbf{H}_{\text{int}})}}{\text{tr} e^{-\beta \mathbf{H}_S}}$$

$$-\partial_\beta \log \mathcal{Z}_P^* = U_P(T) = \langle \mathbf{E}_P^*(T) \rangle$$

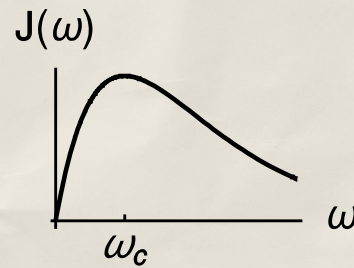
$$Q(\text{tr}_S \boldsymbol{\rho}_T, \mathbf{E}_P^*(T)) \equiv -\frac{1}{2} \int_0^1 da \text{tr} \{ [\mathbf{E}_P^*, \text{tr}_S \boldsymbol{\rho}_T^a] [\mathbf{E}_P^*(T), \text{tr}_S \boldsymbol{\rho}_T^{1-a}] \}$$

$$F_T(\text{tr}_S \boldsymbol{\rho}_T) = \frac{(\Delta \mathbf{E}_P^*)^2 - Q(\text{tr}_S \boldsymbol{\rho}_T, \mathbf{E}_P^*)}{T^4}$$

SENSITIVITY ENHANCED BY DISSIPATION



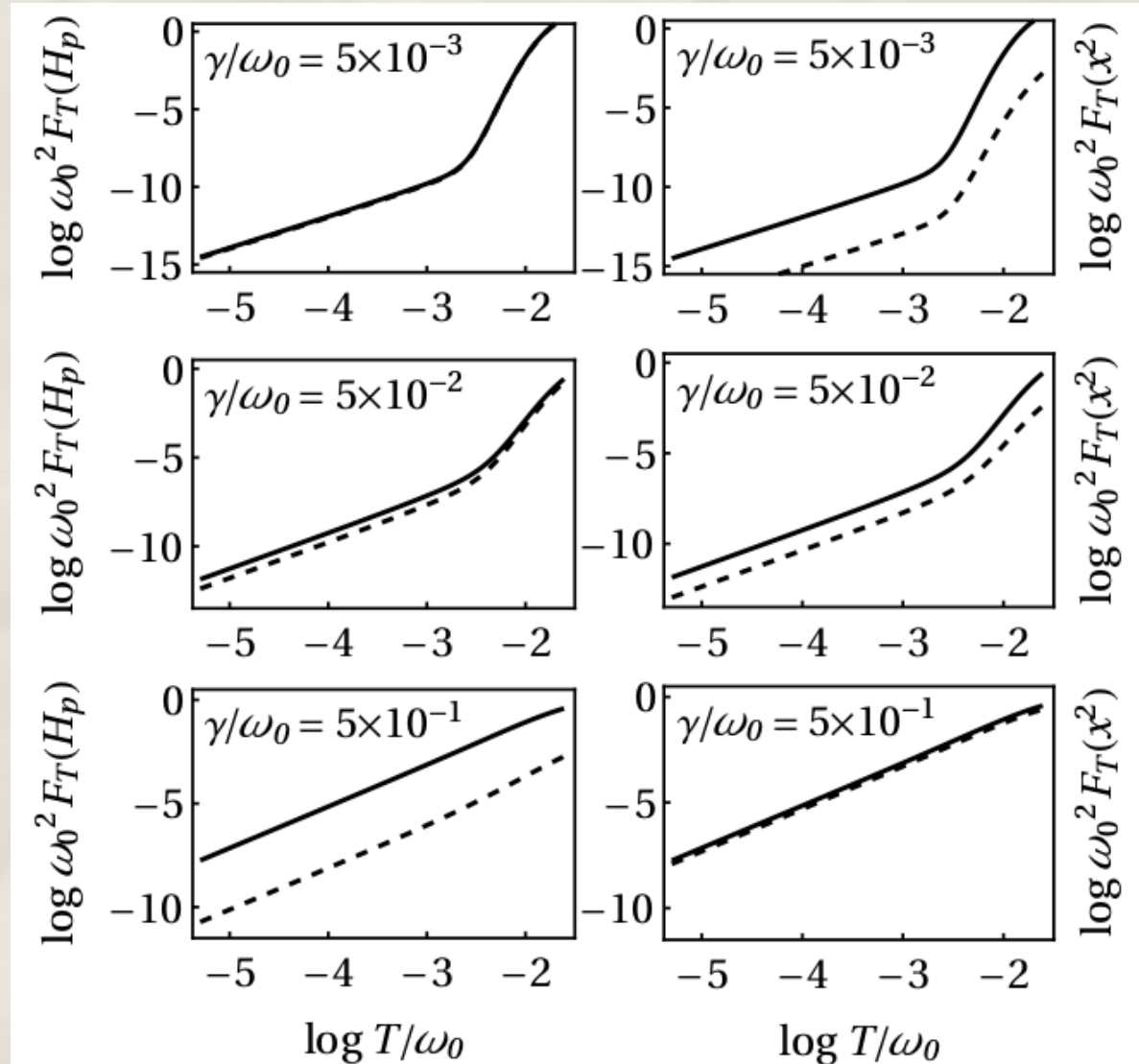
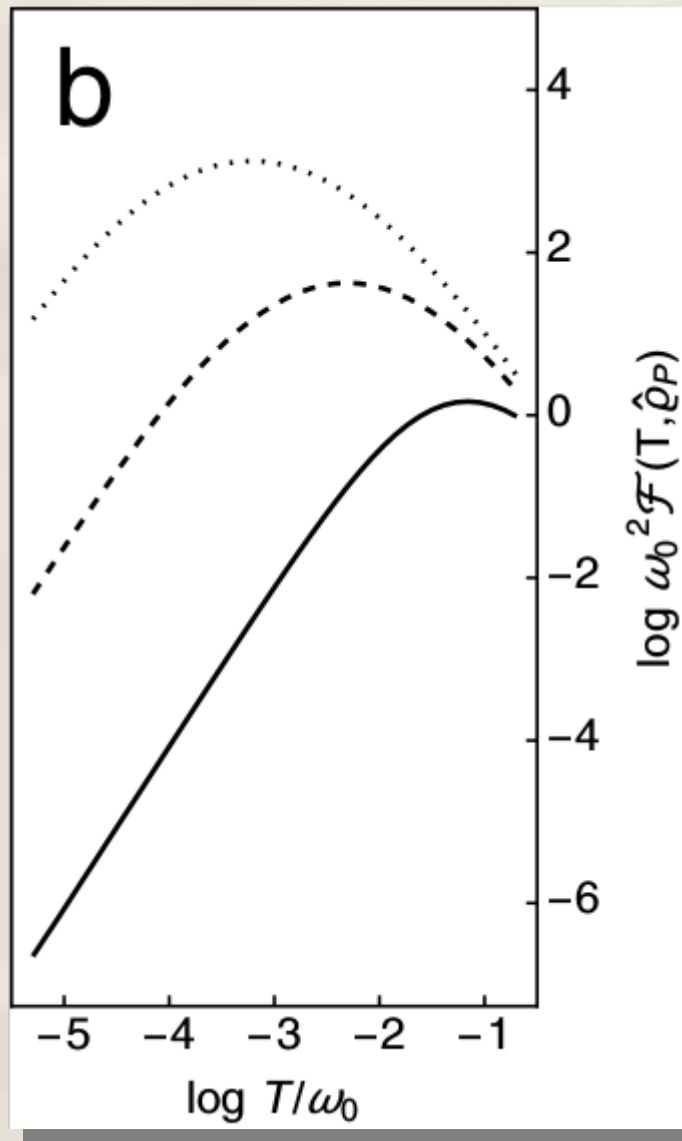
$$J(\omega) \equiv \pi \sum_{\mu} \frac{g_{\mu}^2}{2m_{\mu}\omega_{\mu}} \delta(\omega - \omega_{\mu})$$



