

.AI

Glassy Slowdown and Amorphous Order

Sho Yaida

Complex landscapes in everyday life



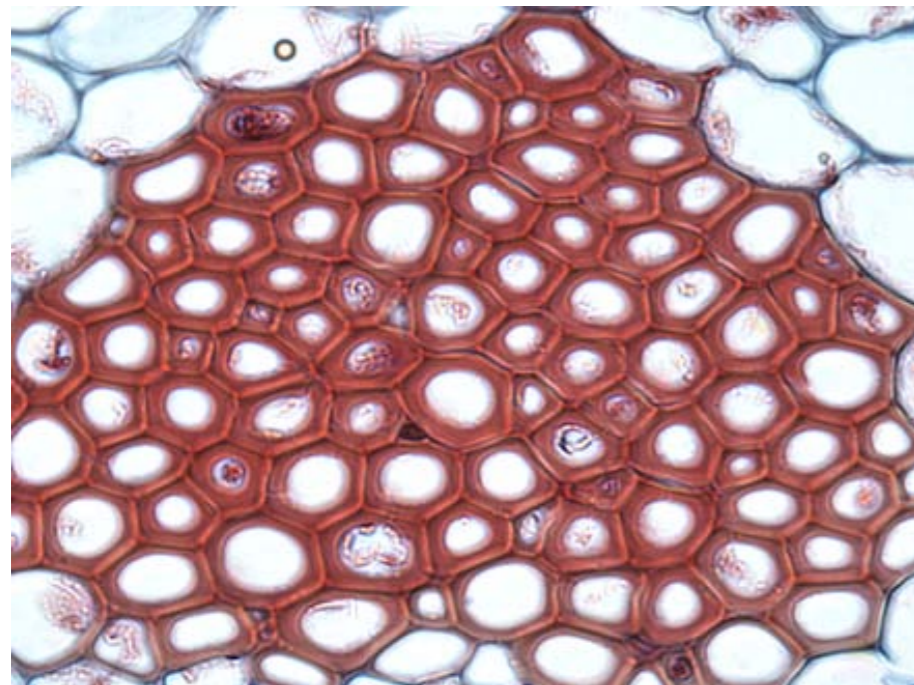
Glass



Plastic



Sand



Tissue



Ecosystem



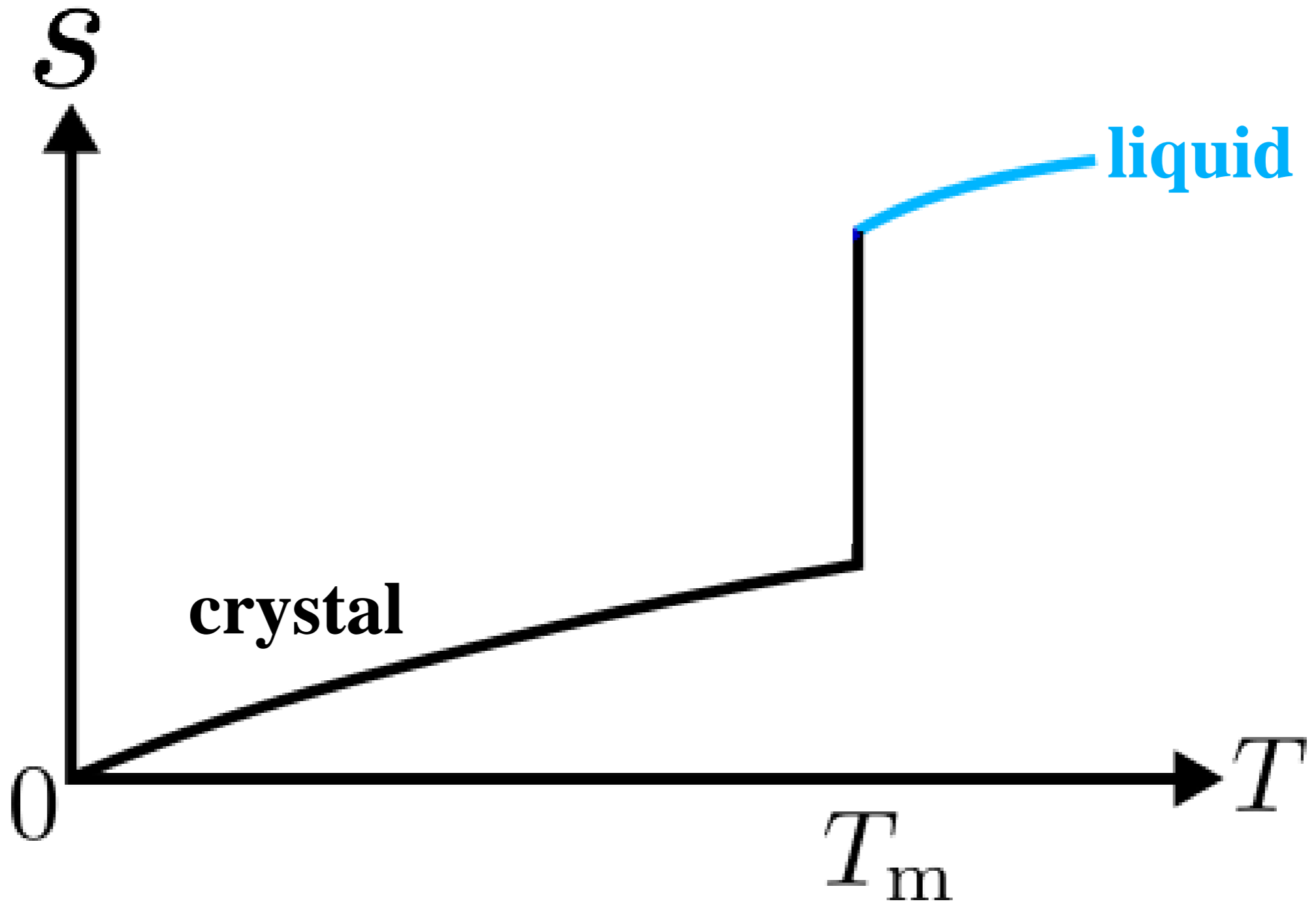
Learning

Outline

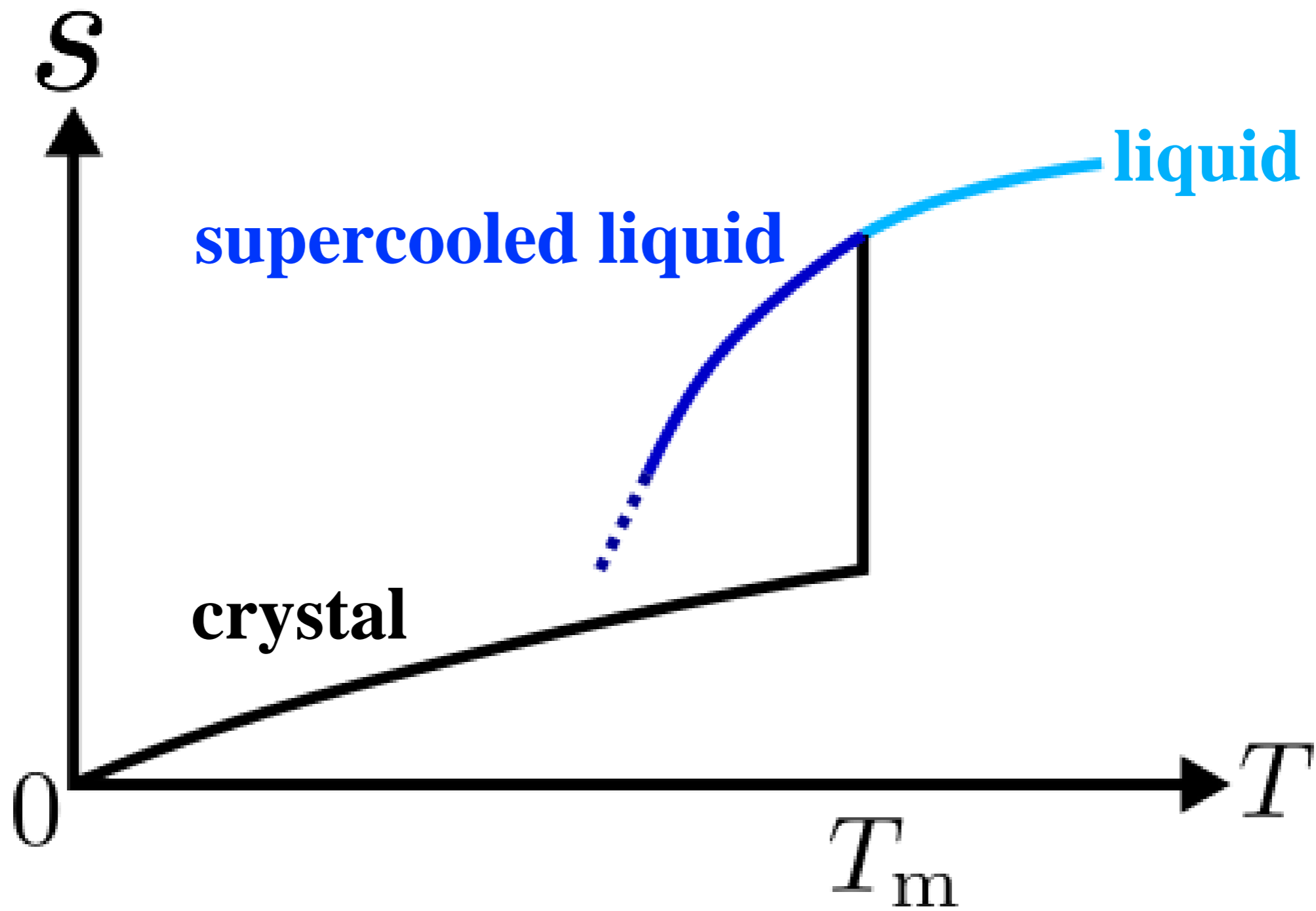
- 1 Review of glass phenomenology**
- 2 Amorphous order**
- 3 Beyond the glass ceiling**

1 Review of glass phenomenology

Glass phenomenology

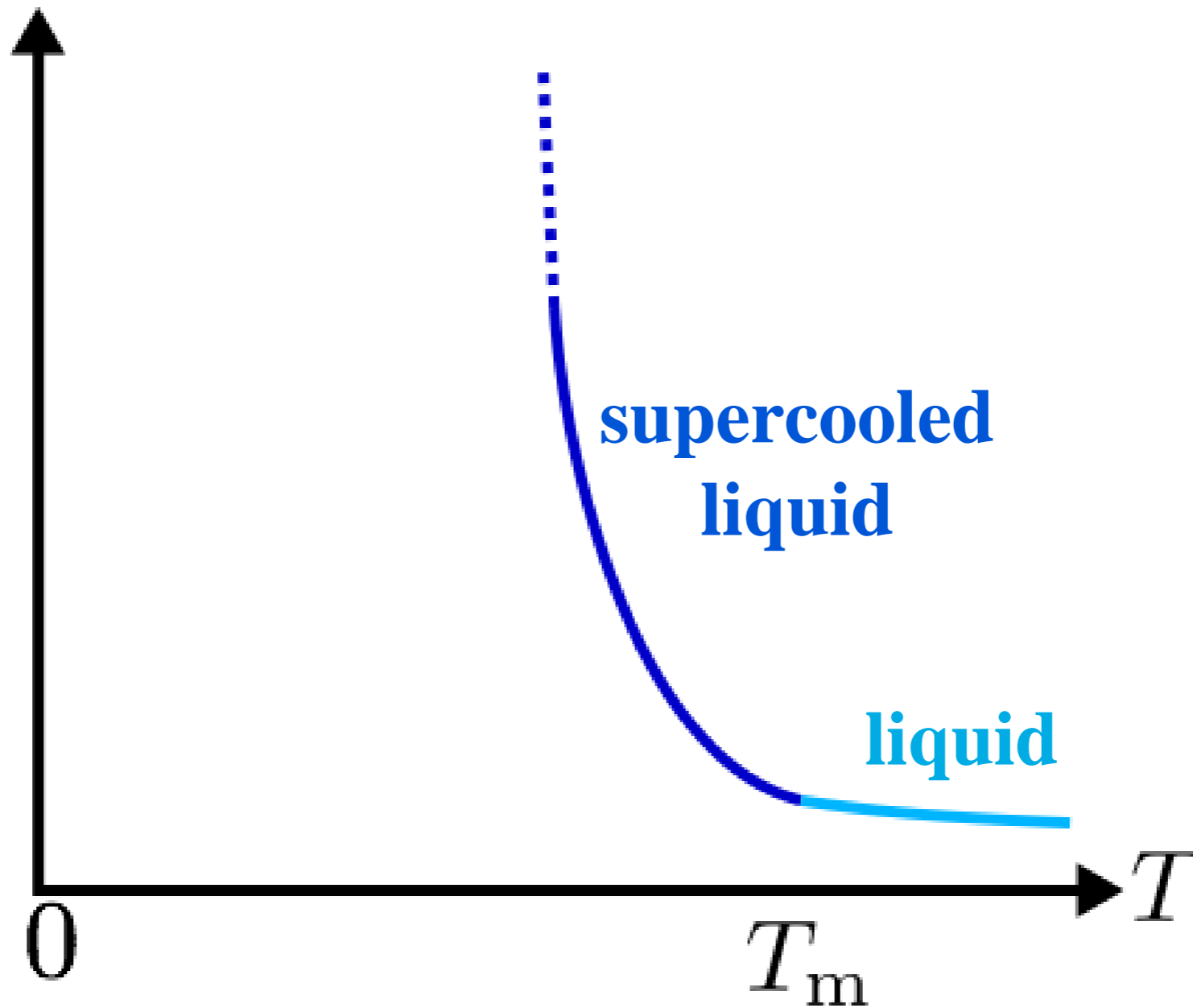


Glass phenomenology



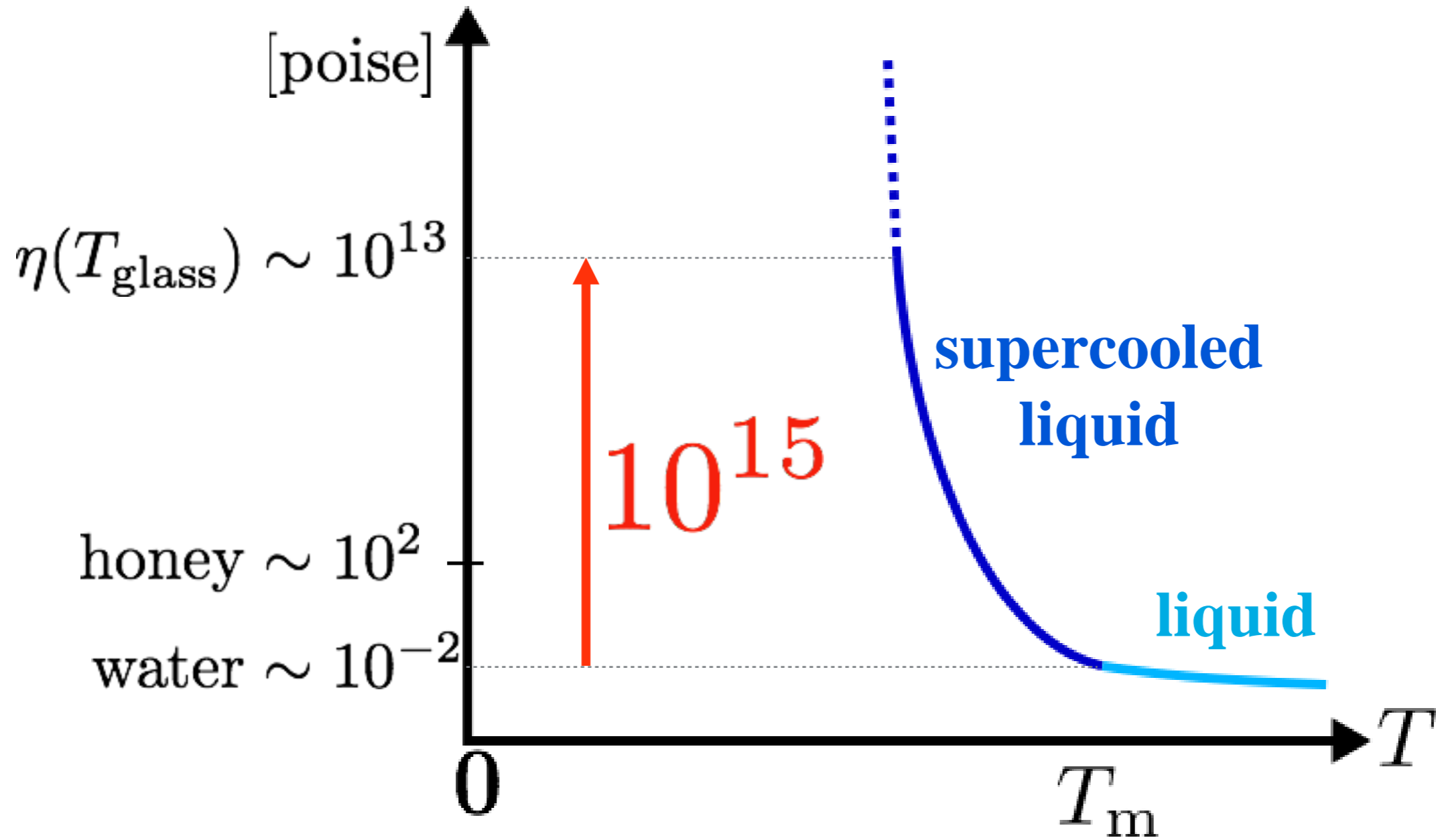
Glass phenomenology

shear viscosity $\eta \sim t_{\text{equilibration}}$



Glass phenomenology

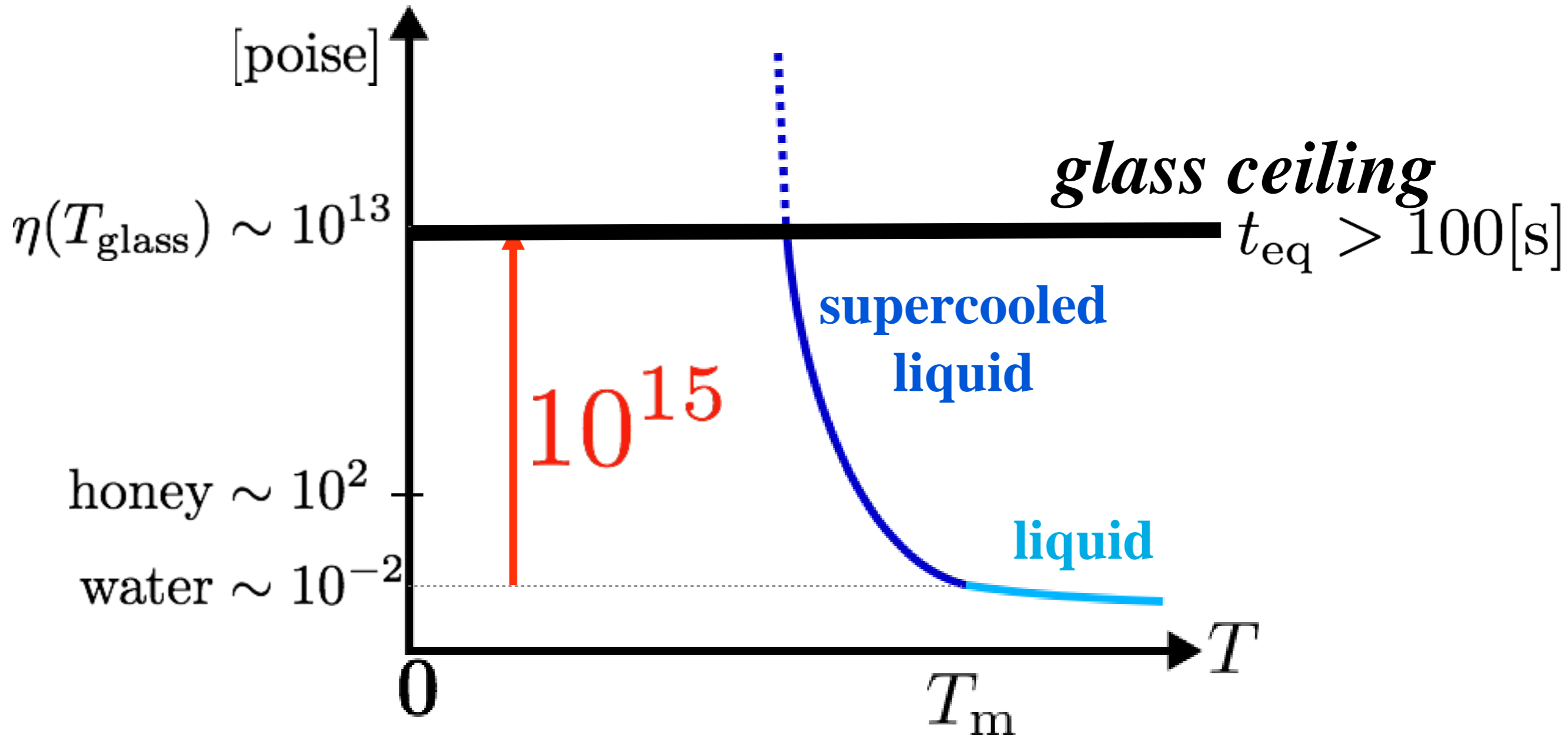
shear viscosity $\eta \sim t_{\text{equilibration}}$



typical molecular liquids

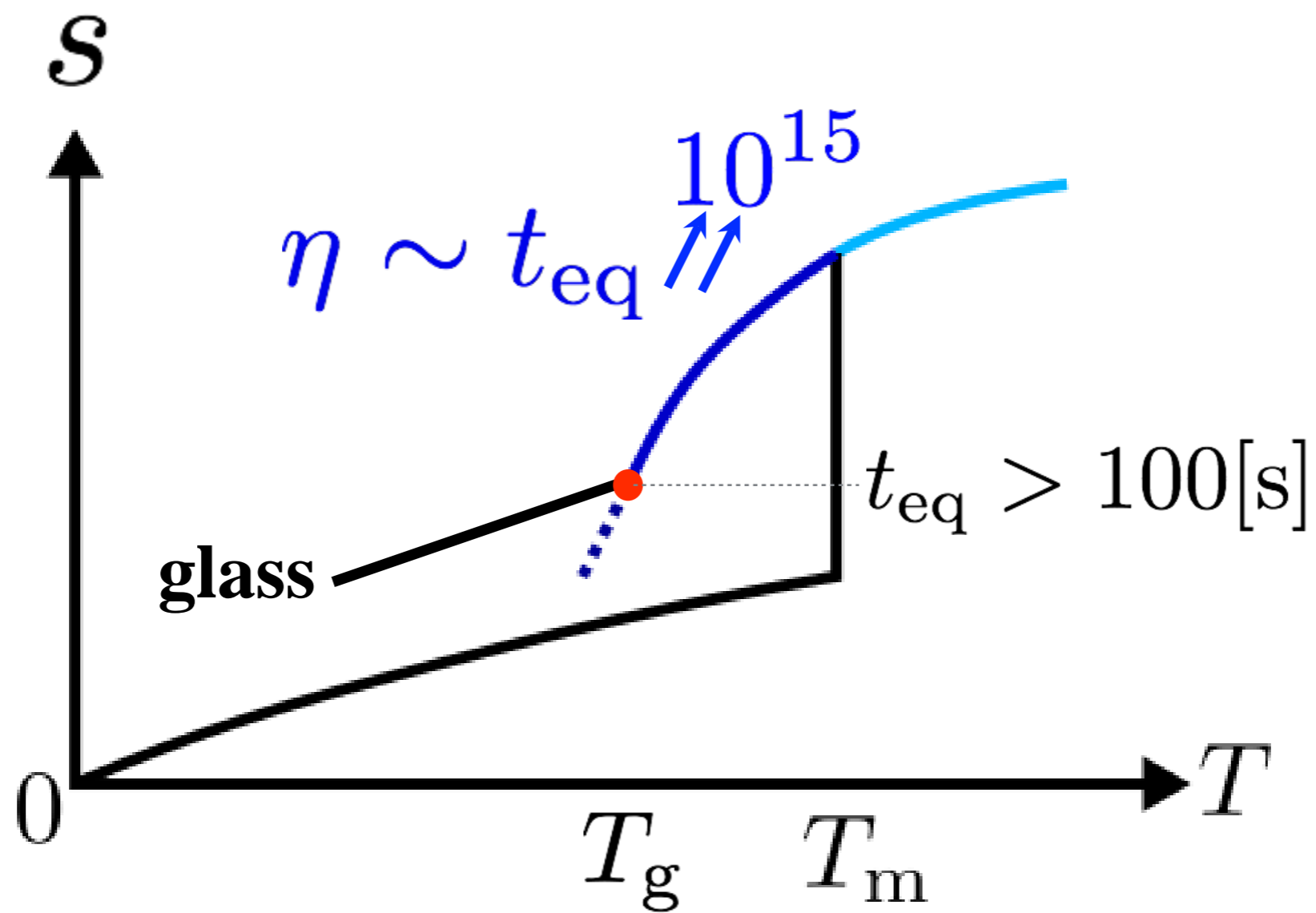
Glass phenomenology

shear viscosity $\eta \sim t_{\text{equilibration}}$

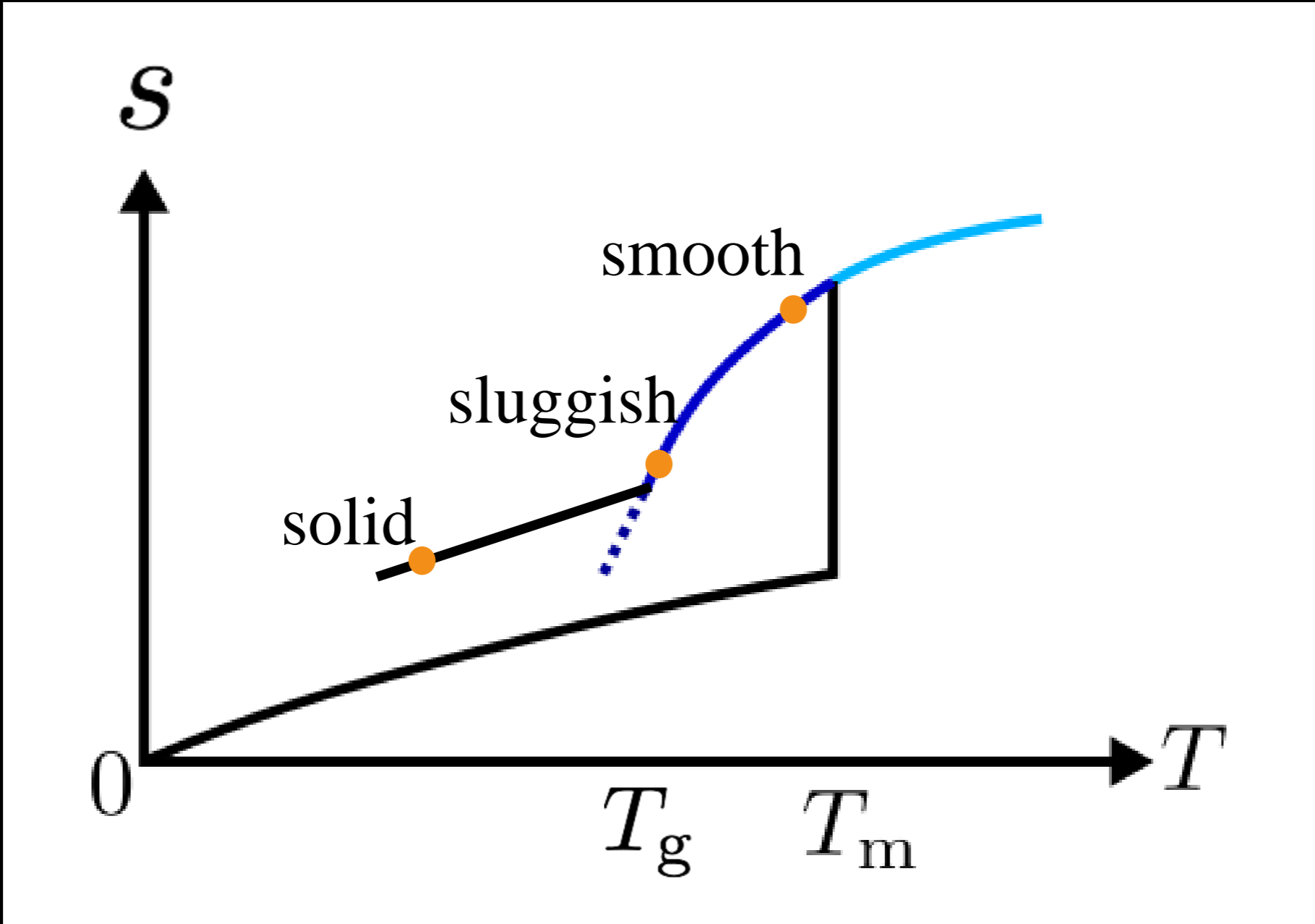


typical molecular liquids

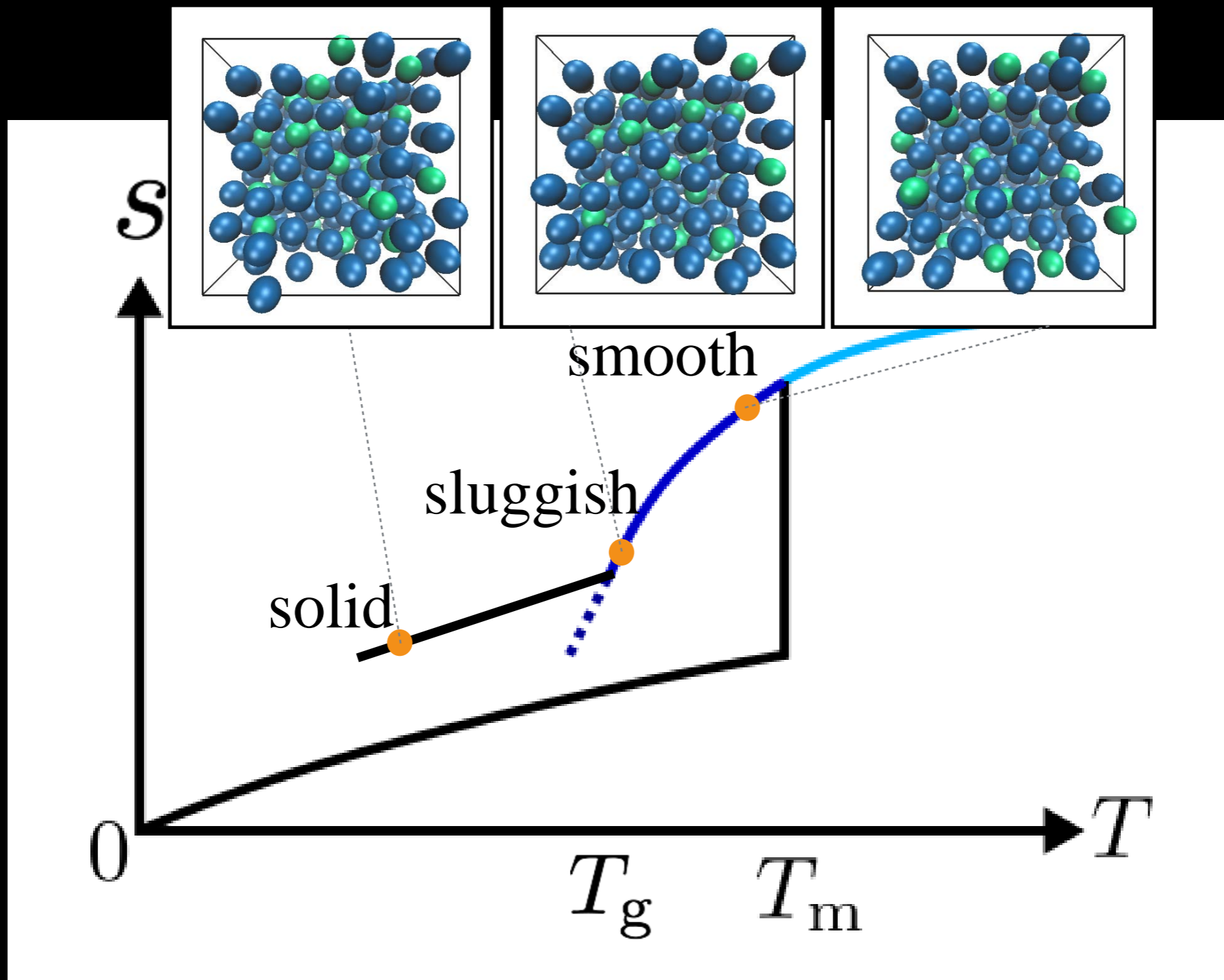
Glass phenomenology



Glass mystery?

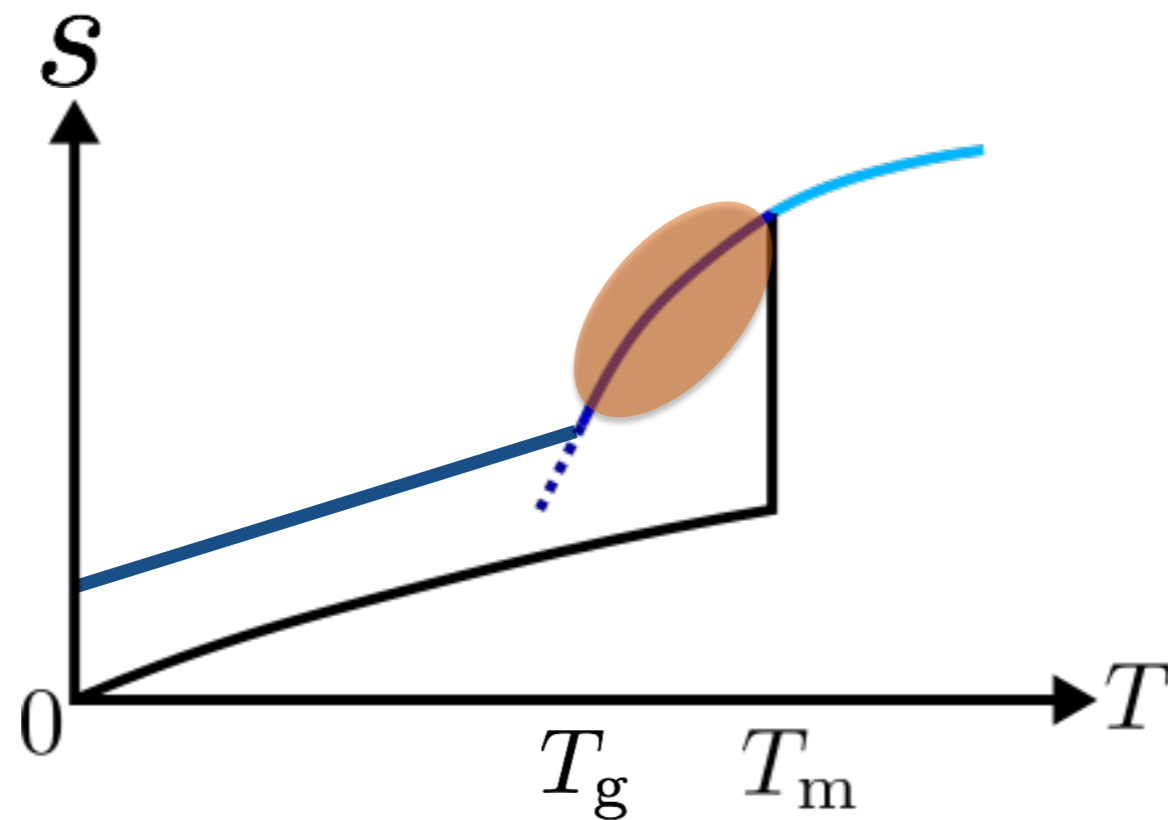


Glass mystery?



They all look the same, amorphous mess!

2 Amorphous order



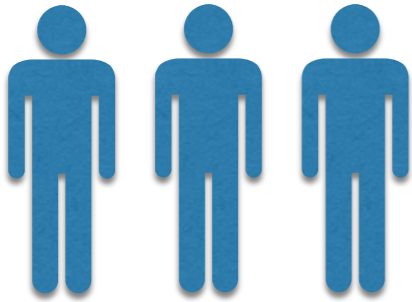
Glass slowdown

Dramatic slowdown without structural change:

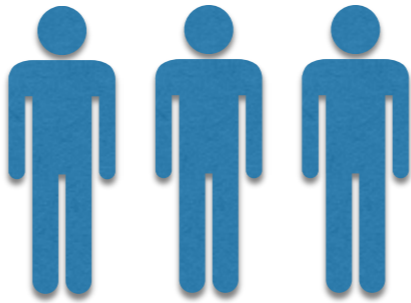
$\eta \nearrow 10^{15}$ while $\langle \rho(\mathbf{r})\rho(\mathbf{0}) \rangle_T$
barely changes

Competing theories for slowdown

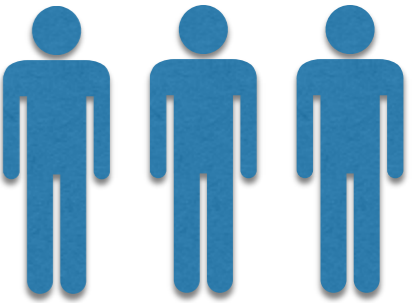
Mode-coupling



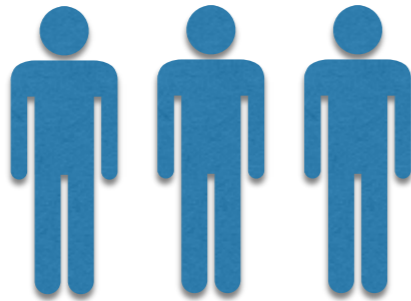
Locally preferred structures



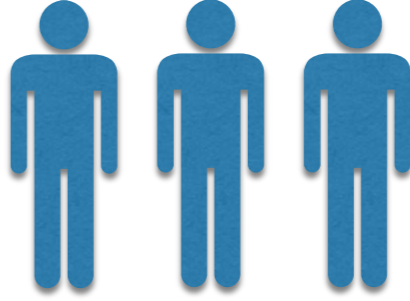
...



Dynamic facilitation

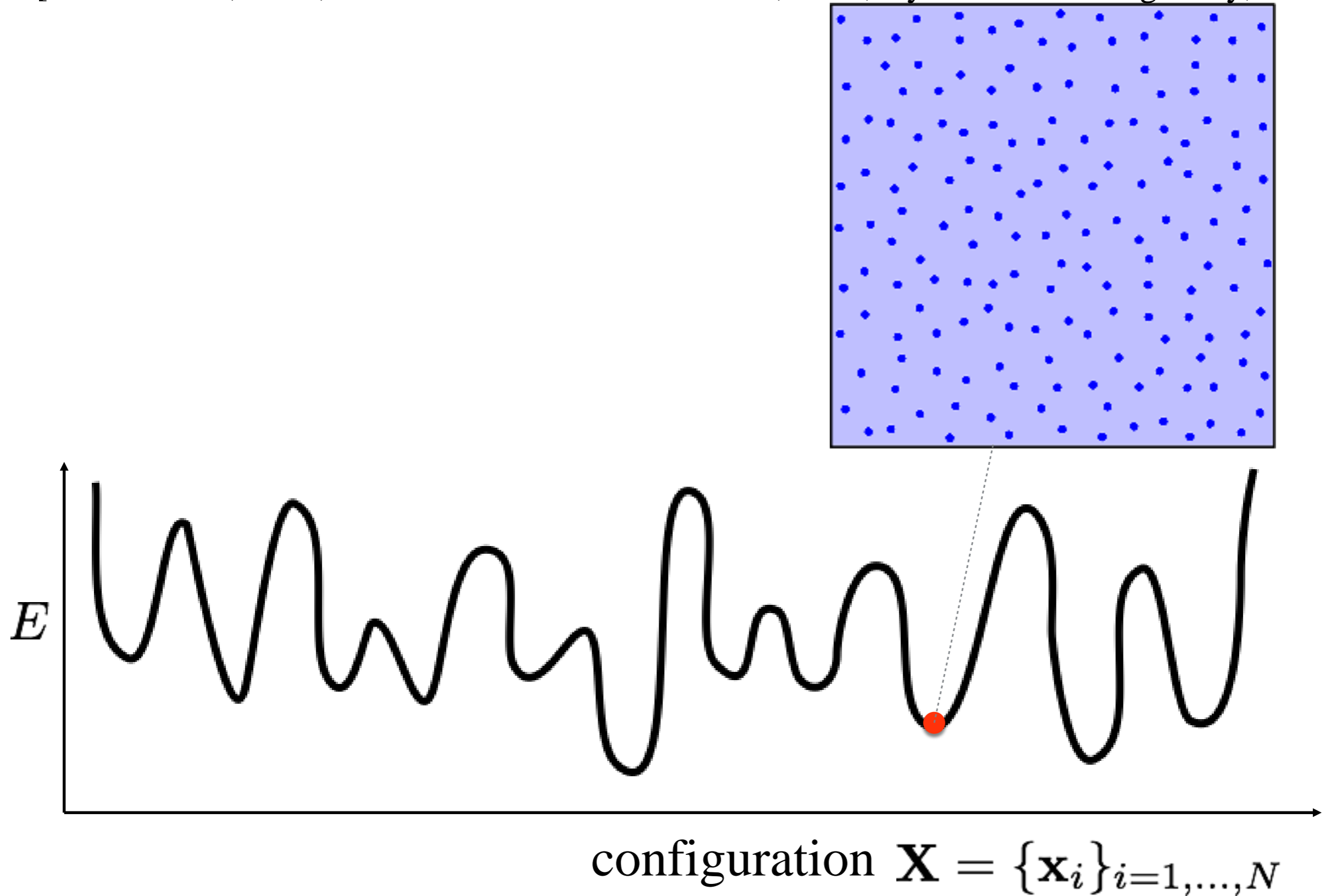


Energy landscape



Energy landscape

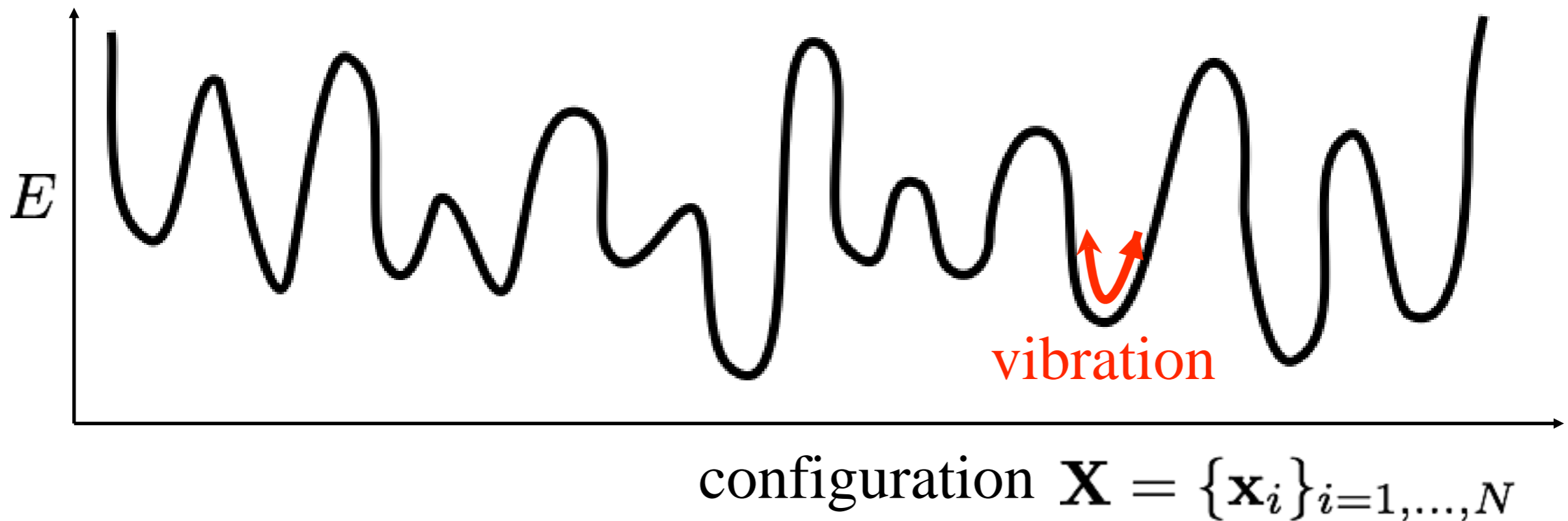
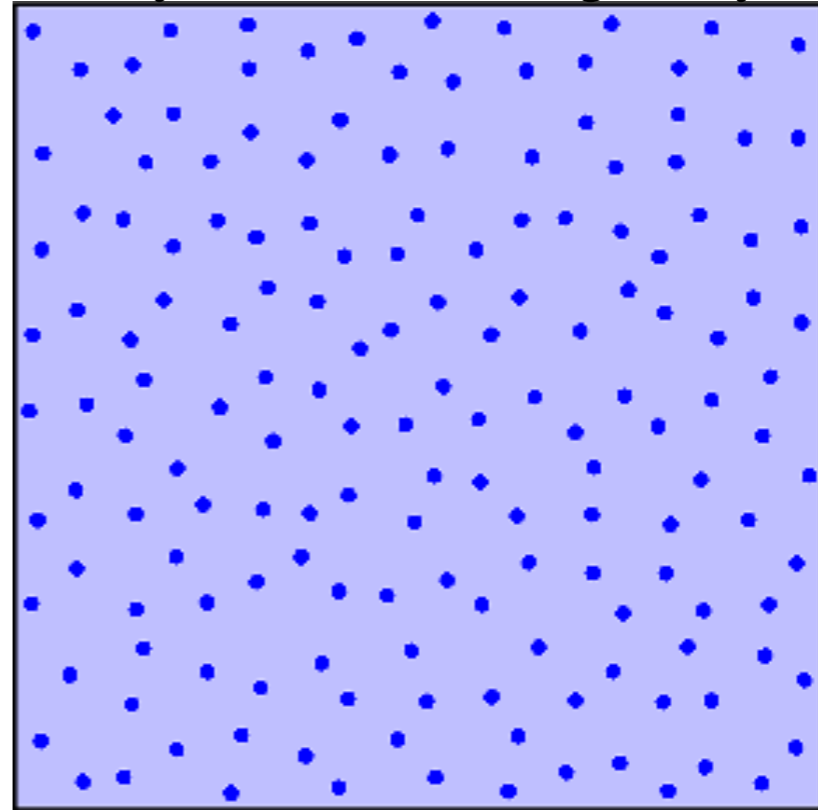
[Adam-Gibbs, 1965; Random First-Order Transition, 1989; Dynamical Heterogeneity, ~2000]