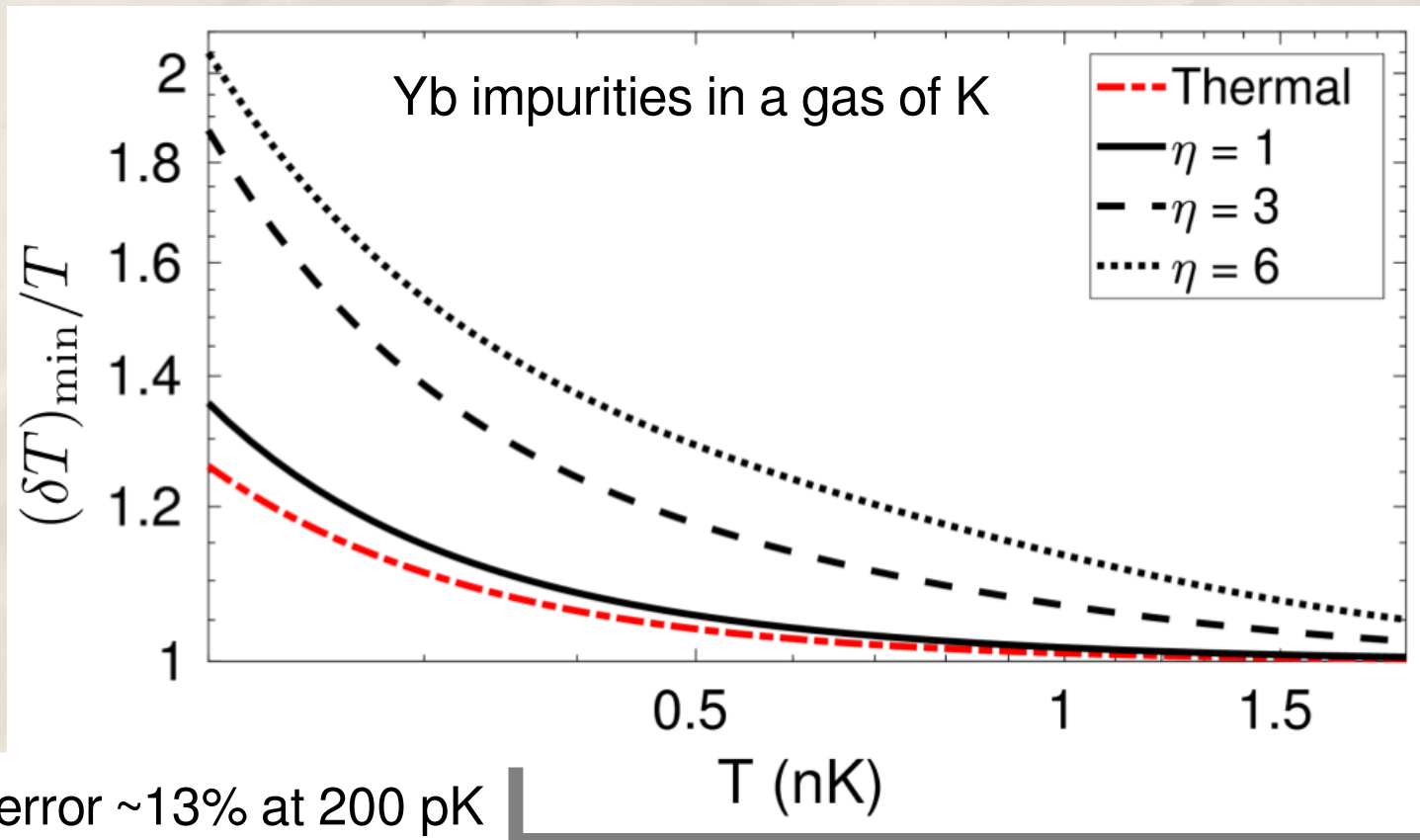
$$\ddot{x}(t) + \tilde{\omega}_0^2 x(t) - x(t) \star \chi(t) = F(t)$$