Energy landscape

[Adam-Gibbs, 1965; Random First-Order Transition, 1989; Dynamical Heterogeneity, ~2000]

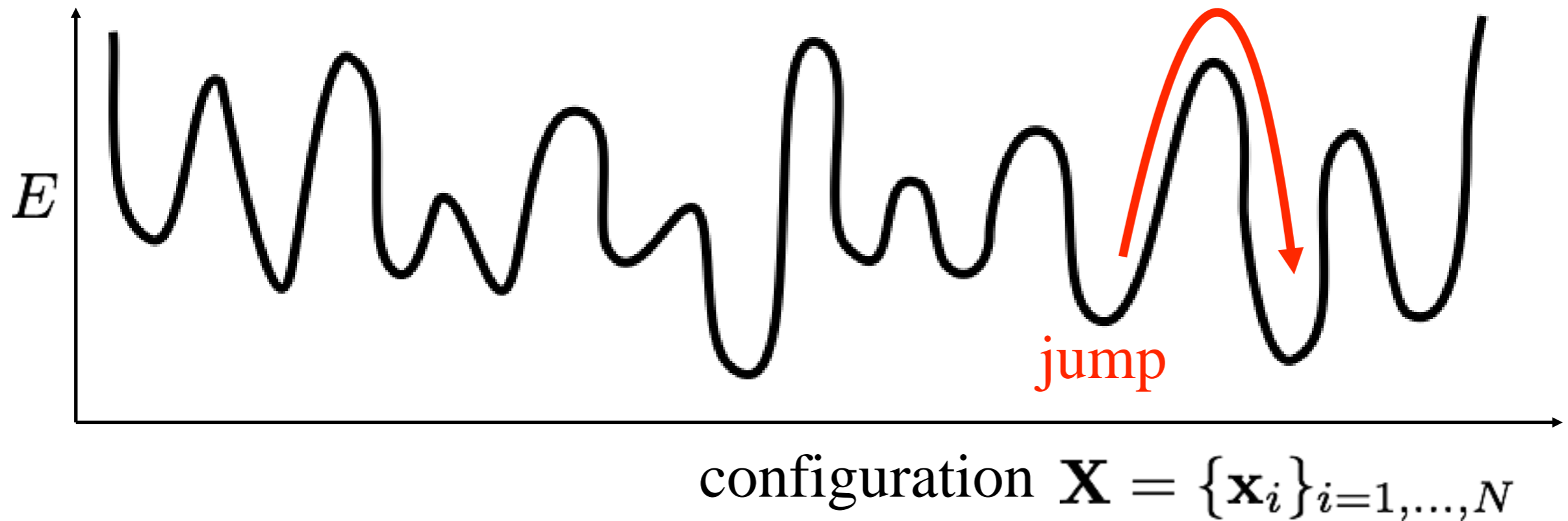
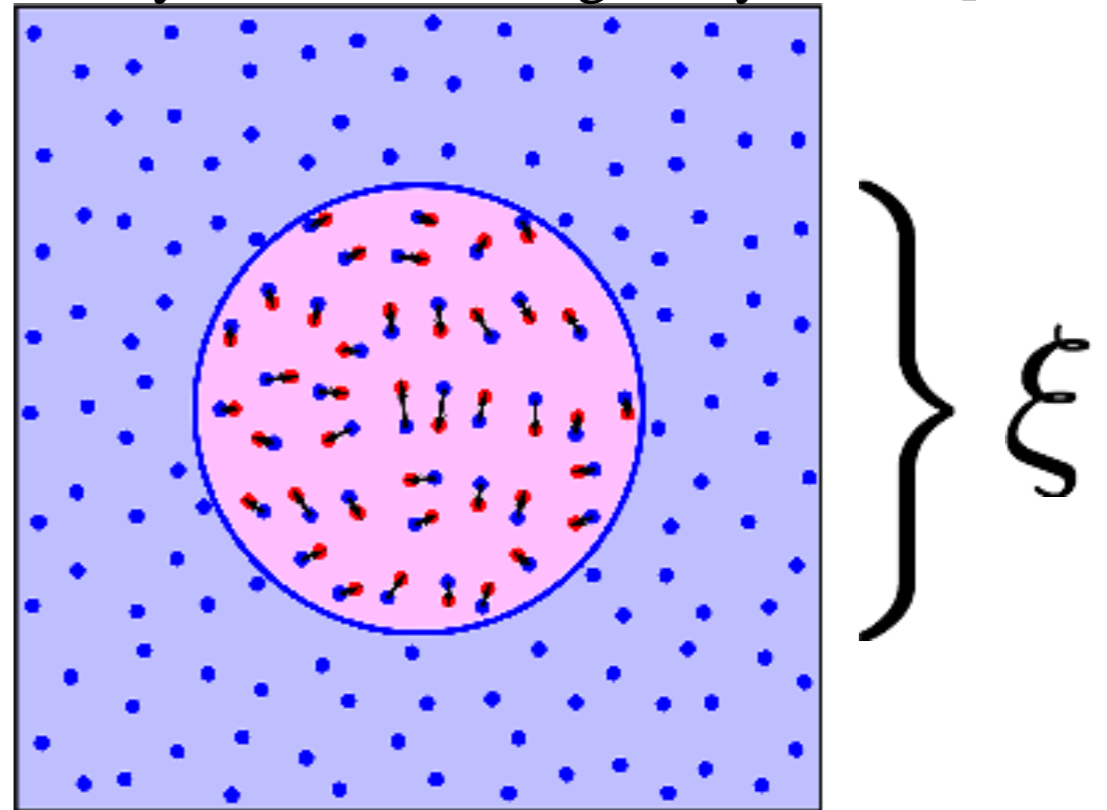
$$t \ll t_{\text{eq}}$$



Energy landscape

[Adam-Gibbs, 1965; Random First-Order Transition, 1989; Dynamical Heterogeneity, ~2000]

$$t \sim t_{\text{eq}}$$

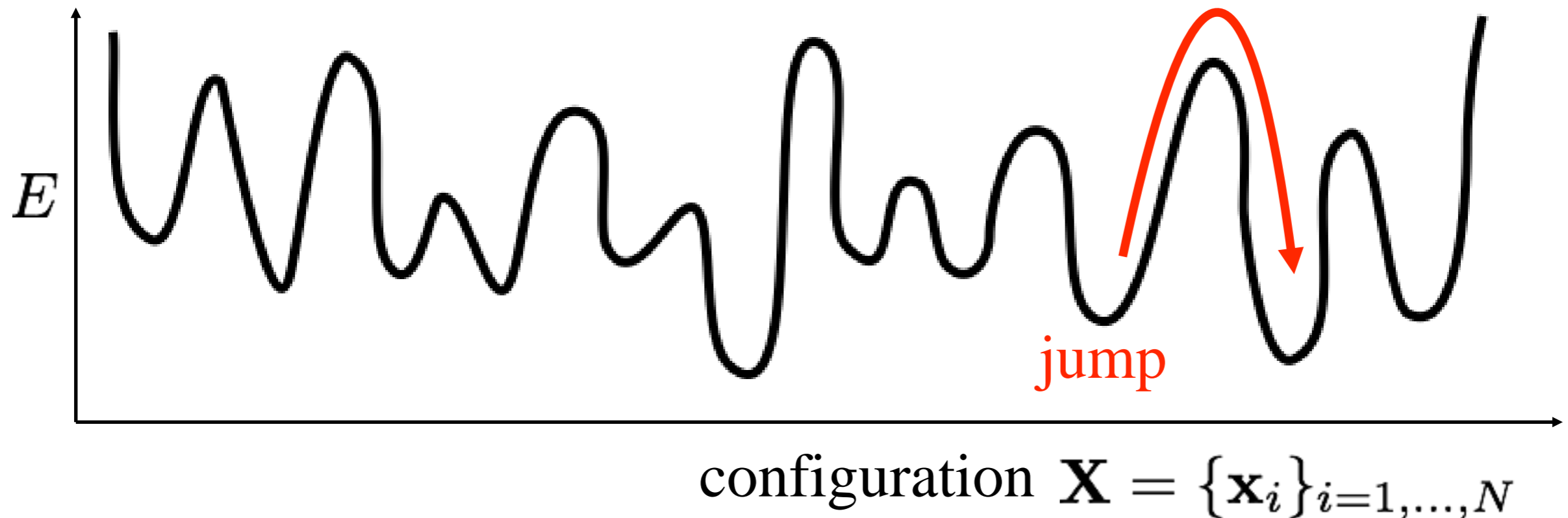
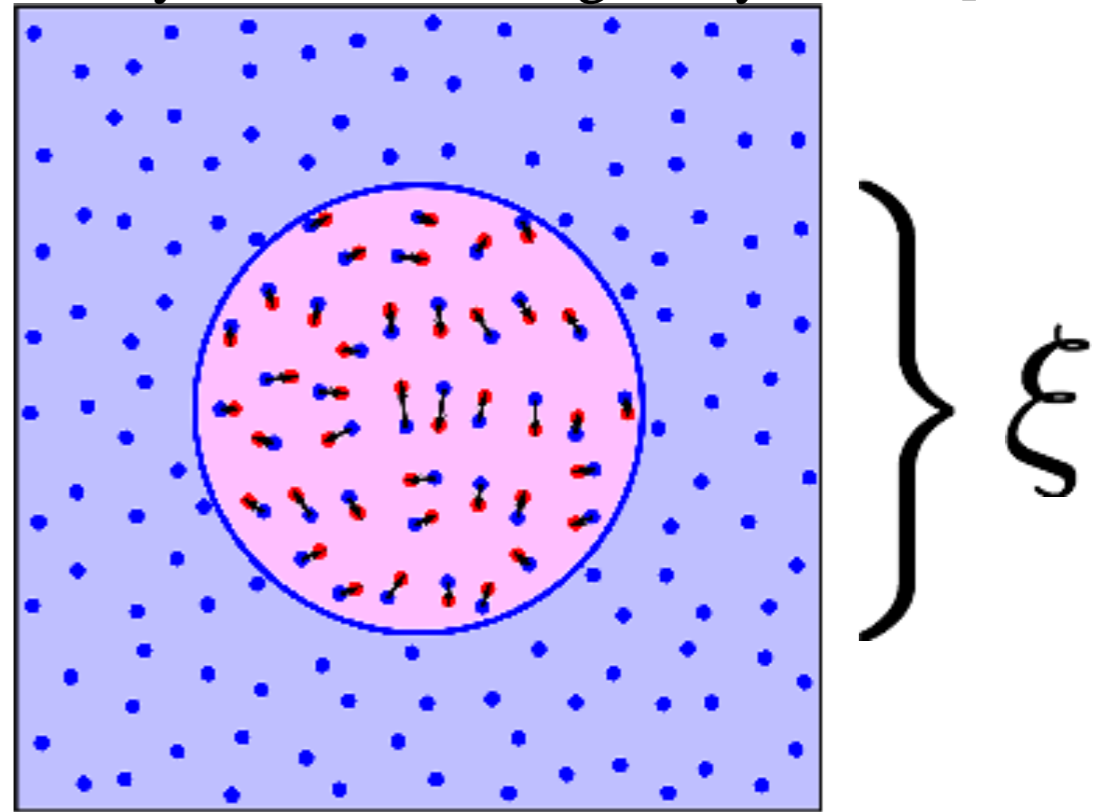


Energy landscape

[Adam-Gibbs, 1965; Random First-Order Transition, 1989; Dynamical Heterogeneity, ~2000]

$$t \sim t_{\text{eq}} \\ \sim e^{\frac{E_{\text{act}}}{T}}$$

$$\text{with } E_{\text{act}} \propto \xi^p$$



Energy landscape

[Adam-Gibbs, 1965; Random First-Order Transition, 1989; Dynamical Heterogeneity, ~2000]

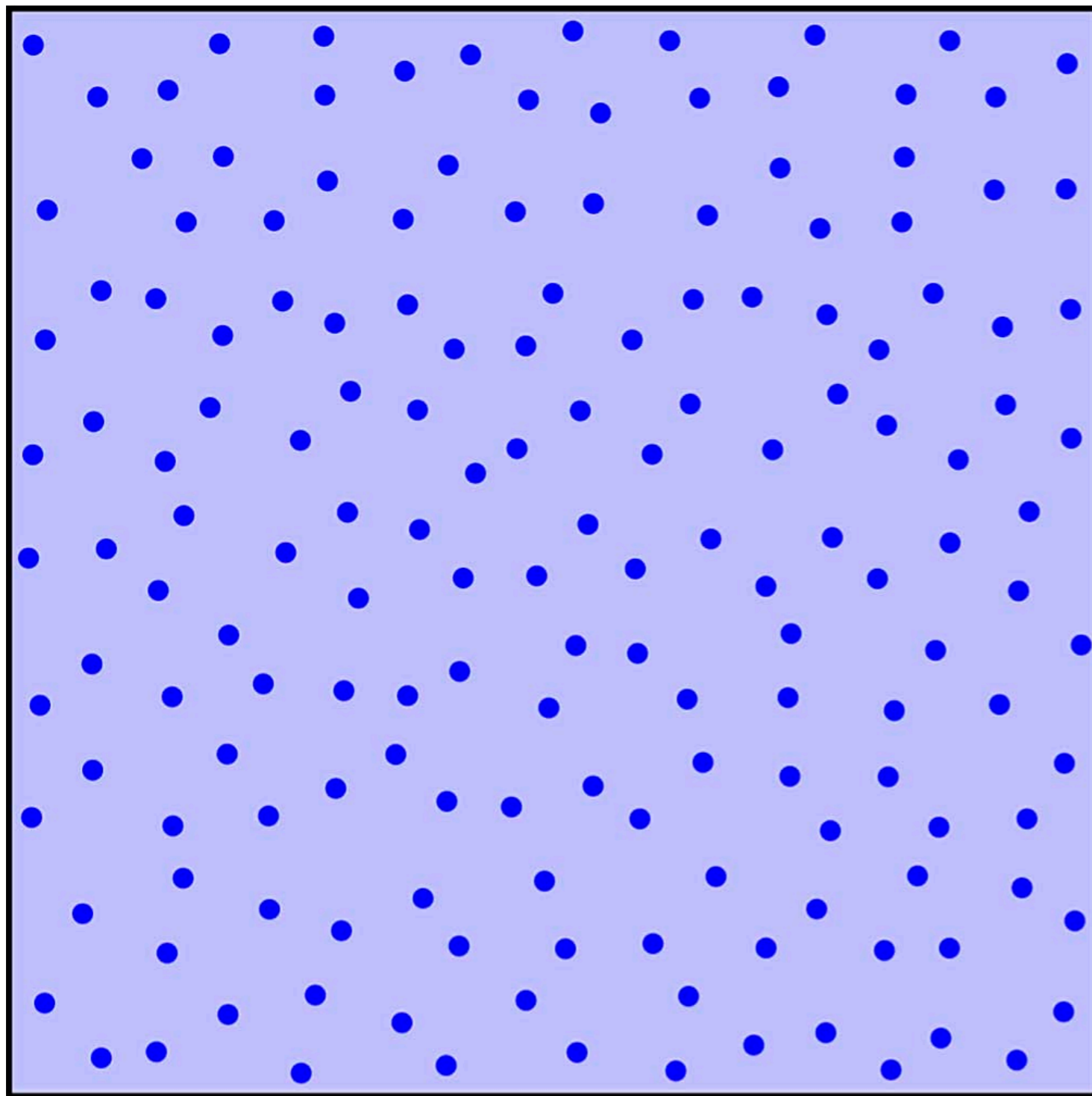
How do we make all this precise?

Point-to-set correlations!!!

[Biroli-Bouchaud, 2004]

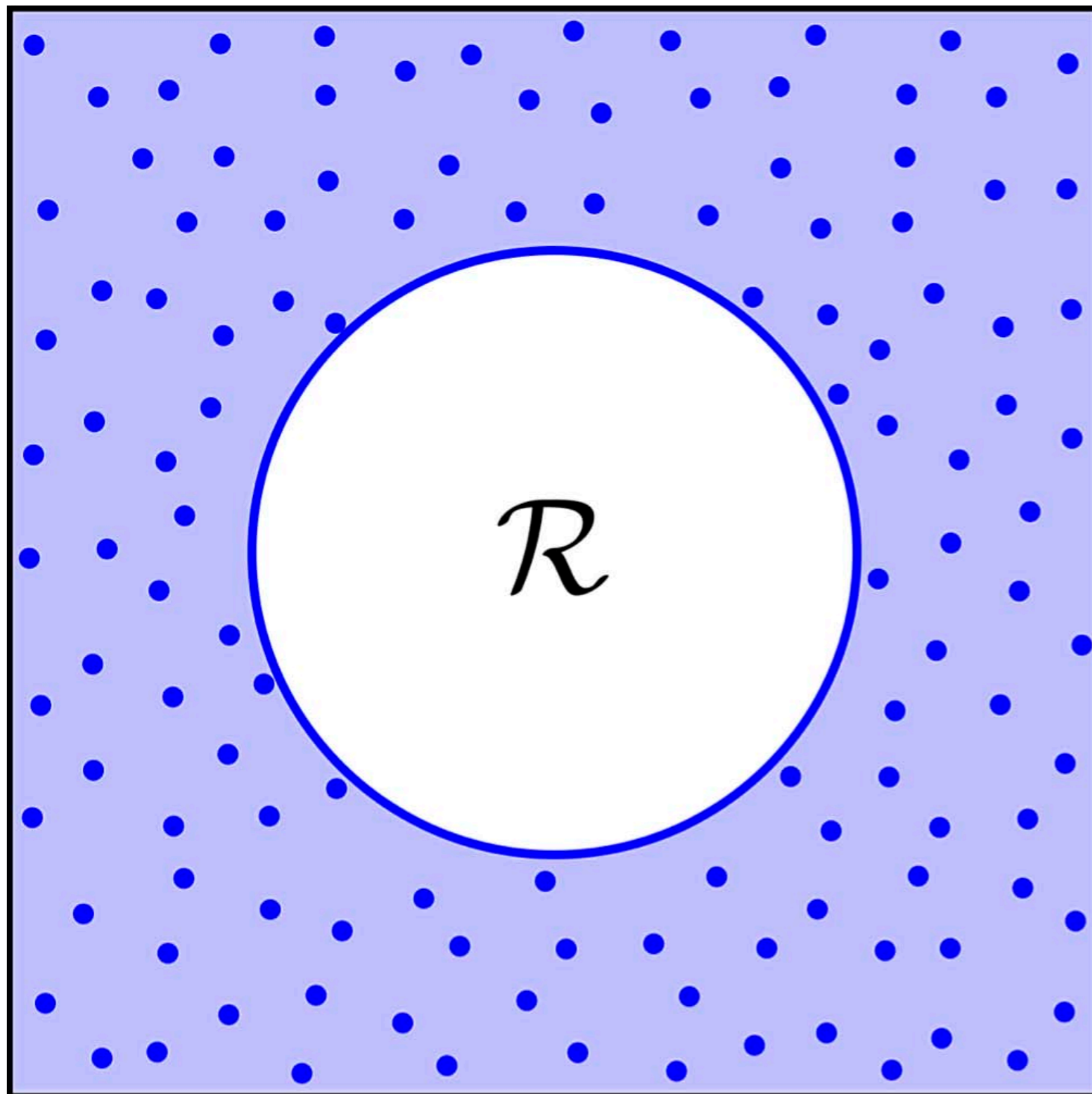
[Montanari-Semerjian, 2006]