Quantum Langevin equation

SENSITIVITY ENHANCED BY DISSIPATION



POLARON THERMOMETRY IN A BEC



Relative error $\sim 13\%$ at 200 pK
with 100 measurements !

$$N = 5000$$

$$\omega_B = 200\pi \text{ Hz}$$

A. Lampo *et al.*, *arXiv:1803.08946* (2018)

M. Mehboudi *et al.*, (in preparation)

QUANTUM THERMOMETRY

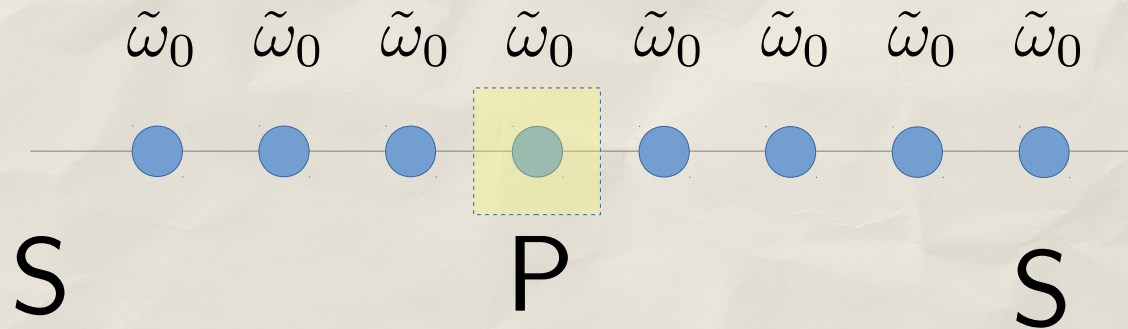
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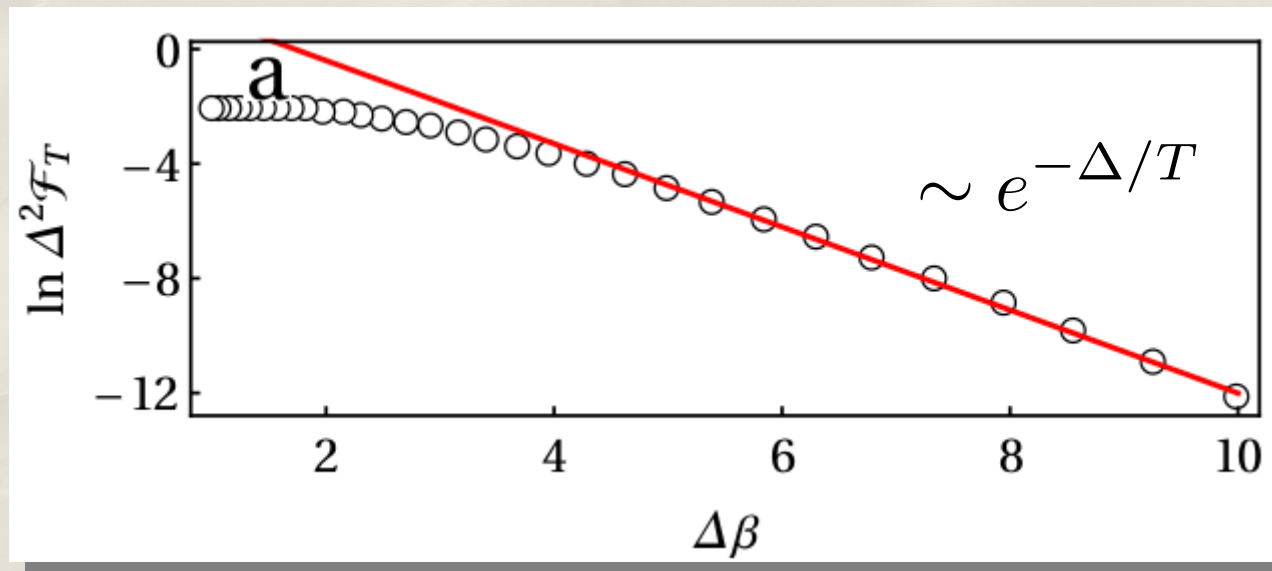
OUTLINE