Point-to-set correlations



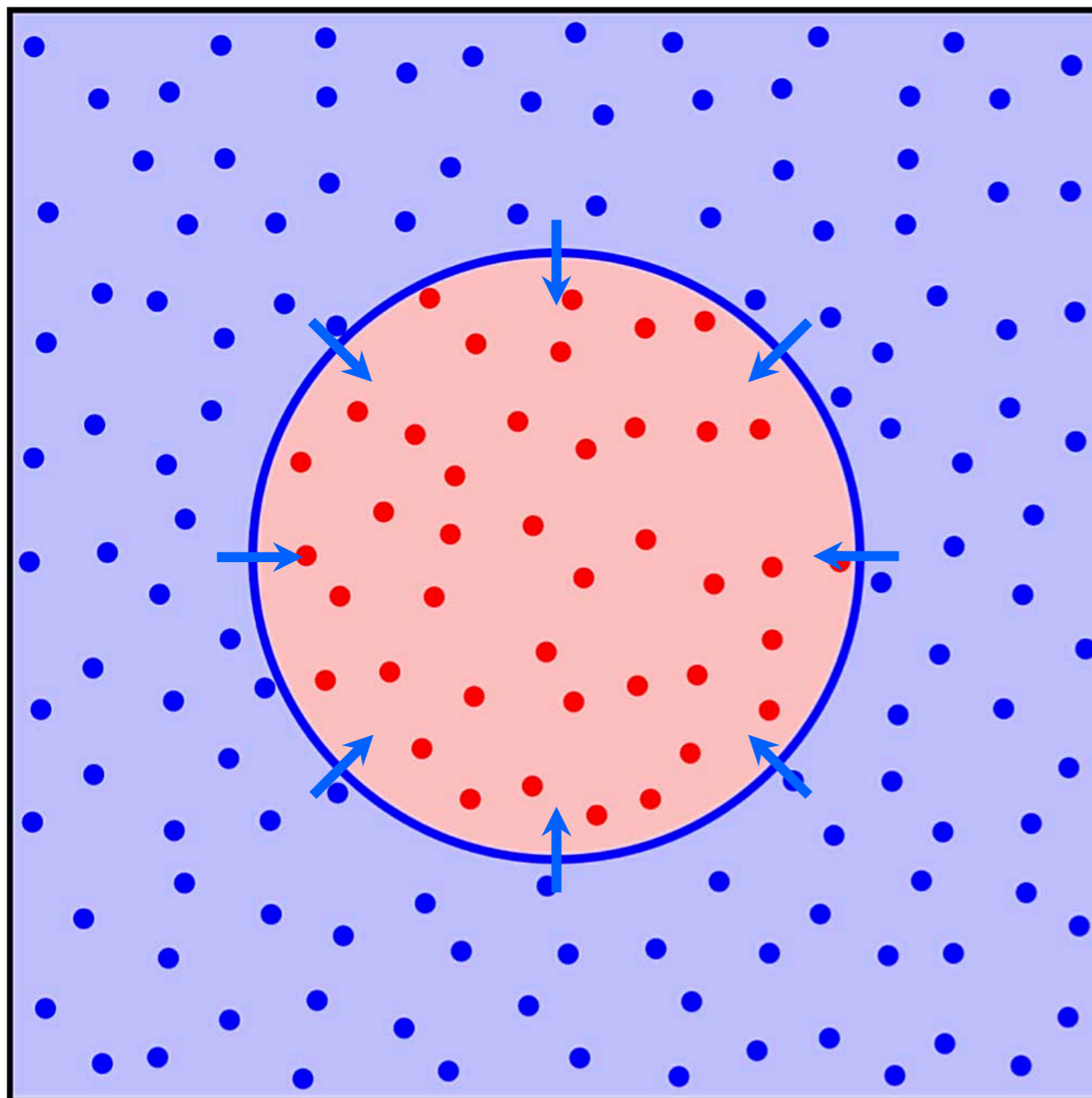
$$\mathbf{X}_1 = \{\mathbf{x}_i\}_{i=1, \dots, N}$$

Point-to-set correlations



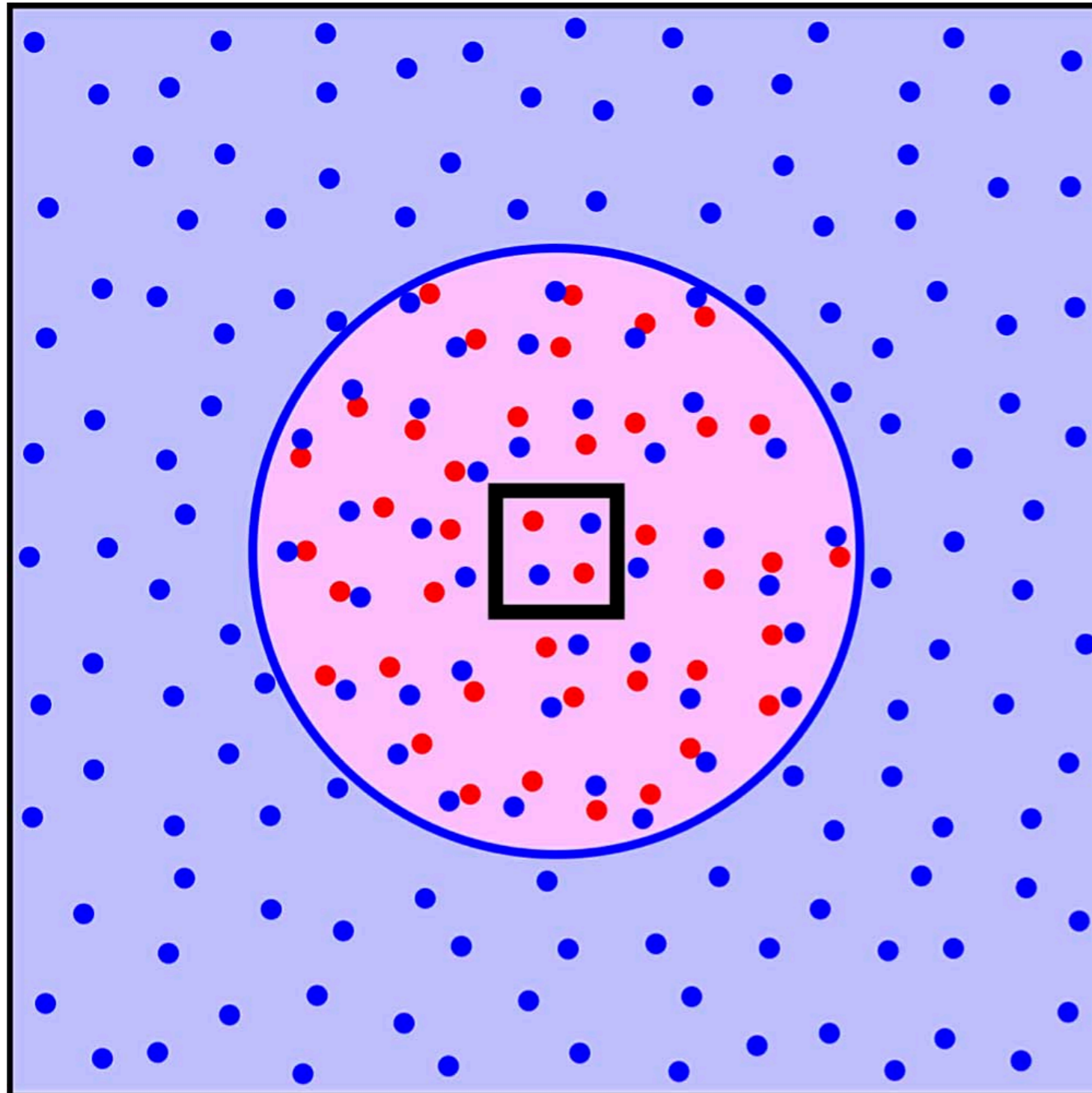
X_1^{out}

Point-to-set correlations



$$\mathbf{X}_2 = \mathbf{X}_1^{\text{out}} + \mathbf{X}_2^{\text{in}}$$

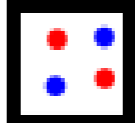
Point-to-set correlations



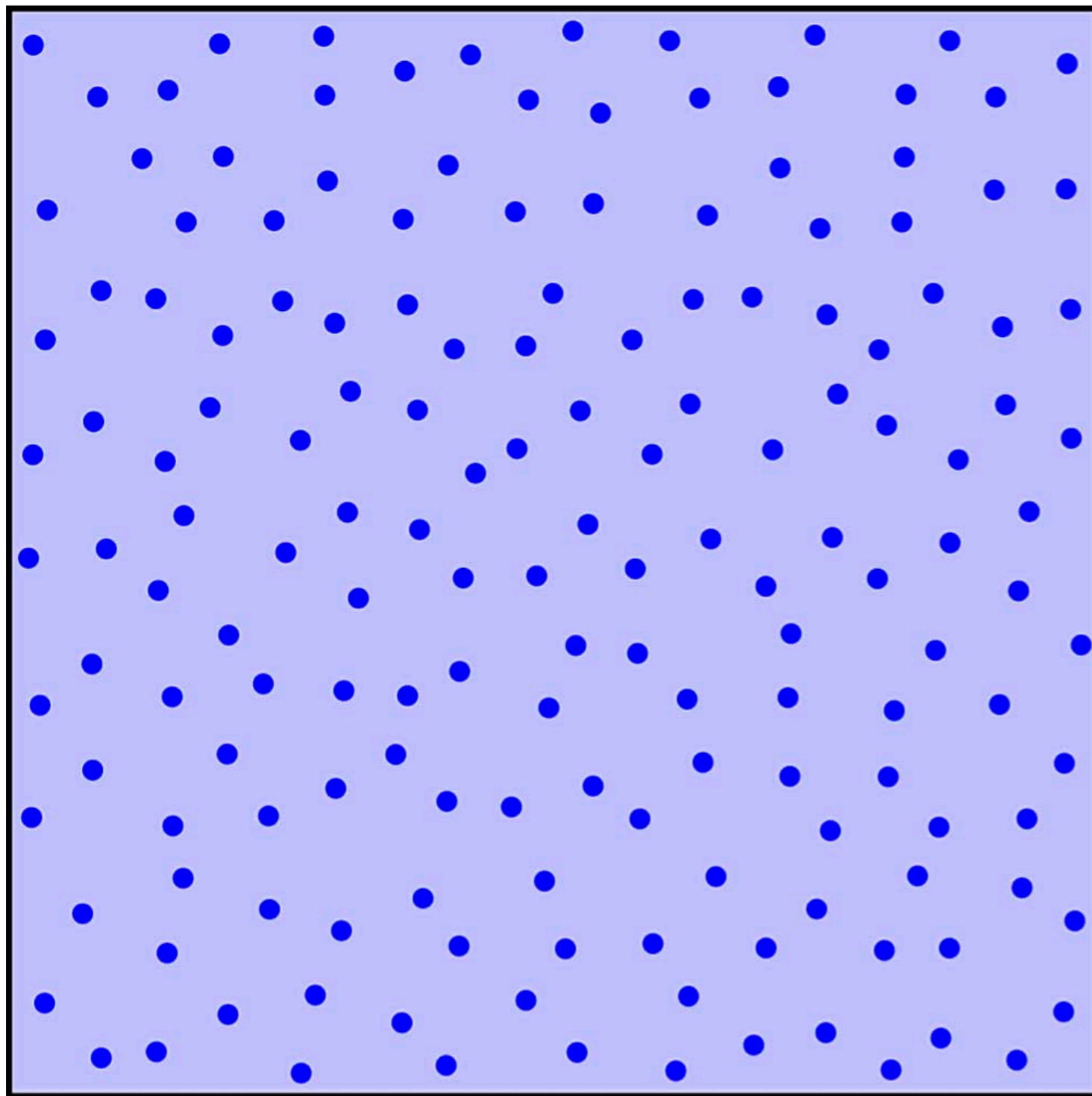
How similar
 X_1 and X_2 ?

Overlap
 \hat{q}_{12}

 similar $q \sim 1$

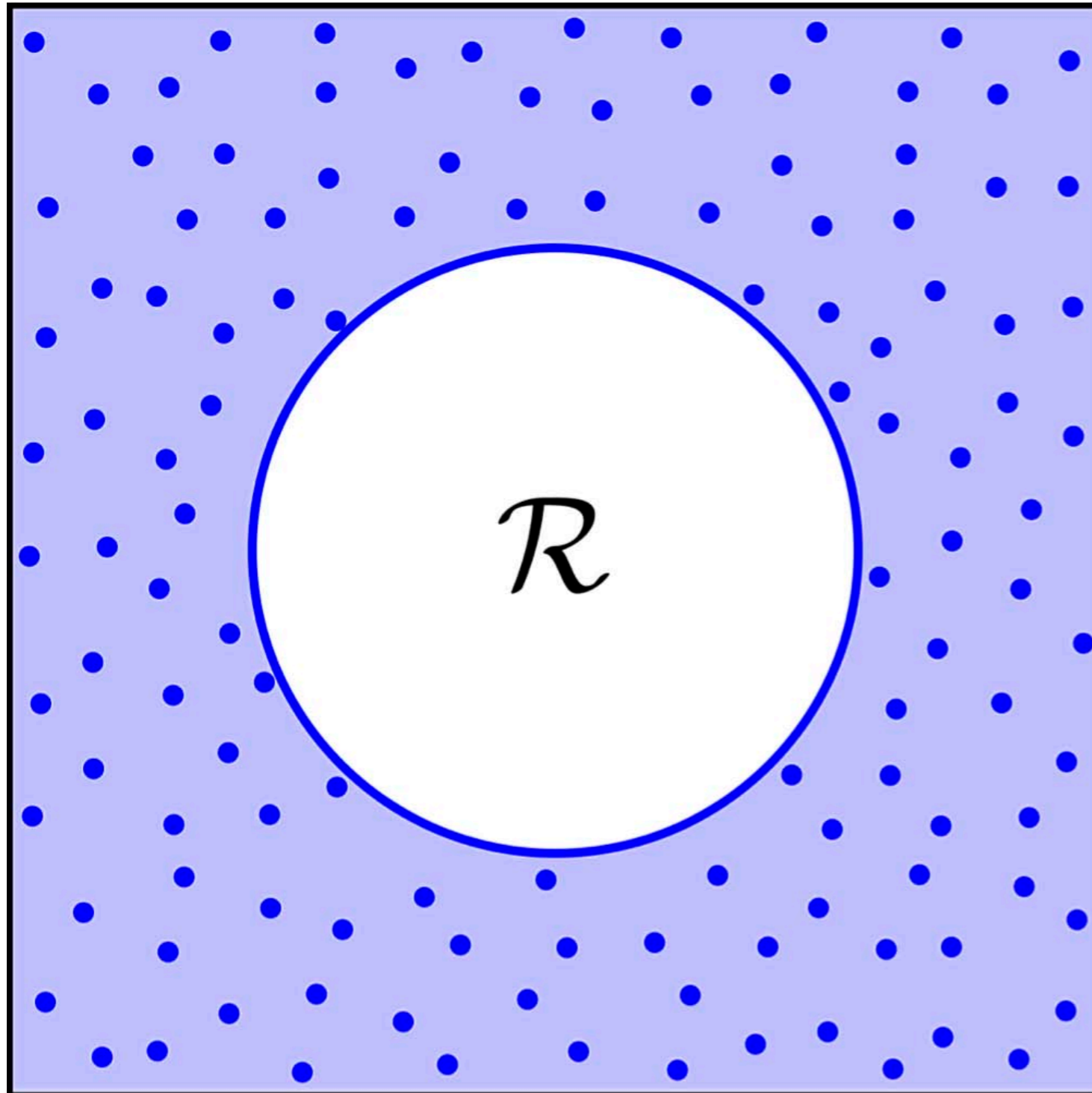
 dissimilar $q \sim 0$

Point-to-set correlations



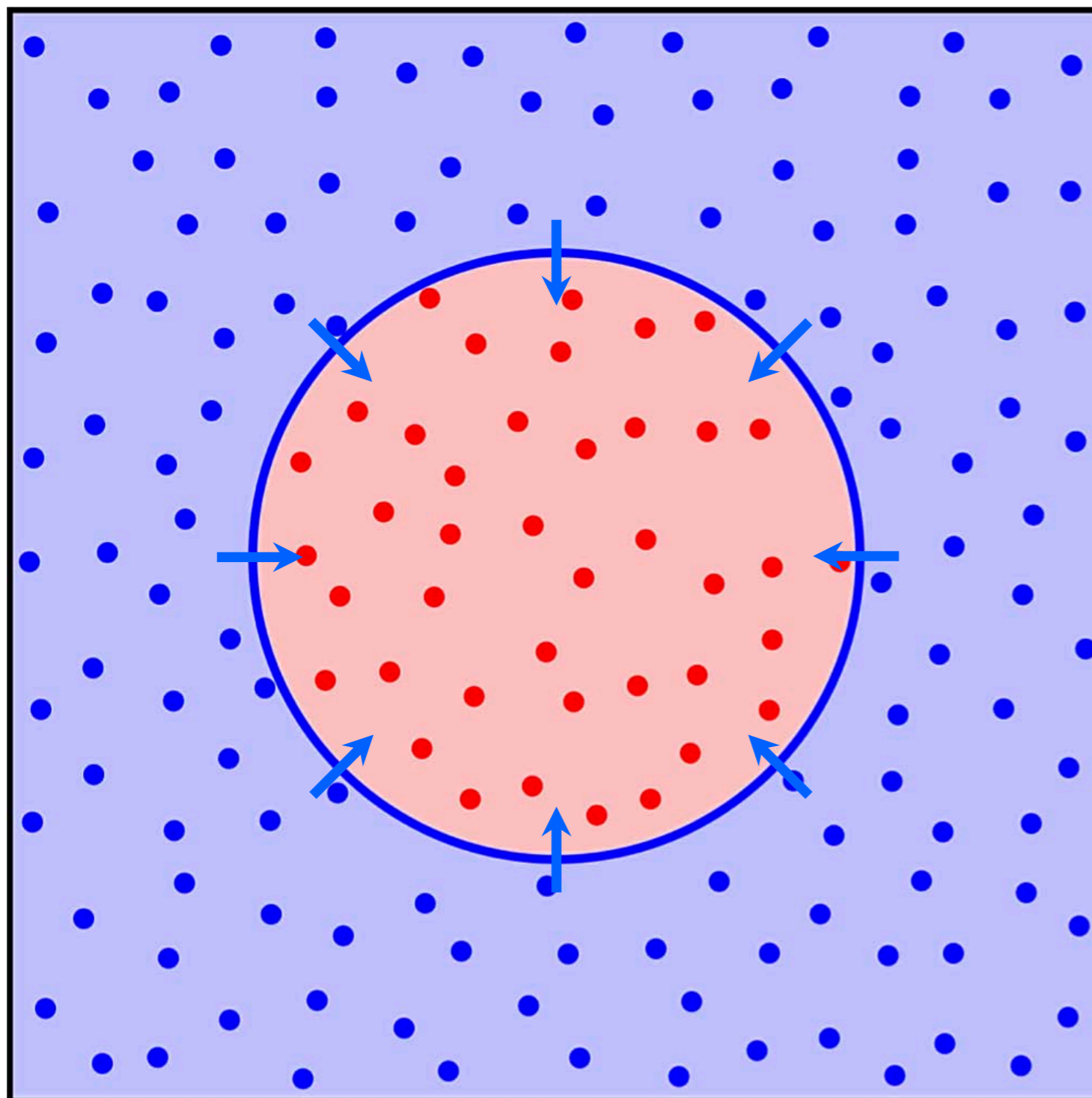
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Point-to-set correlations