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- ~~2. Thermal sensitivity~~
- ~~3. Equilibrium quantum thermometry~~
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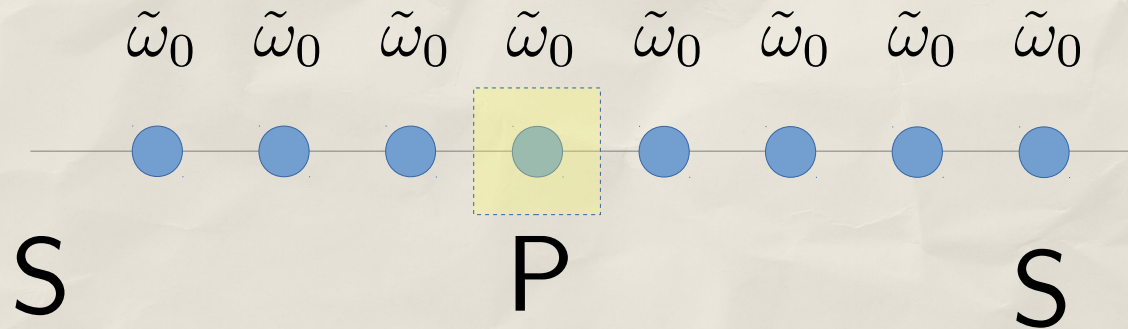
GAPPED VS. GAPLESS



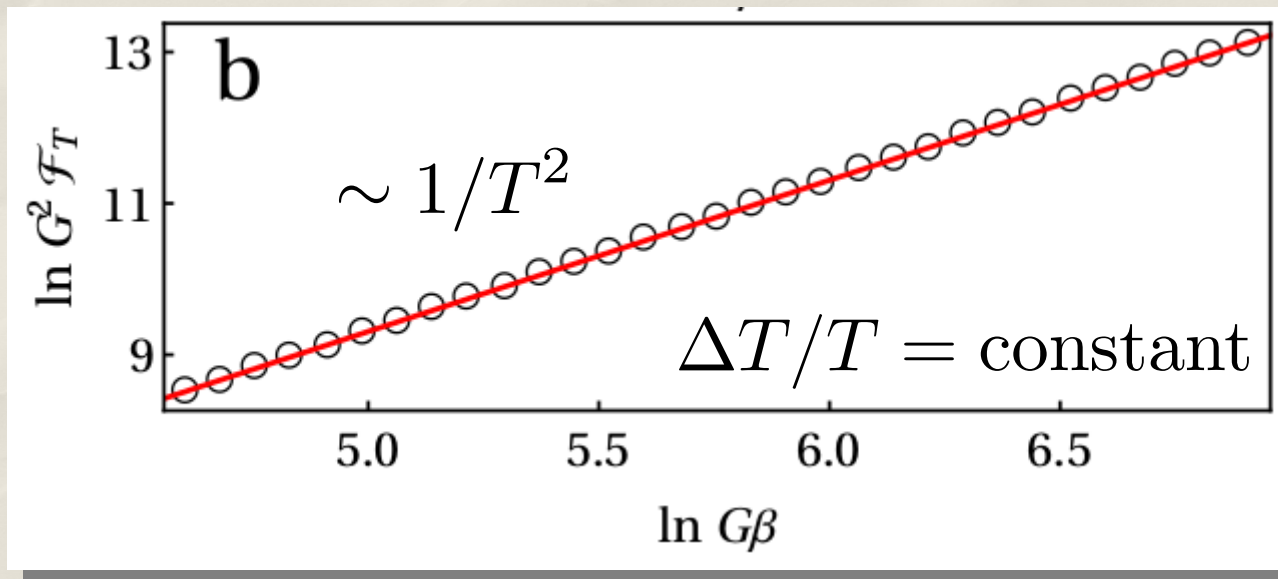
Gapped translationally invariant harmonic chain