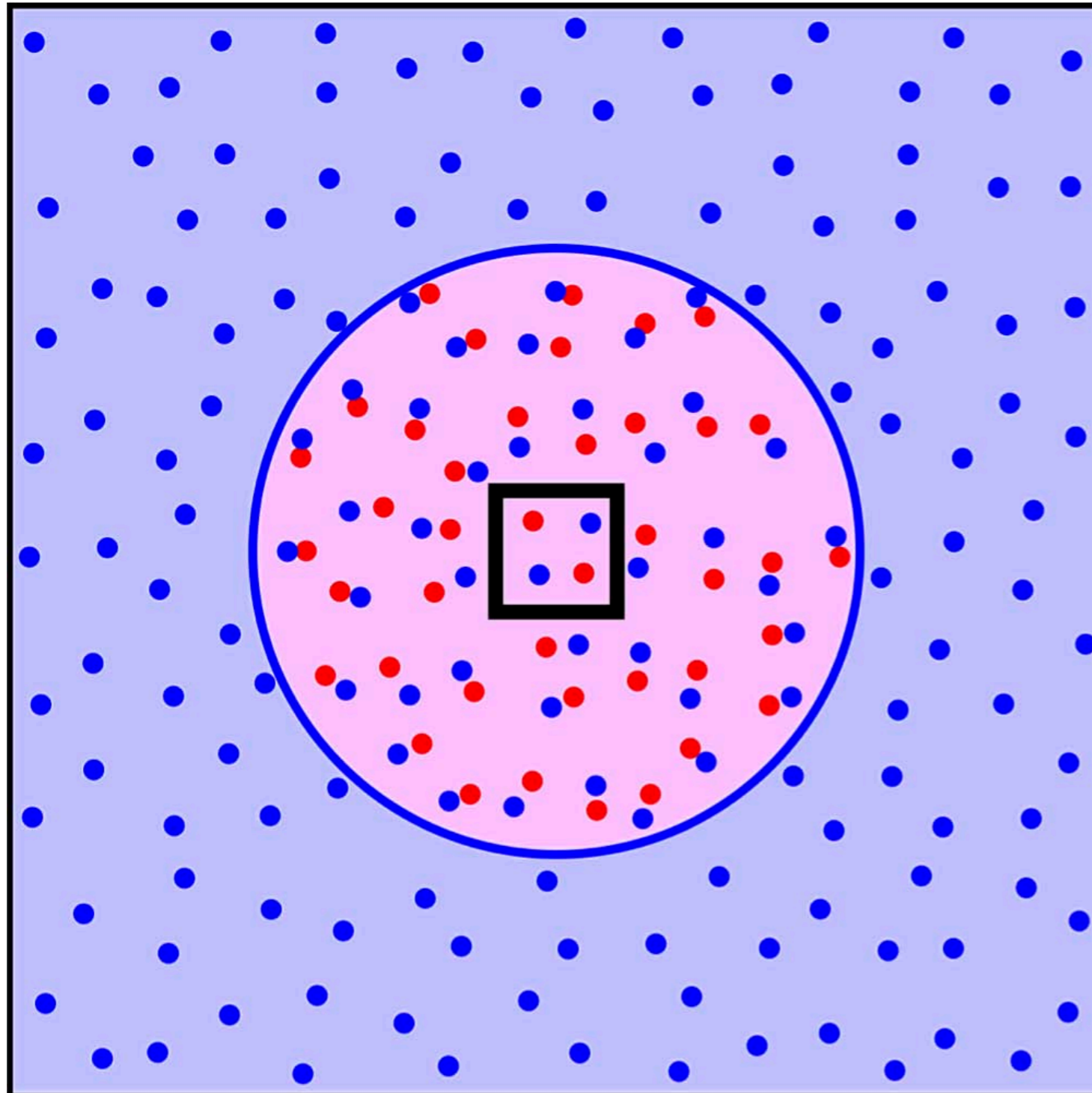
$\mathbf{X}_1^{\text{out}}$

Point-to-set correlations



$$\mathbf{X}_2 = \mathbf{X}_1^{\text{out}} + \mathbf{X}_2^{\text{in}}$$

Point-to-set correlations



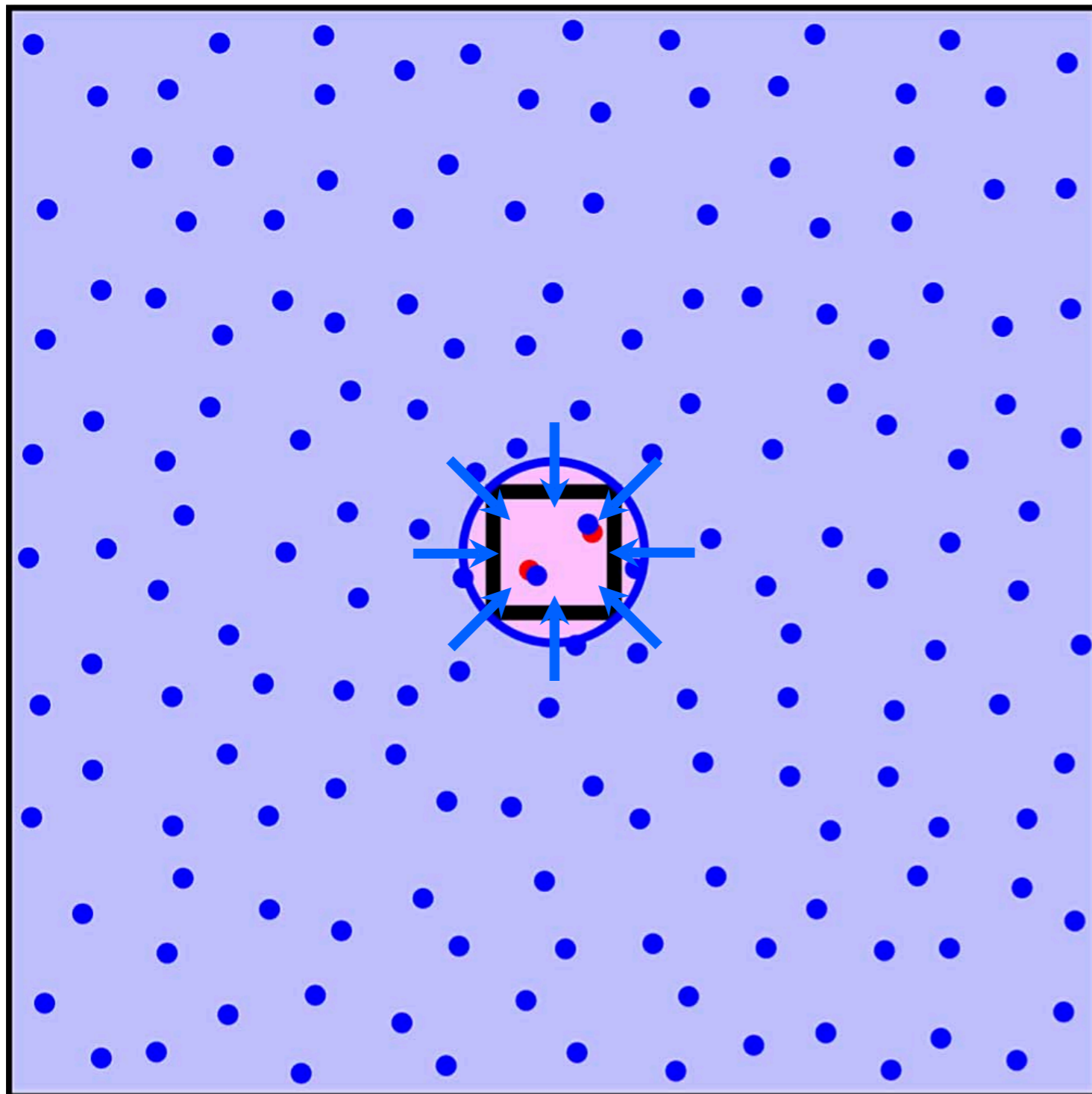
How similar
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 \hat{q}_{12}

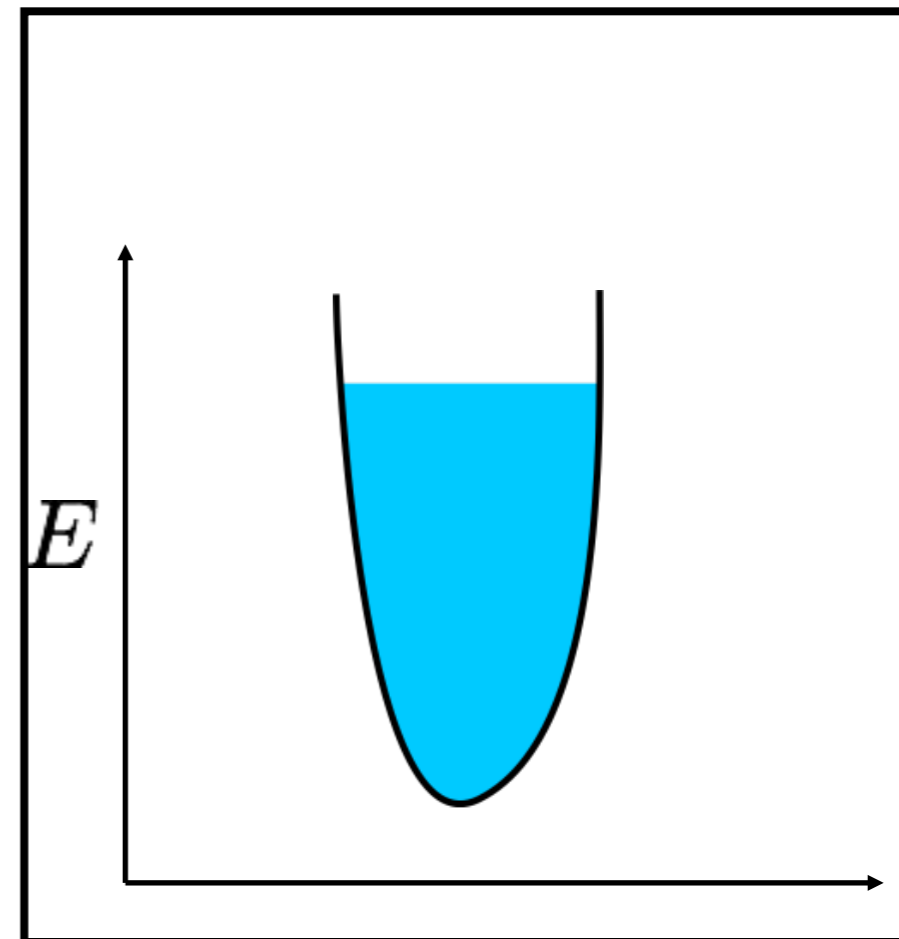
 similar $q \sim 1$

 dissimilar $q \sim 0$

Local free-energy landscape



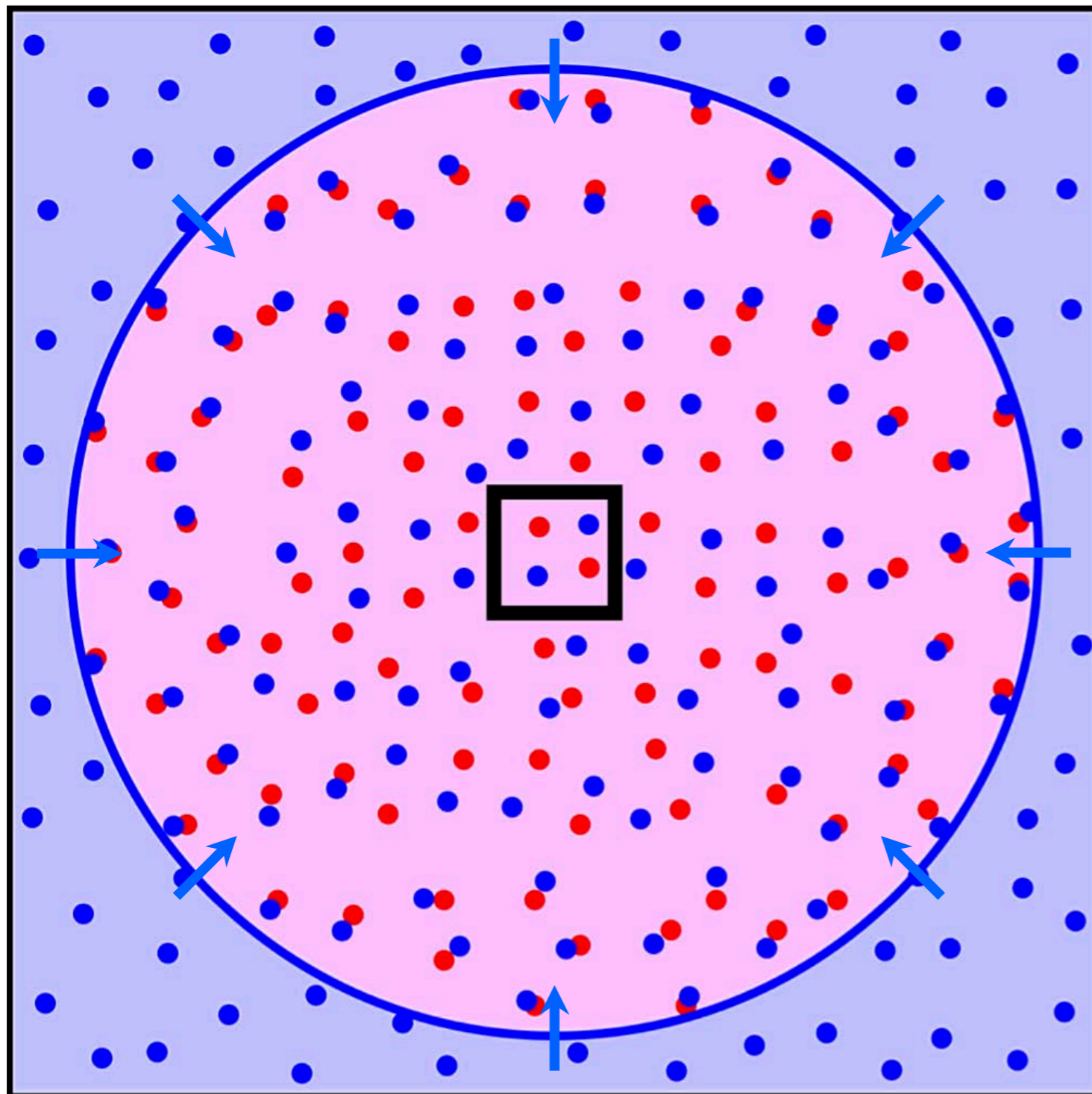
\mathcal{R} small



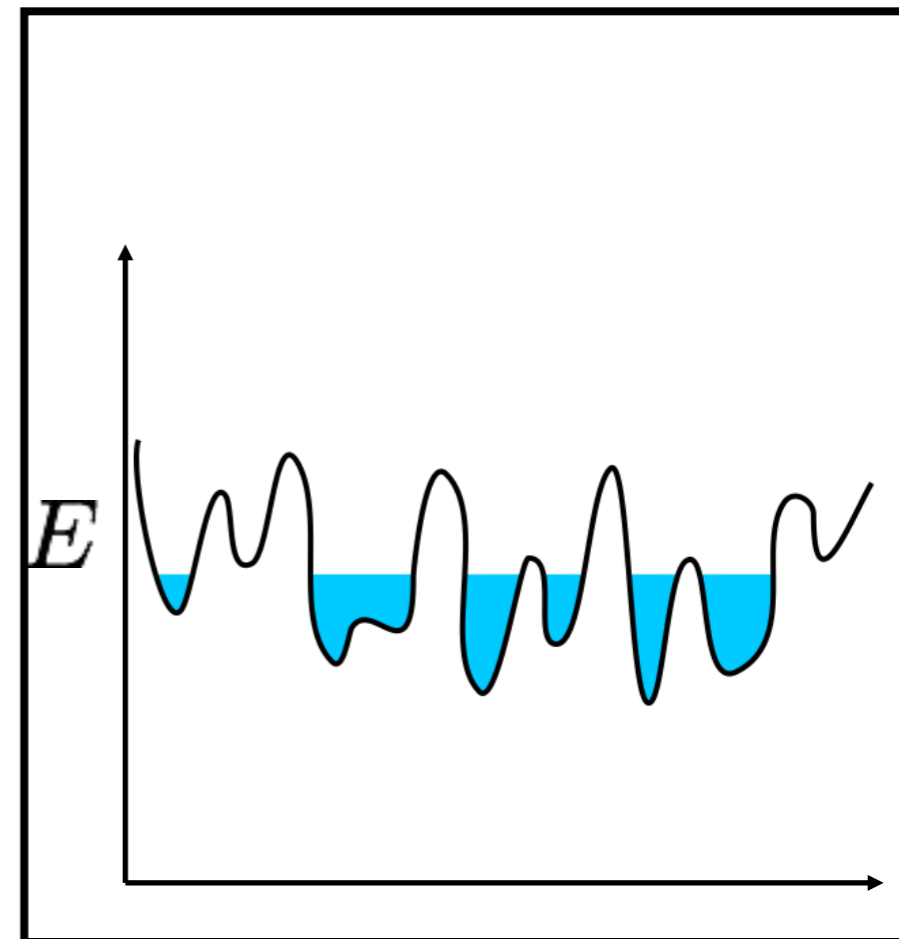
Overlap: High

$$\hat{q}_{12} \sim 1$$

Local free-energy landscape



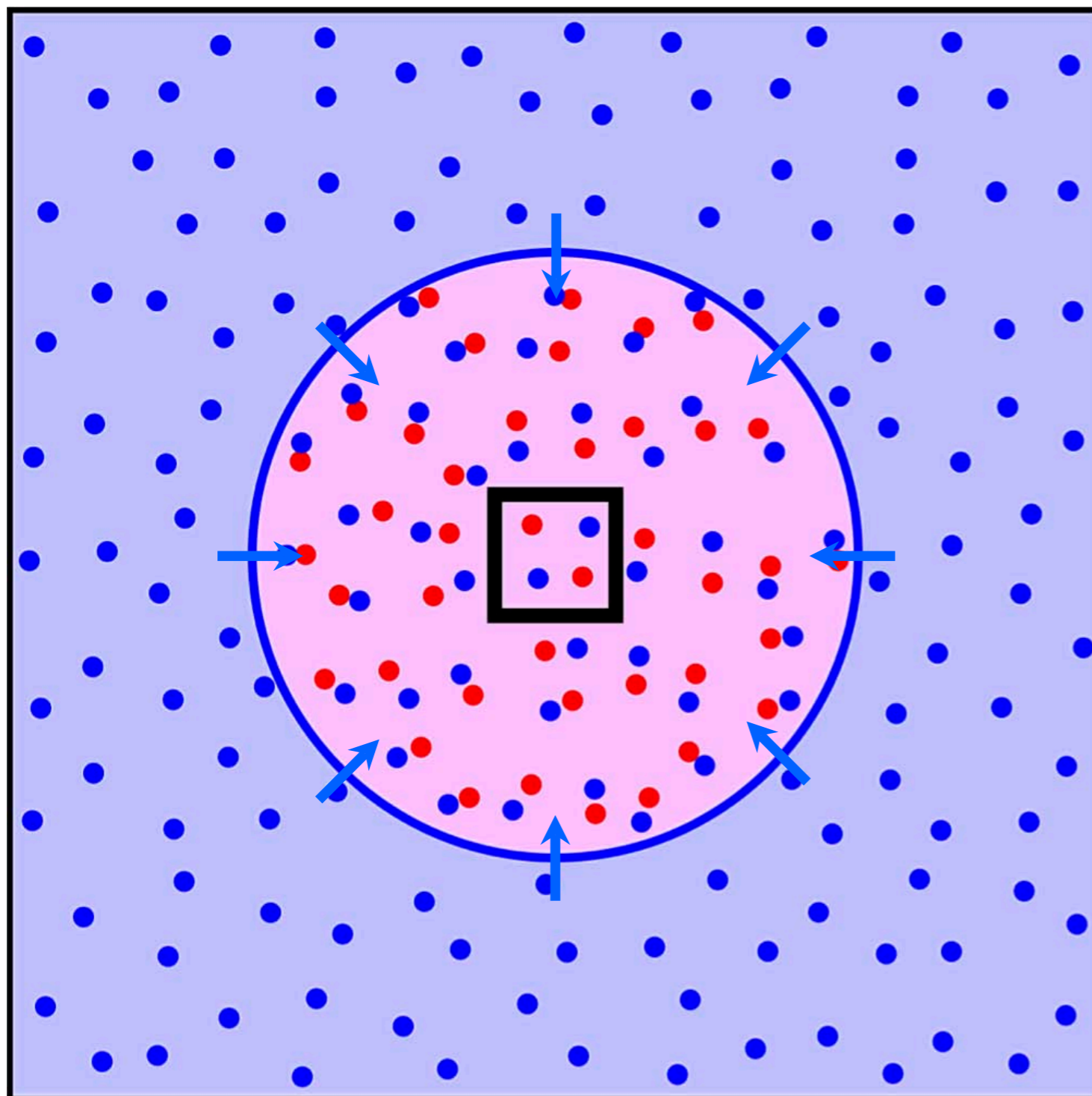
\mathcal{R} big



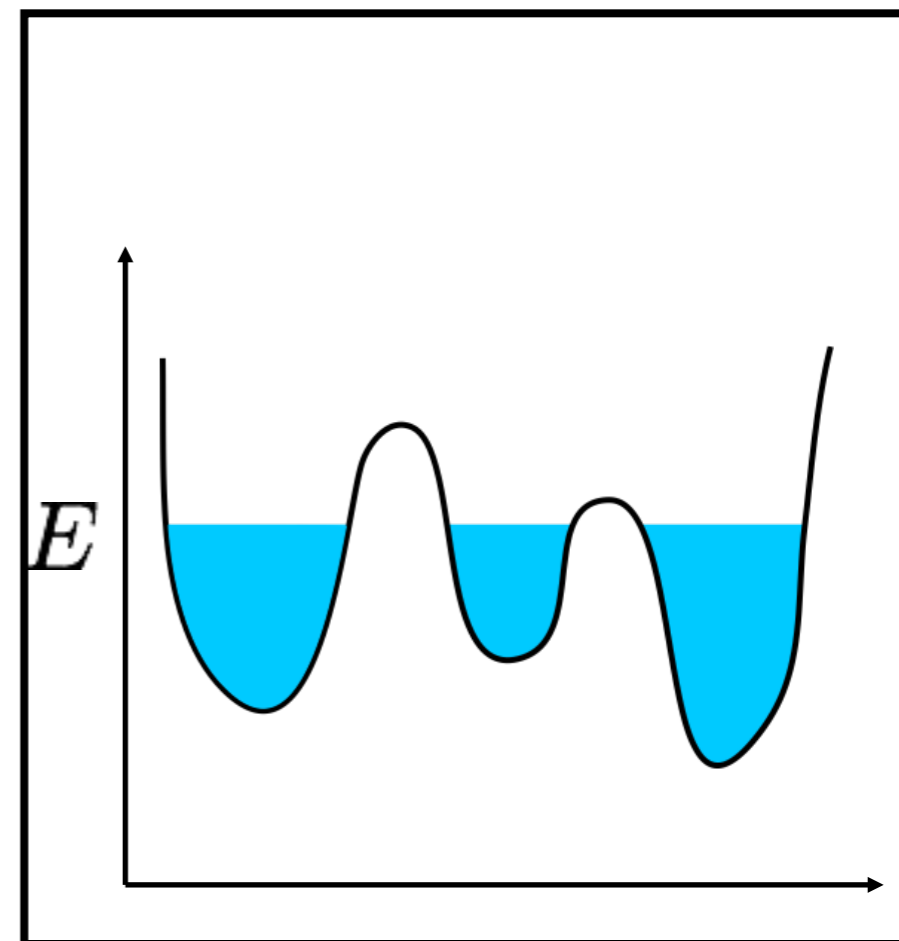
Overlap: Low

$$\hat{q}_{12} \sim 0$$

Local free-energy landscape



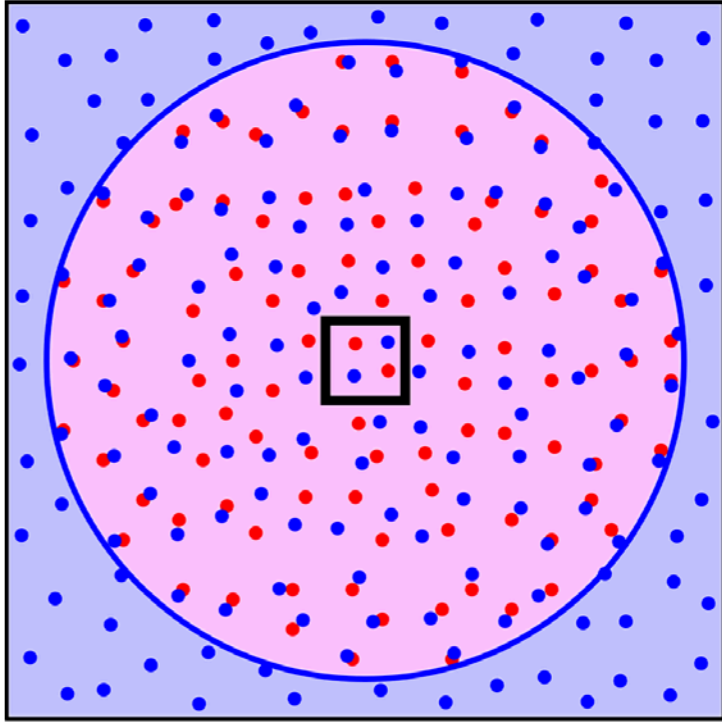
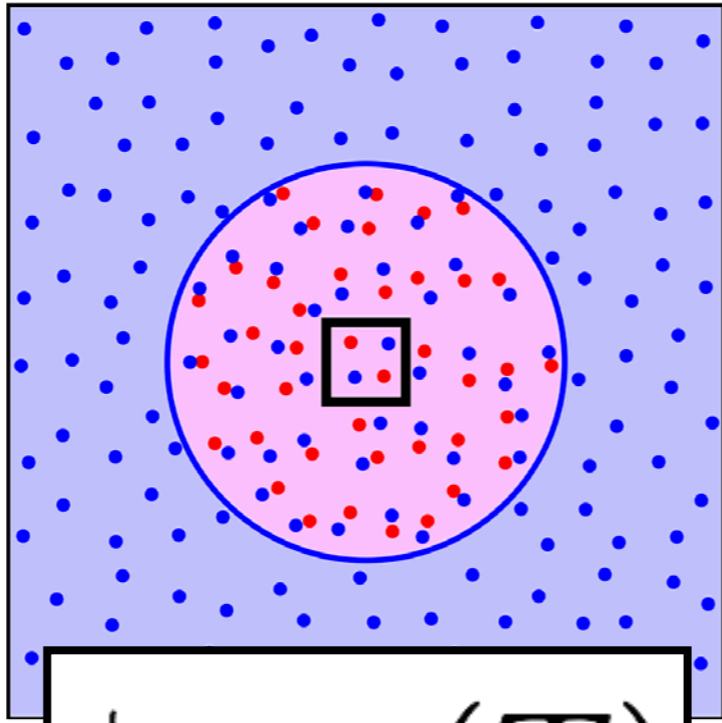
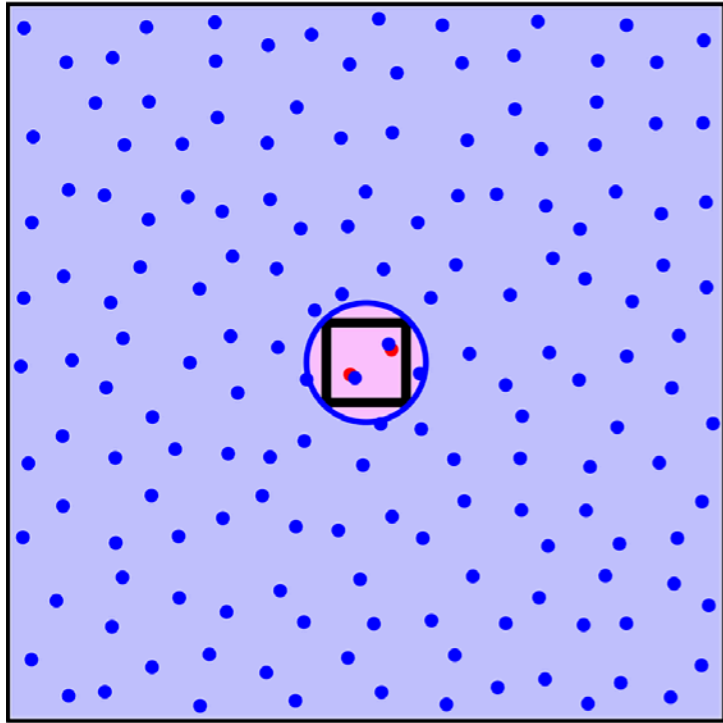
\mathcal{R} intermediate



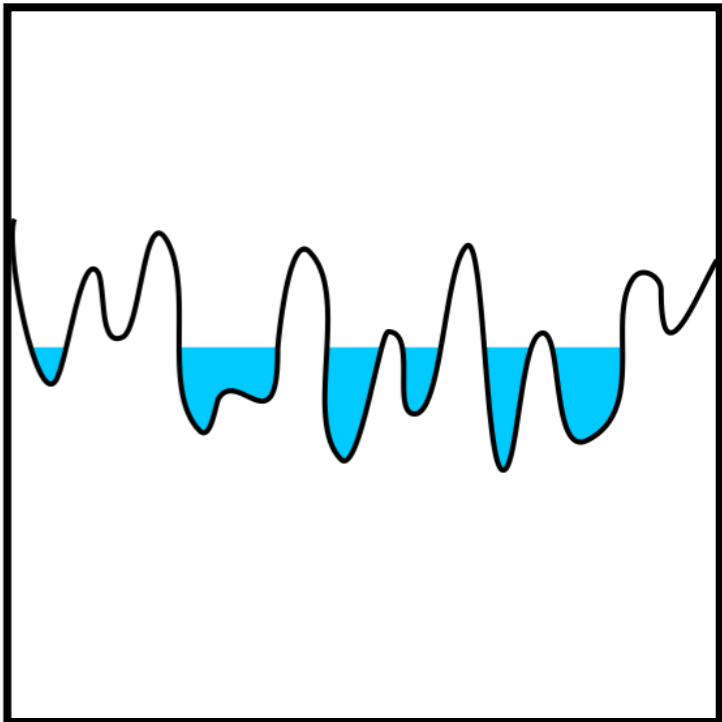
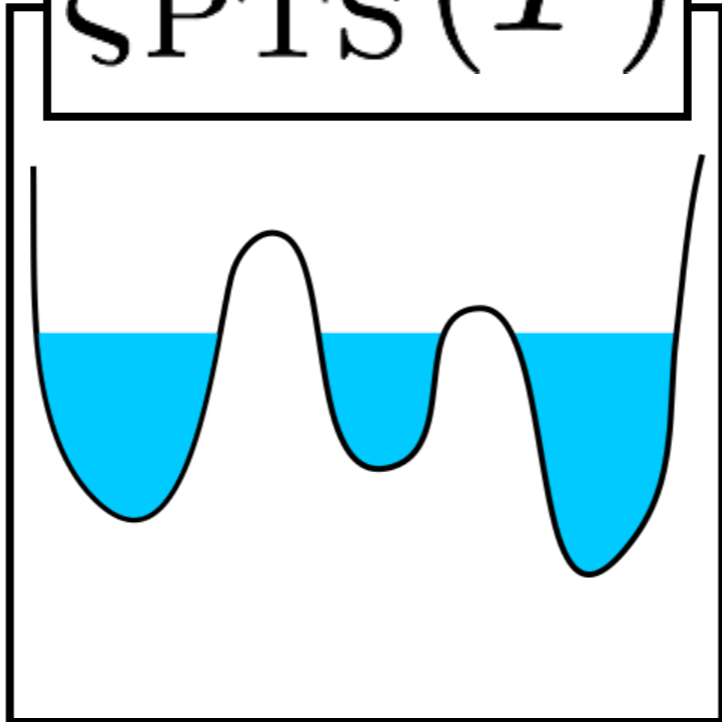
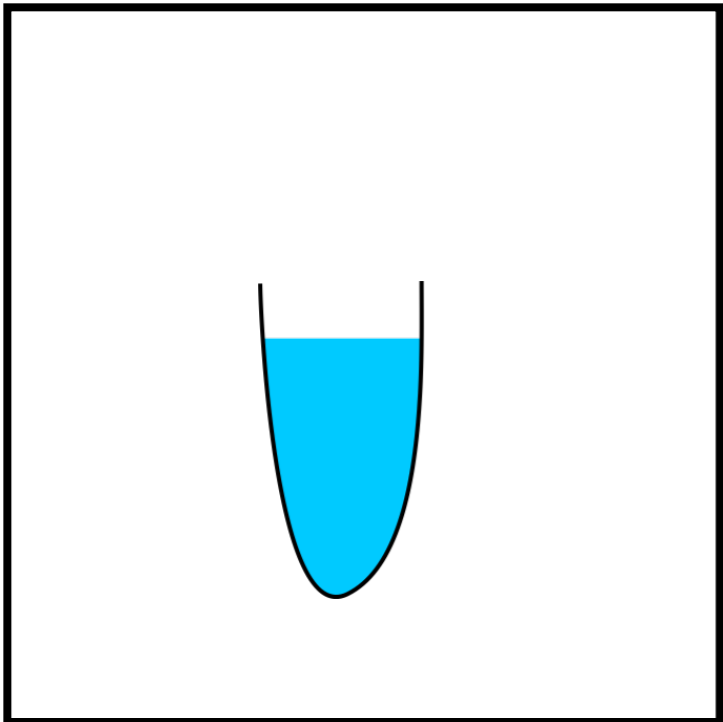
Overlap: Medium

$$\hat{q}_{12} \sim 0.5$$

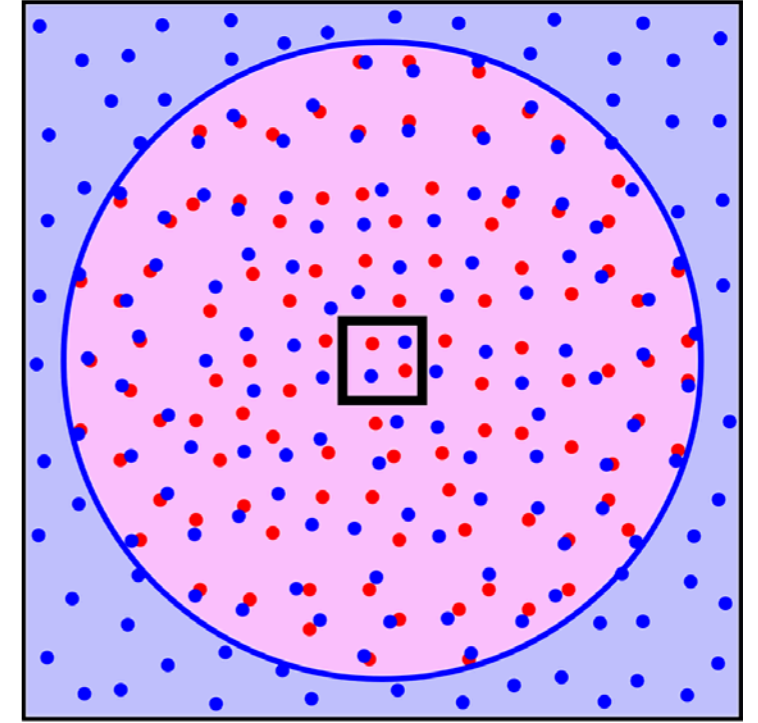
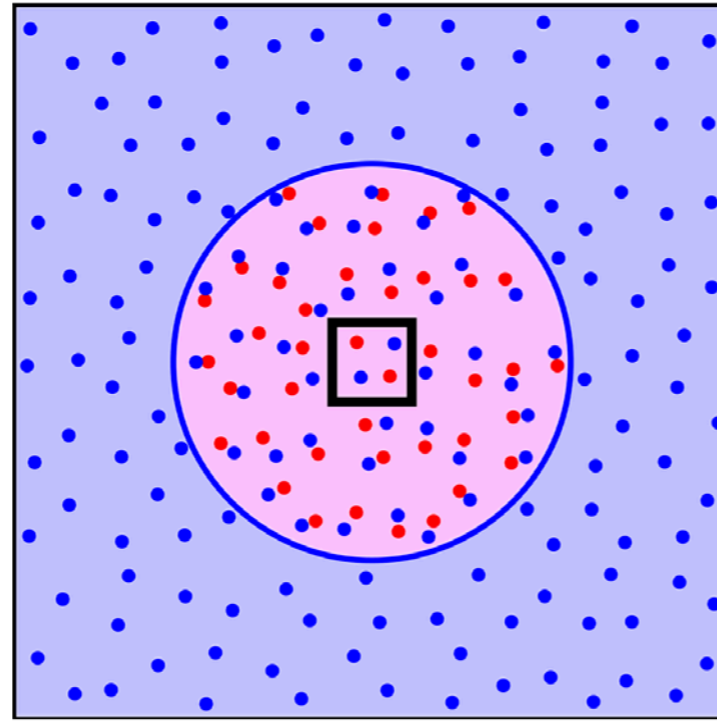
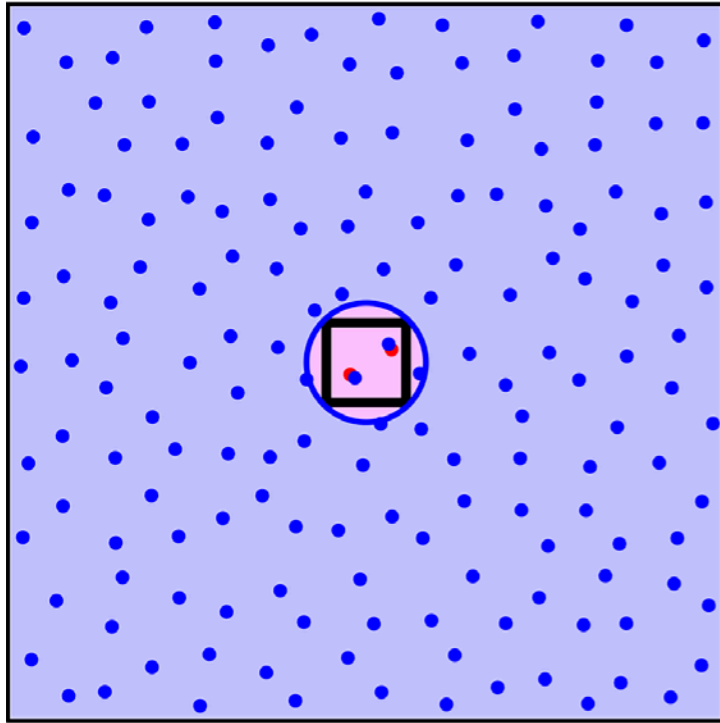
Local free-energy landscape



$\xi_{\text{PTS}}(T)$



Local free-energy landscape



Is this real?

Can we “see” it?



Simulations!

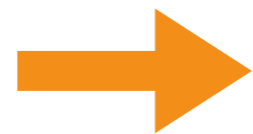
Measuring point-to-set

No reliable measurements for a decade...

2 obstacles:

glassy slowdown

$$t_{\text{eq}} \sim e^{\xi^p}$$



long computational time

confining cavity



even slower...

Measuring point-to-set

Parallel-tempering

...with this

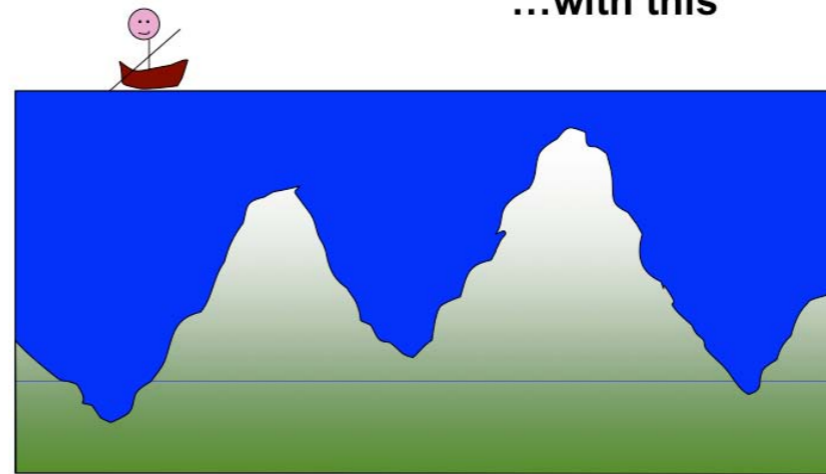
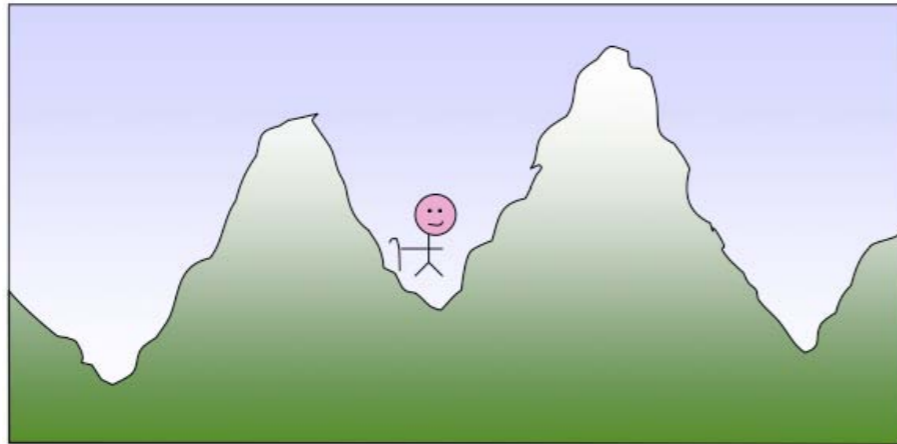


Figure credit:
Daan Frenkel

2 obstacles:

glassy slowdown

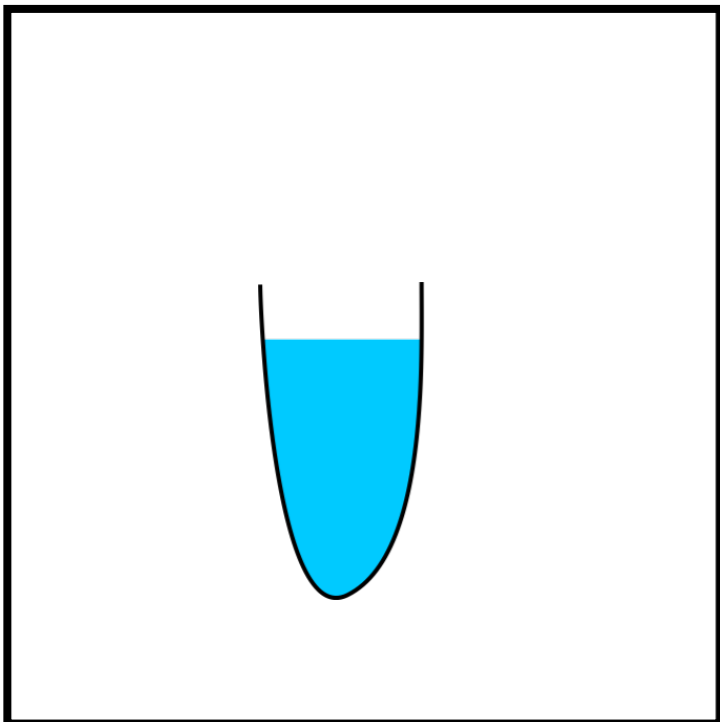
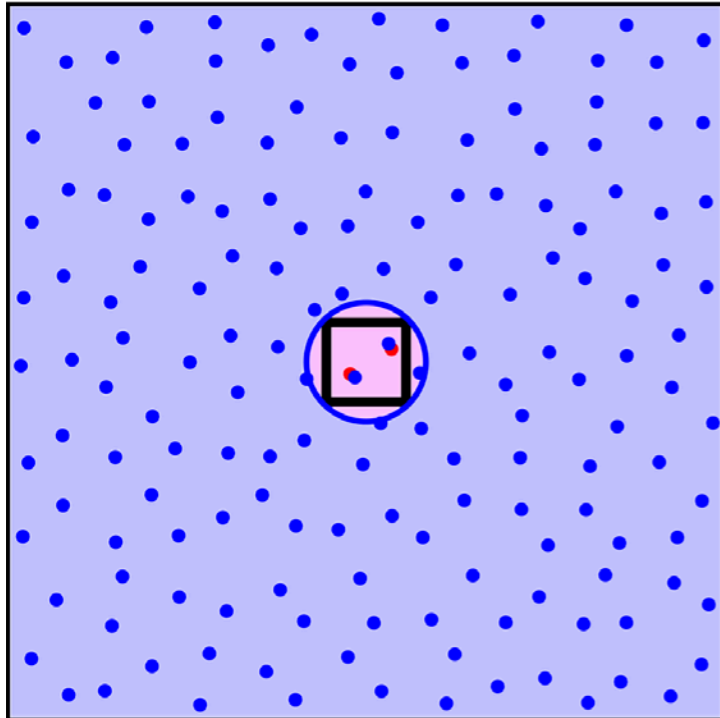
$$t_{\text{eq}} \sim e^{\xi^p}$$

long computational time

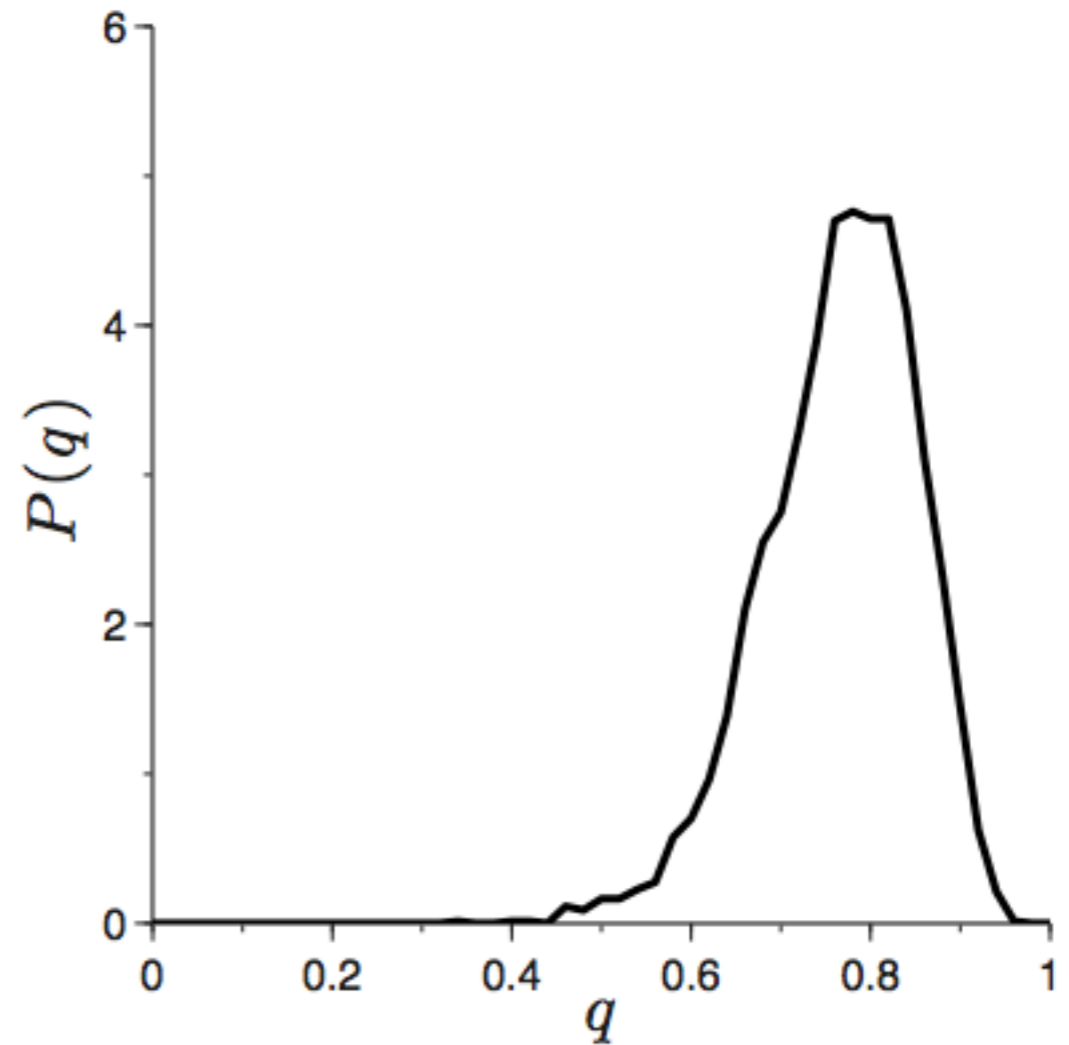
confining cavity

even slower...