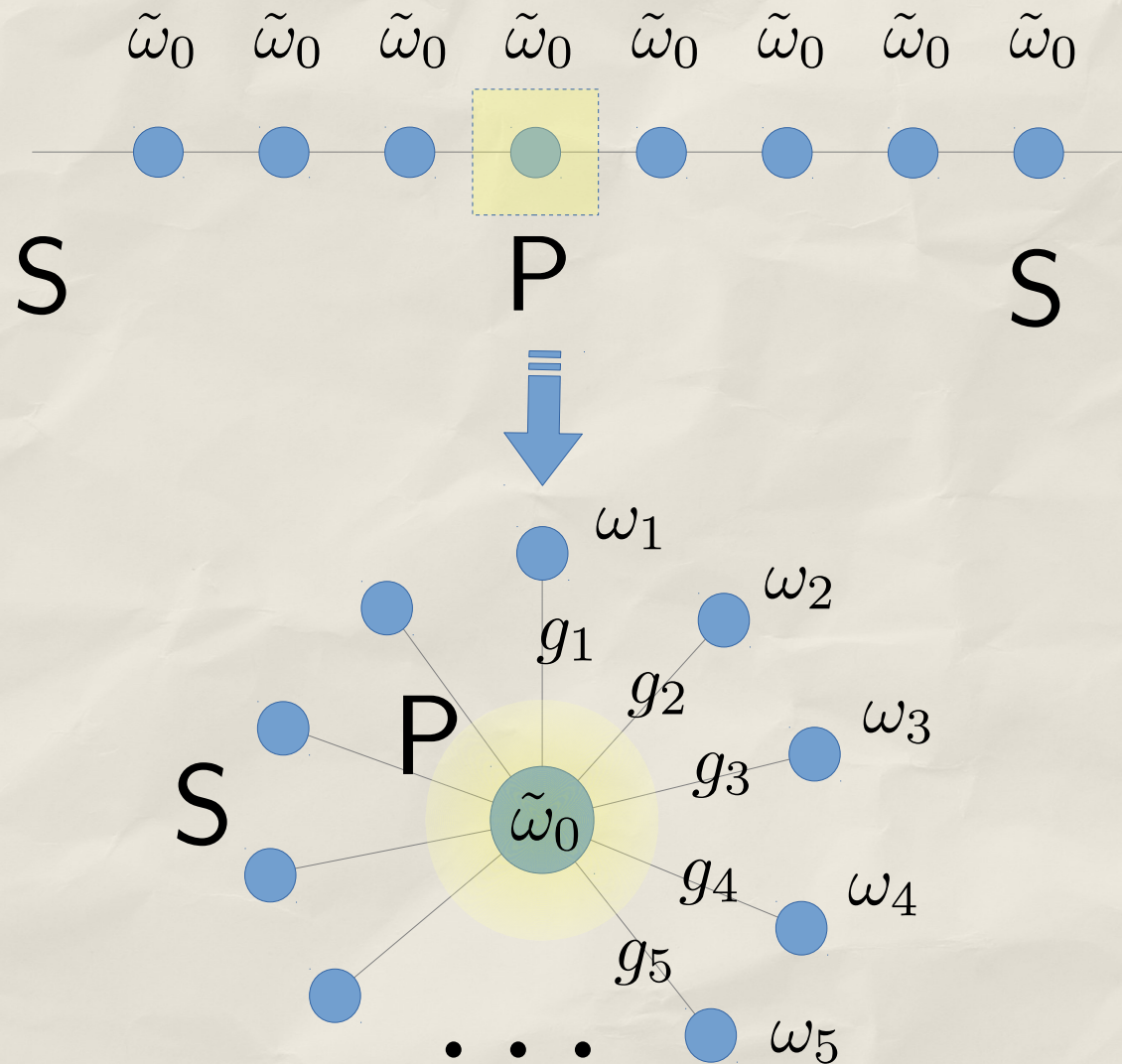
GAPPED VS. GAPLESS



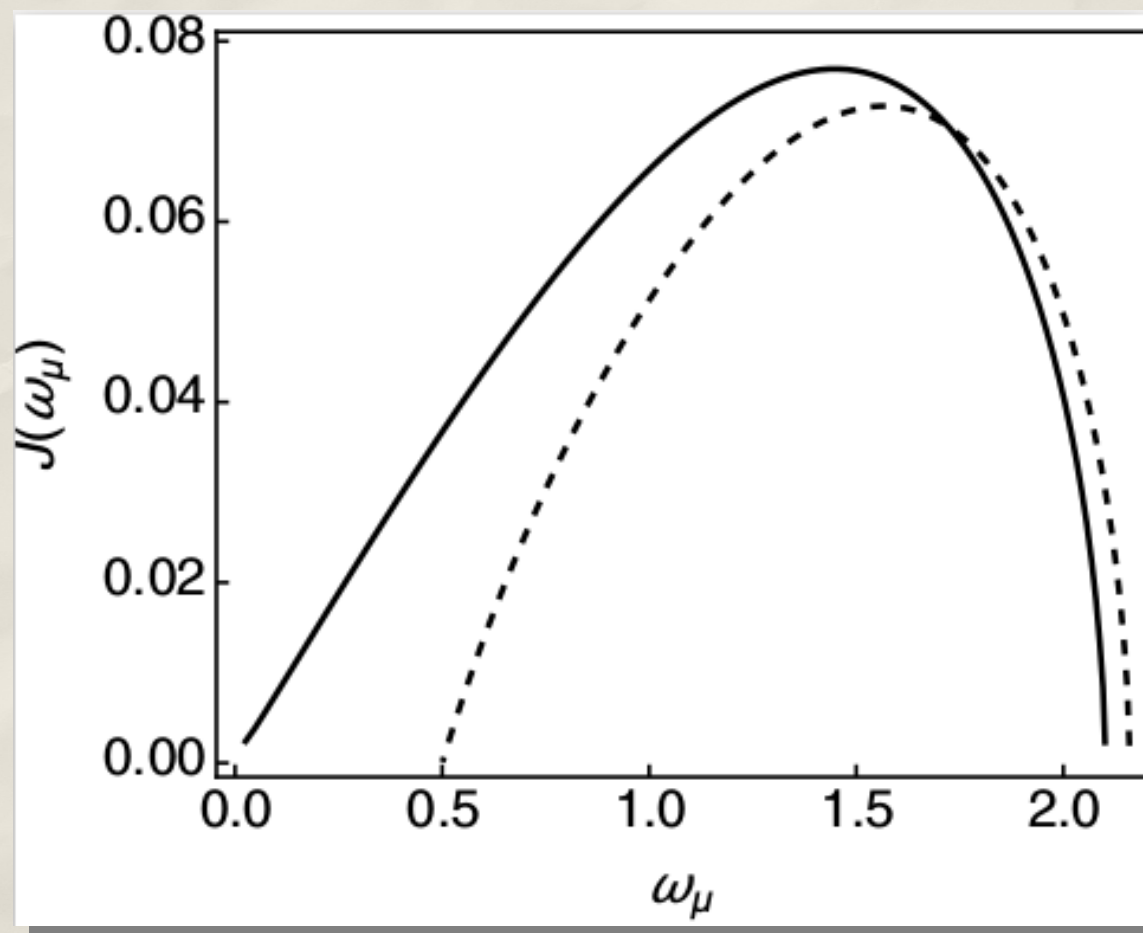
Gapless translationally invariant harmonic chain



AN OPEN SYSTEM APPROACH



AN OPEN SYSTEM APPROACH



WRAP-UP

Equilibrium quantum thermometry

- * The **heat capacity** places the ultimate bound on thermometric precision.
- * Sub-optimal probes can prove **versatile**.

Non equilibrium quantum thermometry

- * **Dissipation** can be exploited as a resource at low T .
- * The bose-polaron model is a good platform to study precise **non-demolition** thermometry on a BEC.

Low-temperature thermometry

- * Local thermometry in a many-body system can be **exponentially or polynomially** inefficient, depending on whether or not the system is gapped.