Seeing landscape

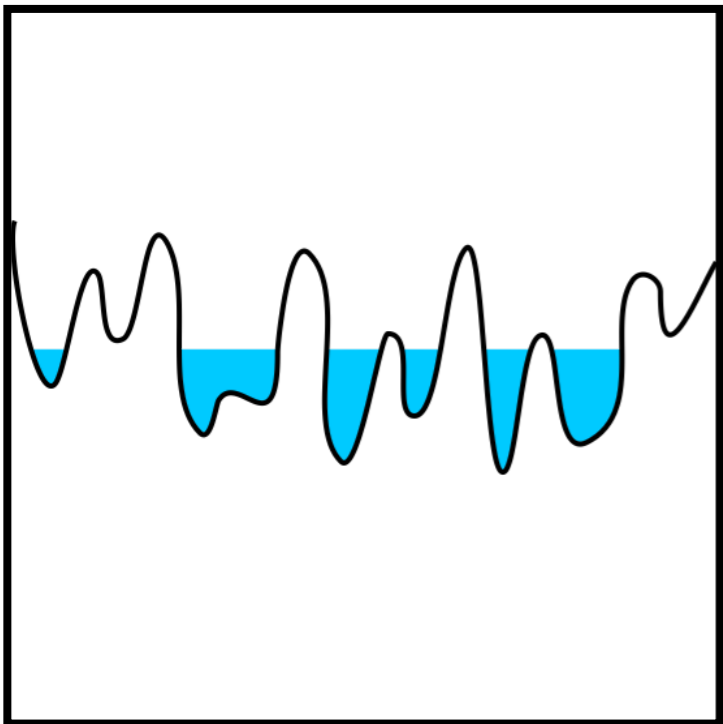
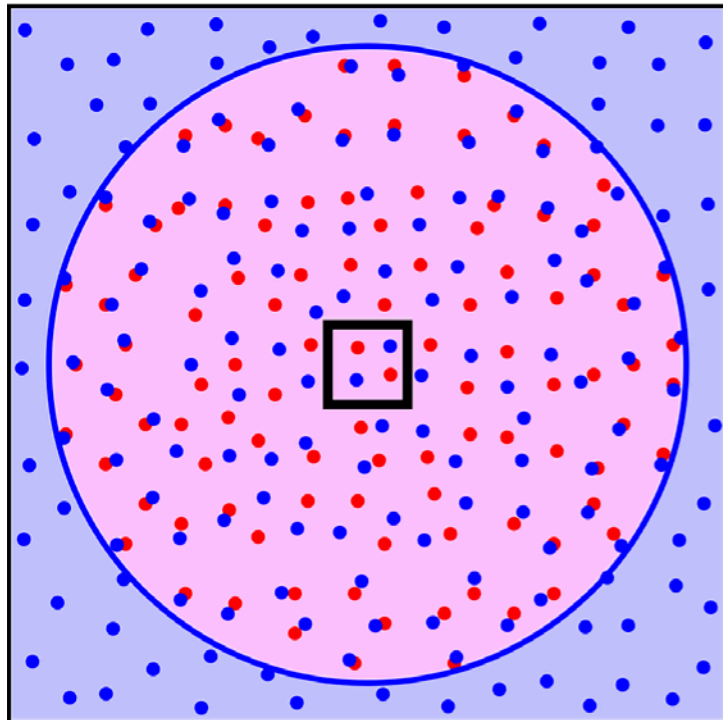


Binary Lennard-Jones liquid (Kob-Andersen @ $T = 0.51$)

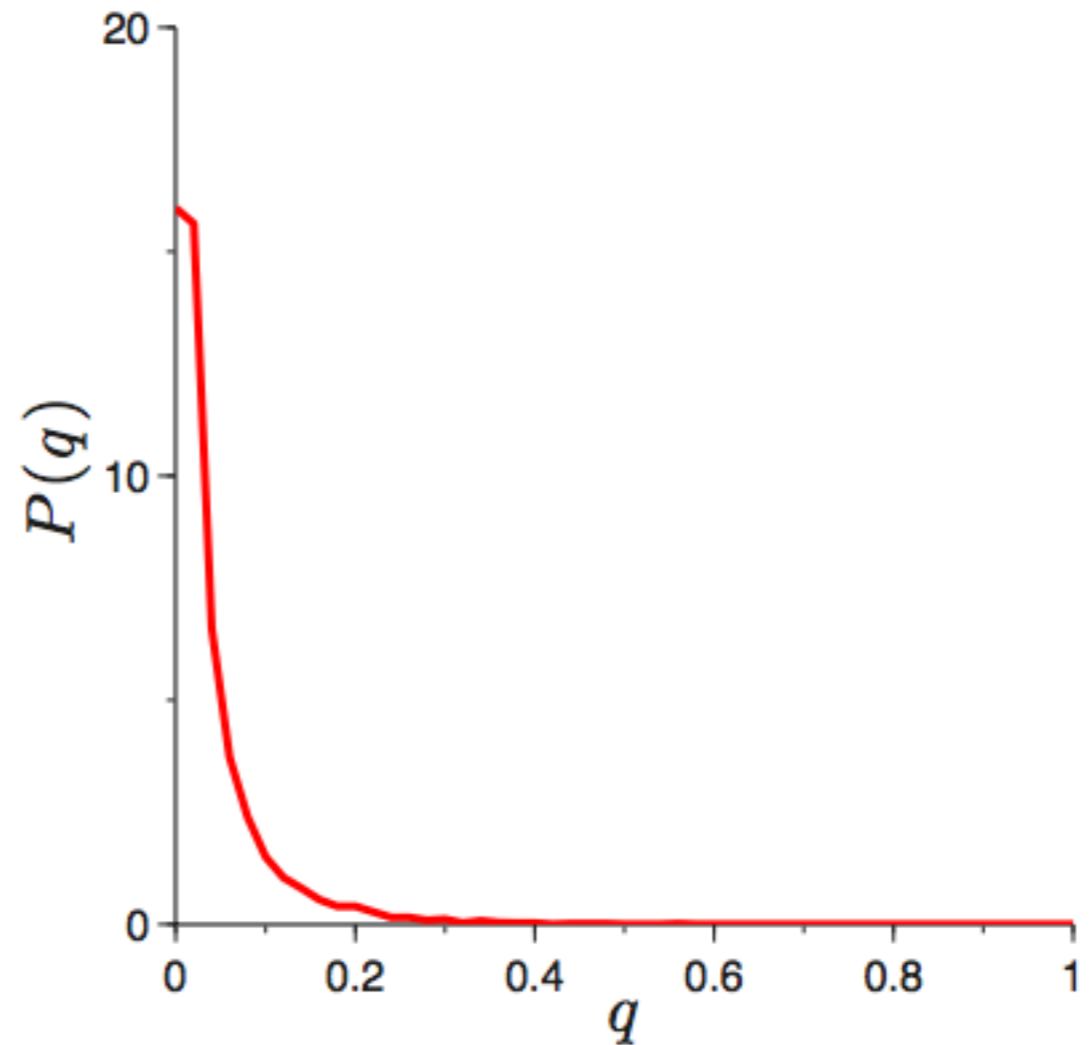


$$\mathcal{R} \ll \xi_{\text{PTS}}$$

Seeing landscape

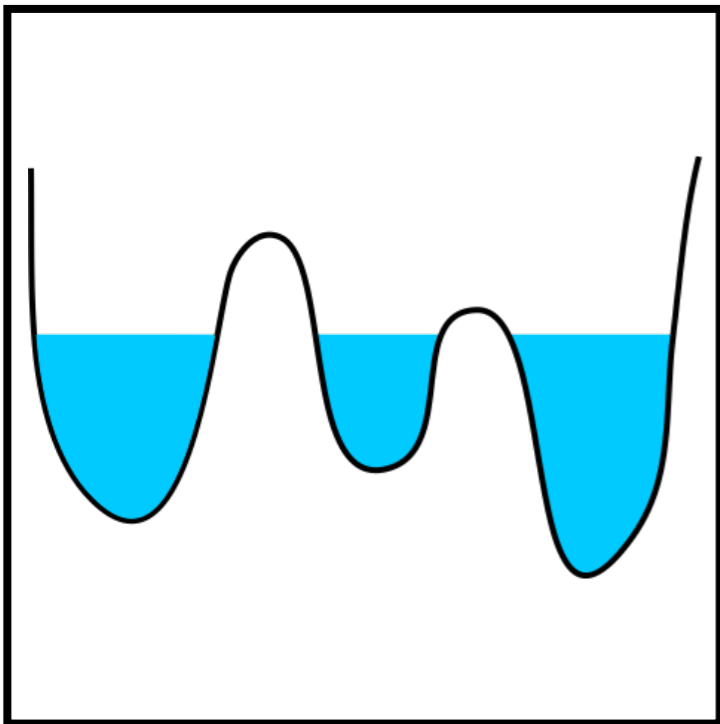
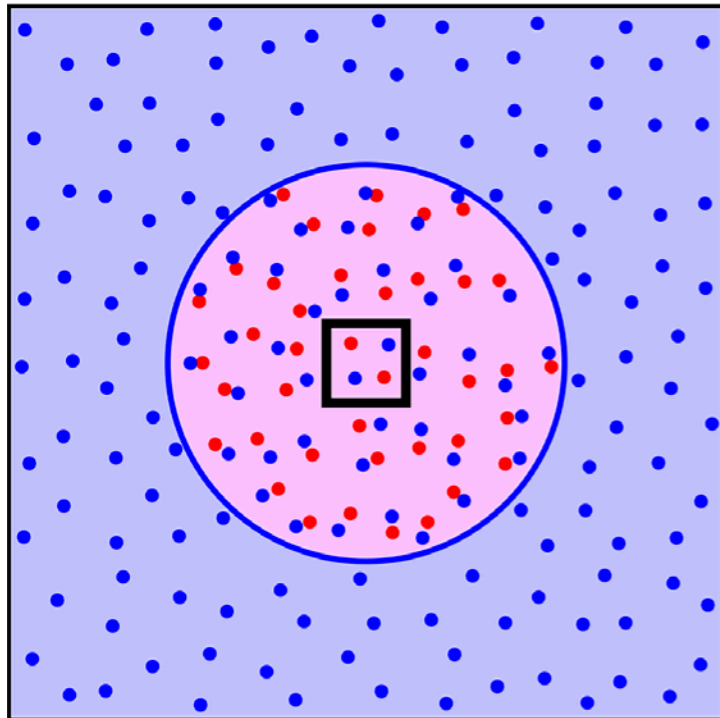


Binary Lennard-Jones liquid (Kob-Andersen @ $T = 0.51$)

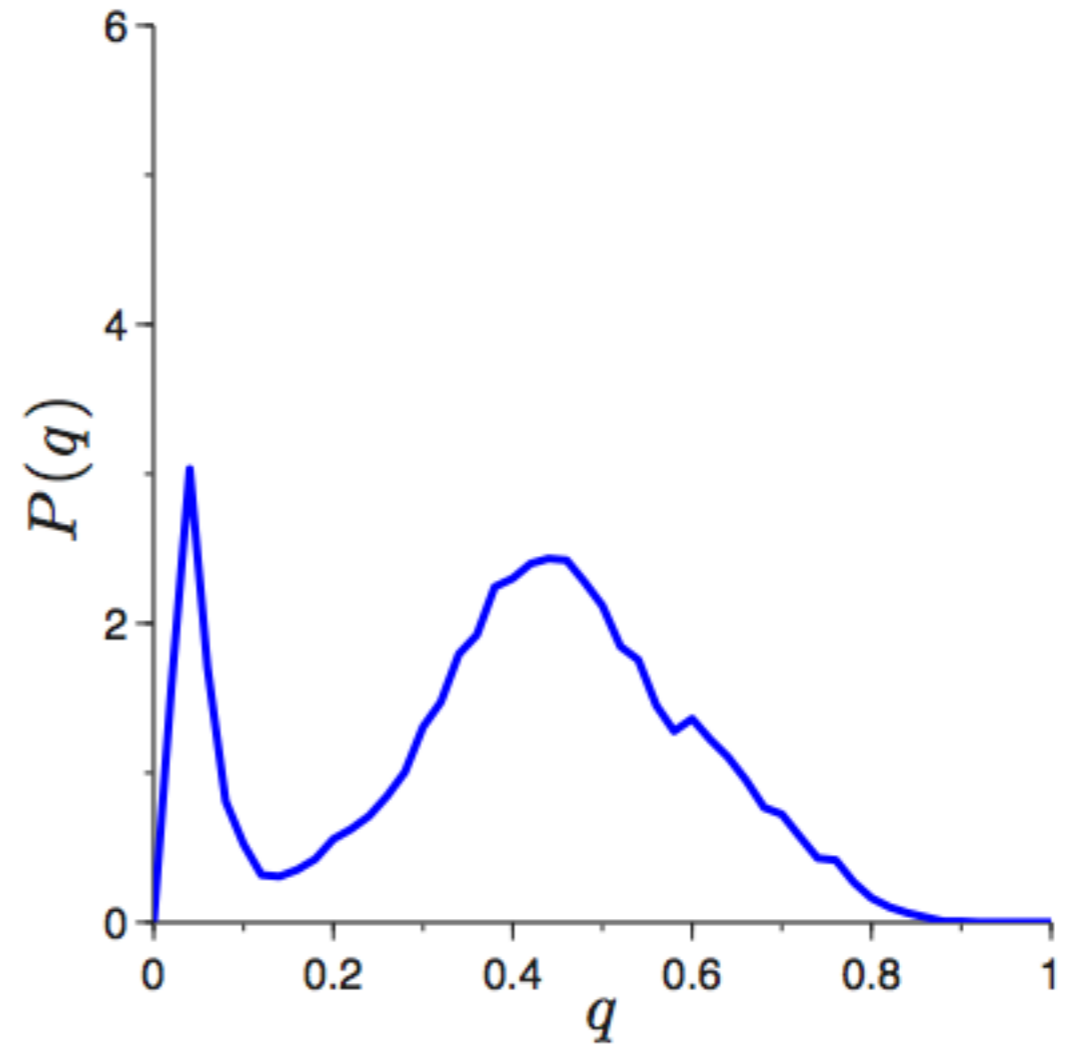


$$\mathcal{R} \gg \xi_{\text{PTS}}$$

Seeing landscape

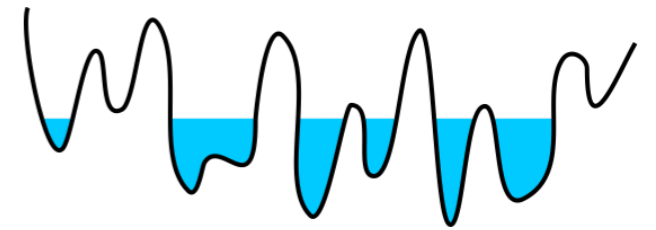
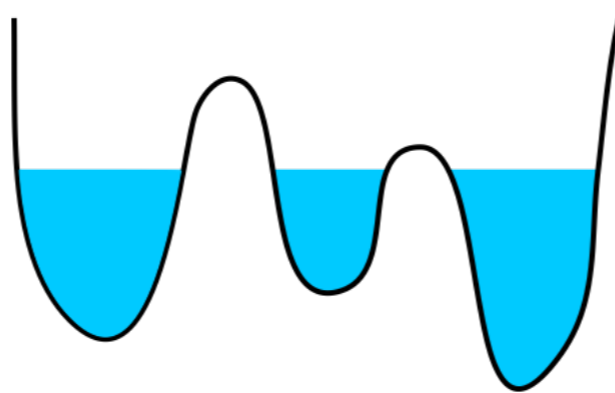
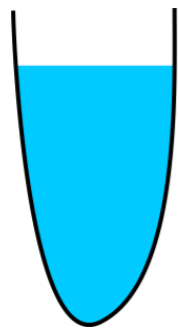
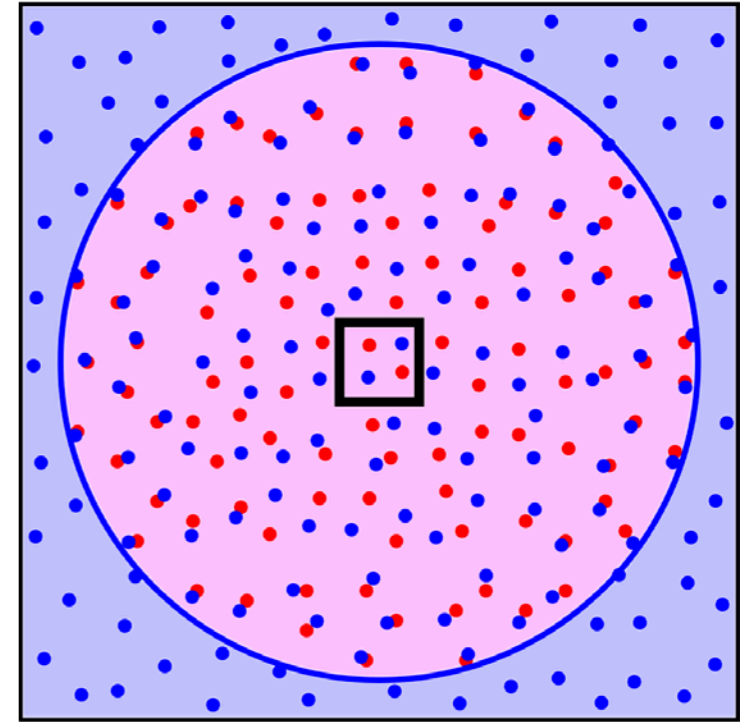
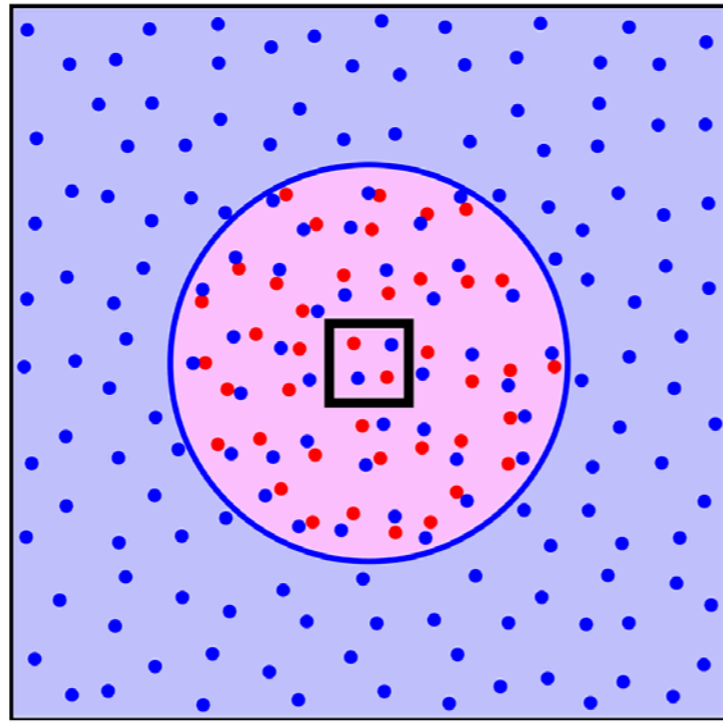
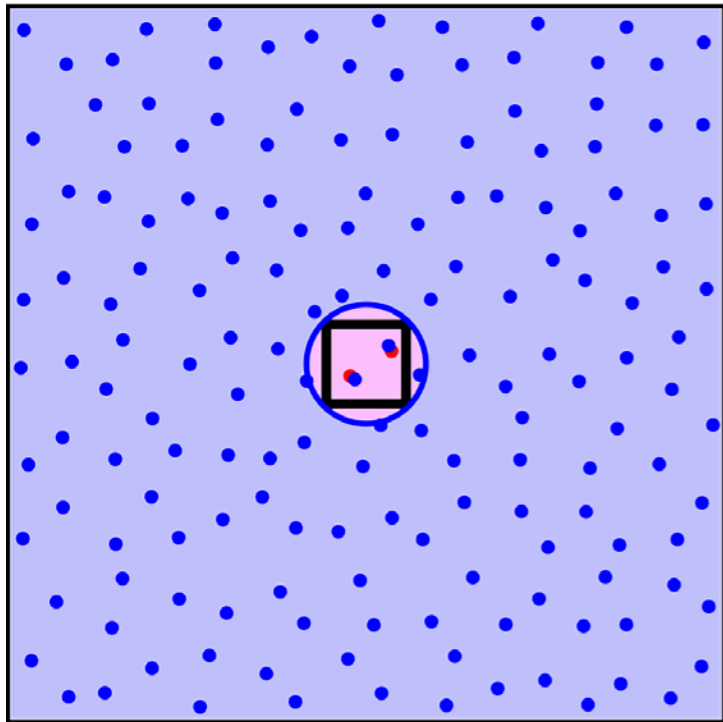


Binary Lennard-Jones liquid (Kob-Andersen @ $T = 0.51$)



$$\mathcal{R} \sim \xi_{\text{PTS}}$$

Seeing landscape



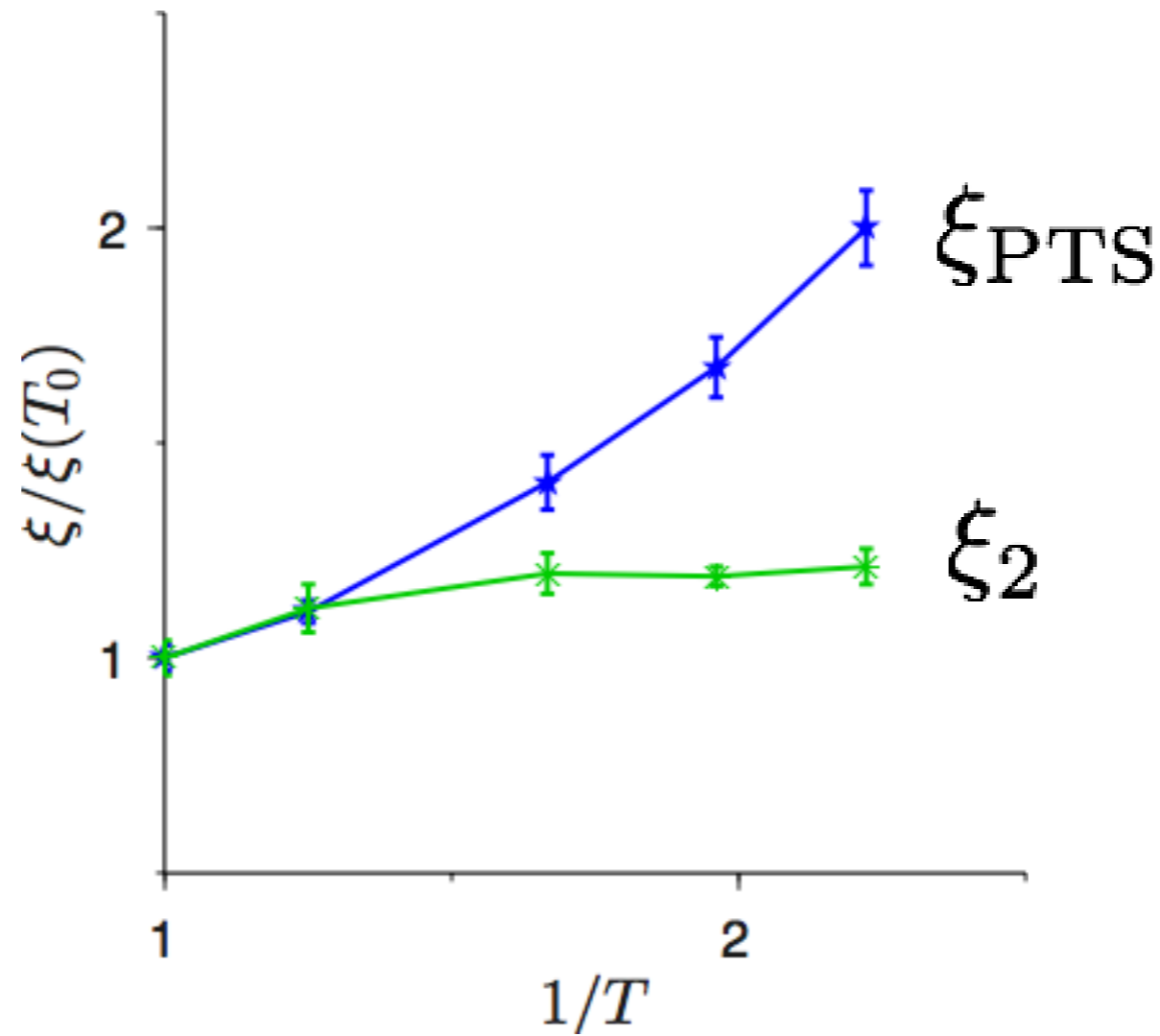
$\mathcal{R} \lesssim \xi_{\text{PTS}}$ governed by “**amorphous order**”

Defining glassiness

Rarefaction of metastable states

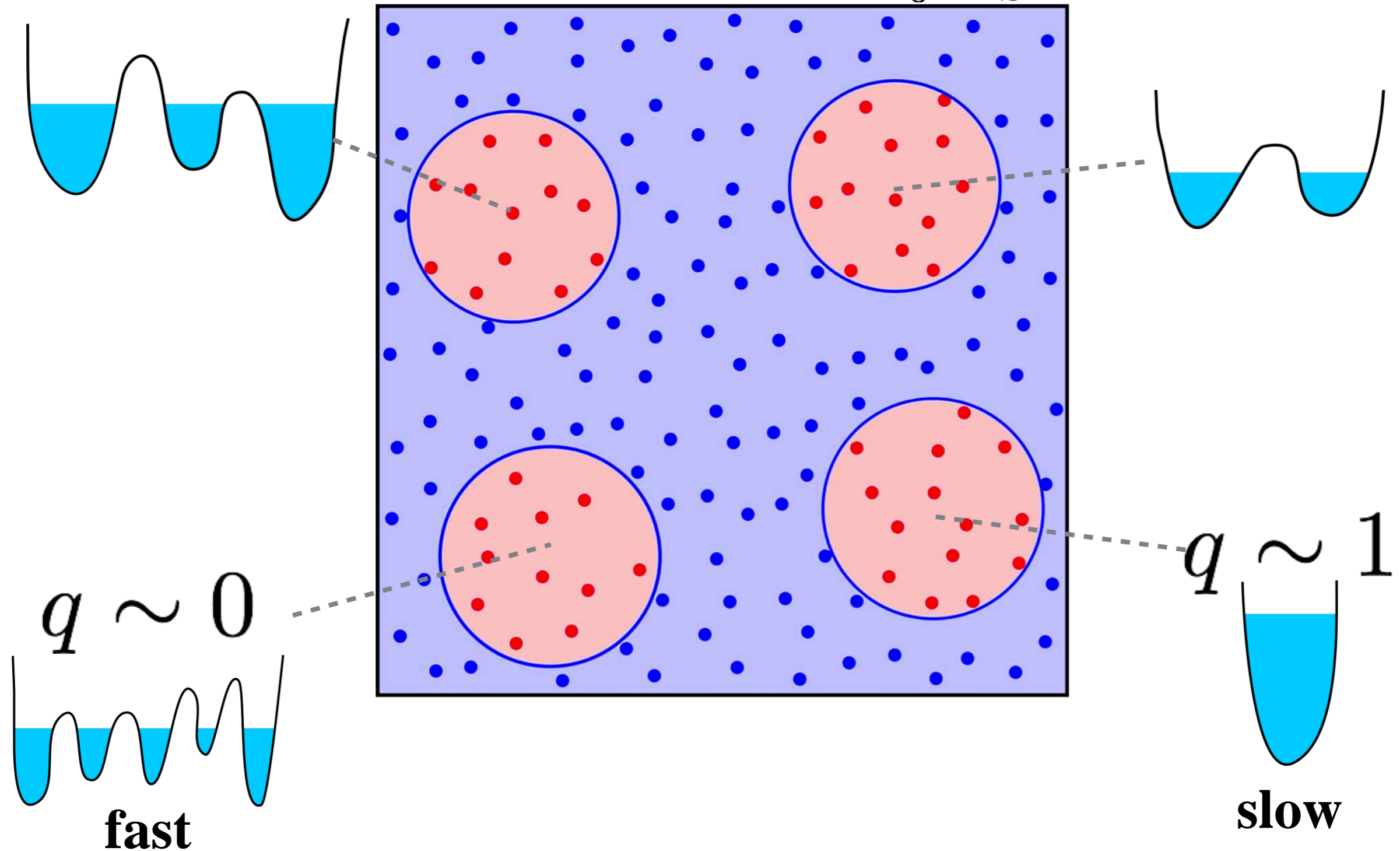
$$\frac{\xi_{\text{PTS}}}{\xi_2}$$

defines glassiness!



Statics and dynamics

Probe with $\mathcal{R} = \xi_{\text{PTS}}$



[Hocky *et al.*, PRL (2014)]

[Charbonneau-Dyer-Lee-Yaida, JSTAT (2016)]

**Rarefaction of metastable states
in a complex landscape**



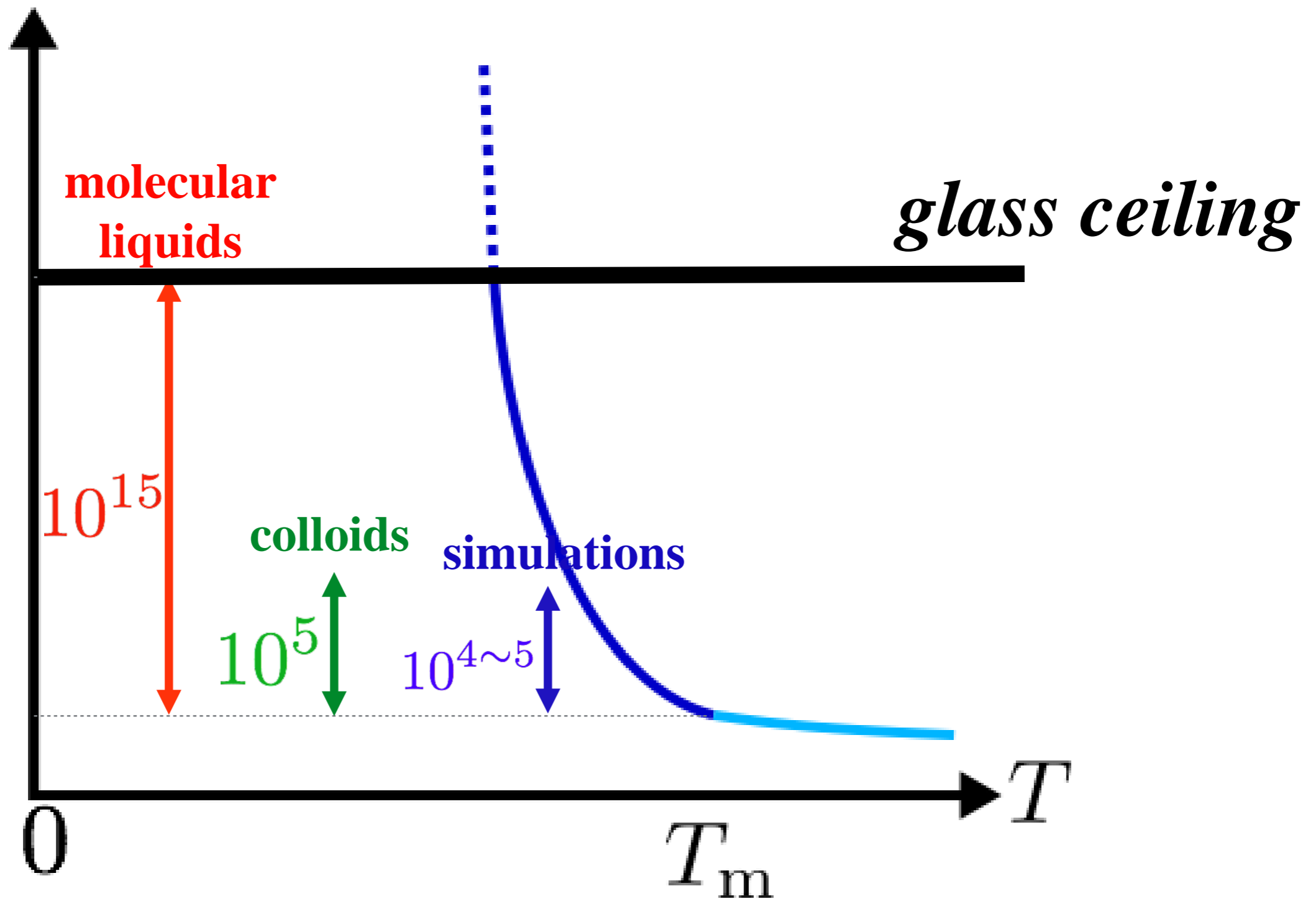
**point-to-set
correlations**

Dramatic slowdown of glassy systems

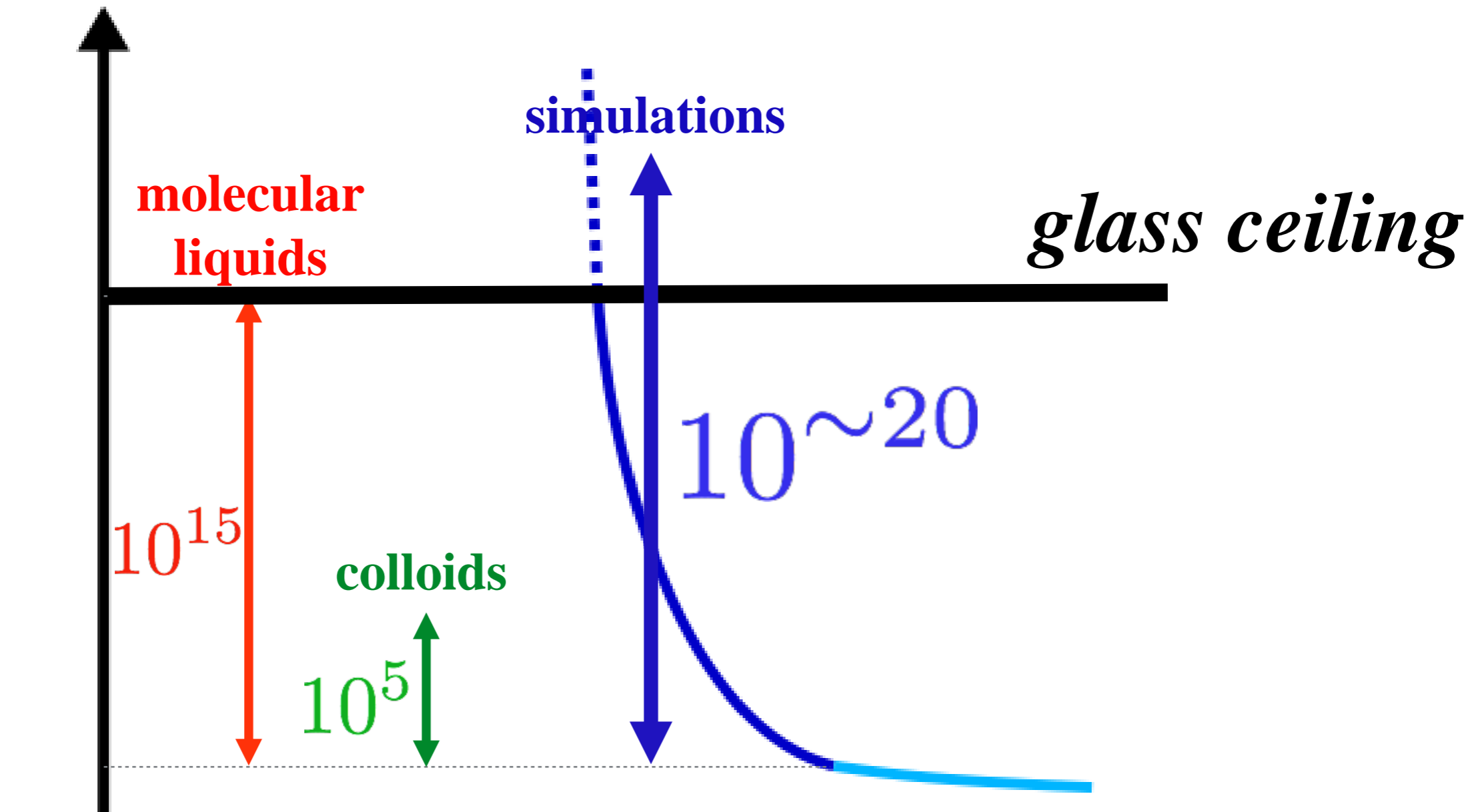


3 Beyond the glass ceiling

Breaking the glass ceiling

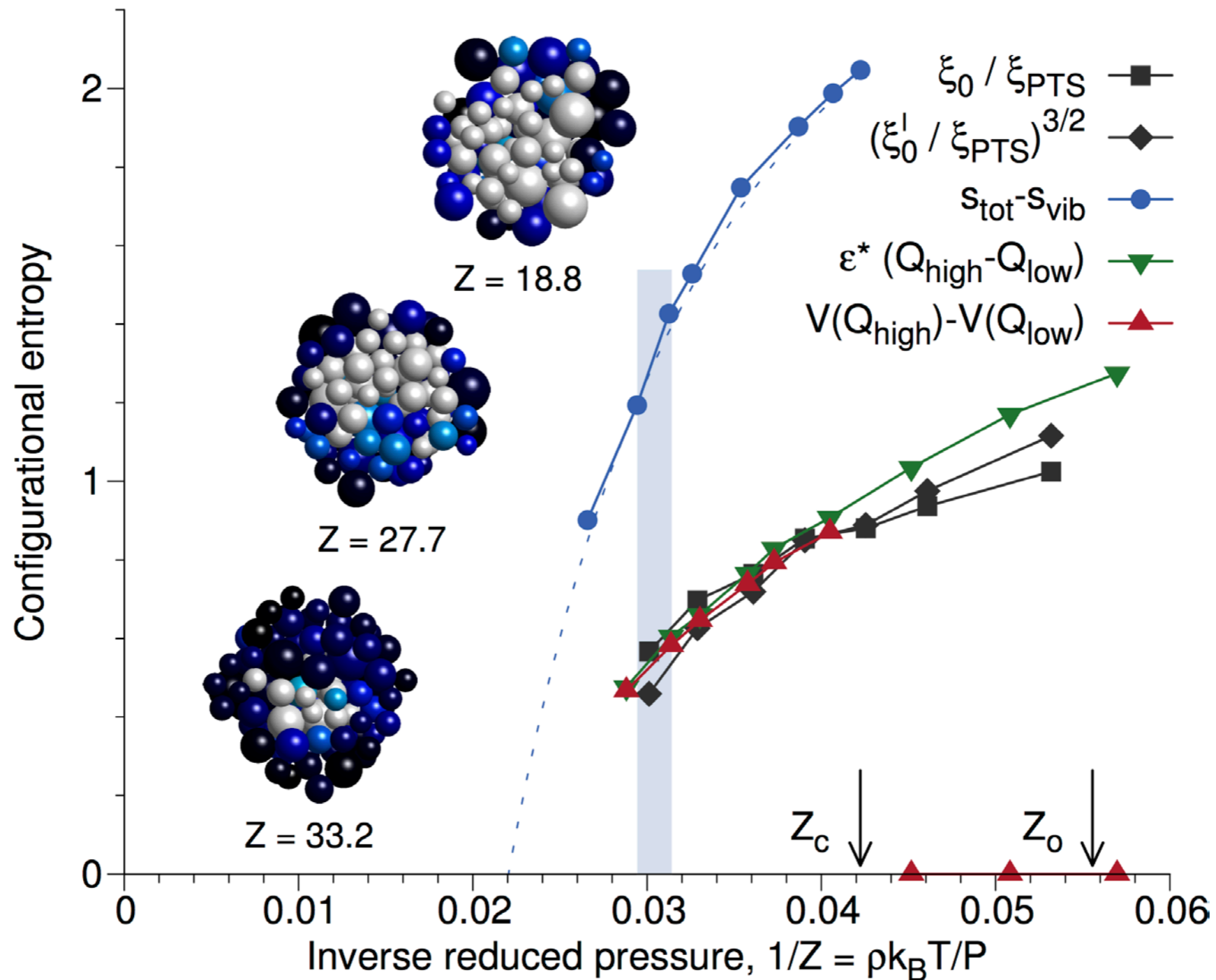


Breaking the glass ceiling



Tuning both algorithm AND system
WHILE also retaining glassy physics

Breaking the glass ceiling

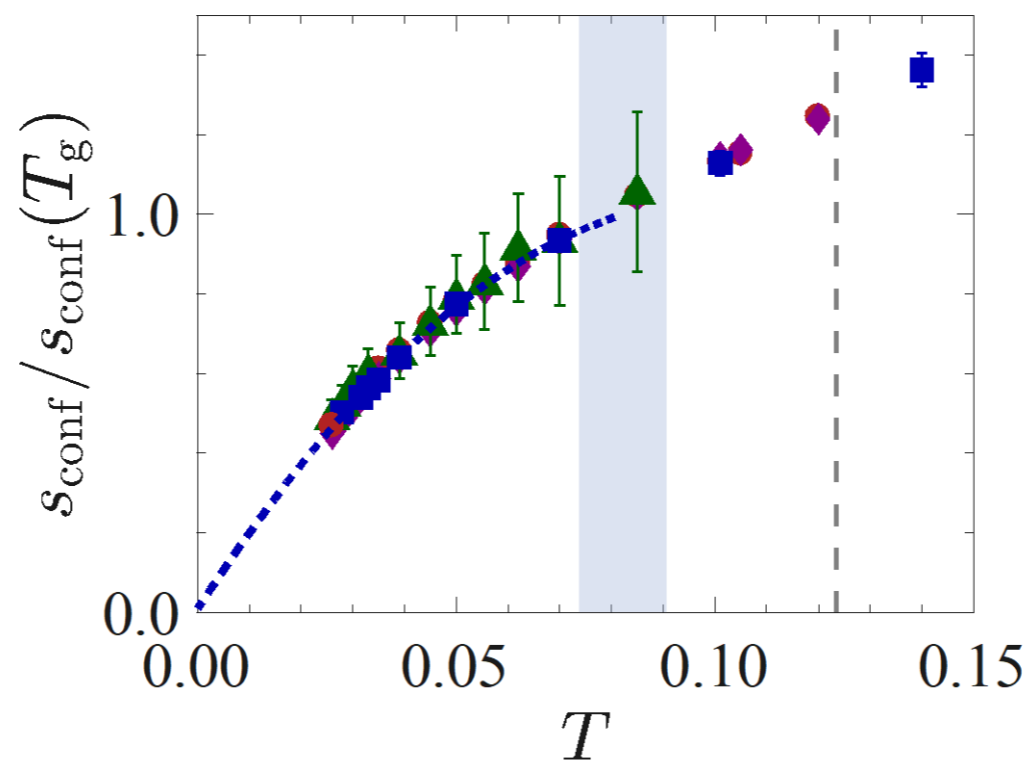


Rarefaction of metastable states **keeps driving** glassy slowdown

Breaking the glass ceiling

Difference between glassy liquids in
 $d = 2$ and $d \geq 3$

- [Flenner-Szamel, Nat. Comm. (2015)]
- [Vivek-Kelleher-Chaikin-Weeks, PNAS (2017)]
- [Illing-Fritschi-Kaiser-Klix-Maret-Keim, PNAS (2017)]

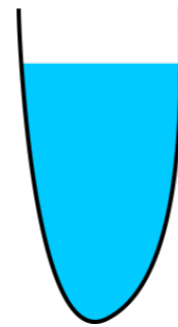


$$T_K = 0 \text{ in } d = 2$$

Summary

Glassy slowdown is cool and puzzling

Point-to-set correlations capture **amorphous order** determined by surroundings for $\mathcal{R} \lesssim \xi_{\text{PTS}}$



Rarefaction of metastable states drives $\frac{\xi_{\text{PTS}}}{\xi_2} \nearrow$ & $\eta \nearrow \nearrow$

Local amorphous order is a good predictor of local dynamics

Point-to-set correlation length **universally** grows

Recently tested **beyond the glass ceiling**