

Self-regulated star formation: concepts and computations

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Outline

- Self-regulated SF in turbulence-dominated systems: concept and formulation
- Computational results: starbursts and outer disks
- Self-regulation: radiation effects
- Computational results: star-forming clouds with IR radiation

- I. Self-regulation: the concept
- II. Starburst and outer-disk simulations
- III. Continuum radiation effects
- IV. RHD simulations of SF clouds

Turbulent driving and dissipation

Consider system of mass M , size L^3 , turbulence v

- Assume SF feedback momentum/mass is p_*/m_*

- Momentum input rate is $\dot{p}_{driv} = \frac{p_*}{m_*} \dot{M}_*$

- Momentum dissipation rate is $\dot{p}_{diss} \sim \frac{vM}{t_{dyn}} \sim \frac{v^2 M}{L}$

- Balancing, $\dot{M}_* \sim \frac{v^2 M}{L p_*/m_*}$

- For system in dynamical equilibrium $v^2 \sim GM_{tot}/L$



$$\dot{M}_* \sim \frac{GM_{tot}M}{L^2 p_*/m_*}$$

Self-regulated
star formation

Star-forming equilibrium gaseous system supported by driven turbulence

- $M_{tot} \sim M \Rightarrow \dot{M}_* \sim \frac{v^2 M}{L p_* / m_*} \rightarrow \frac{GM^2}{L^2 p_* / m_*}$

- Can also define: $\dot{M}_* \equiv \epsilon_{ff} \frac{M}{t_{ff}} \sim \epsilon_{ff} \frac{vM}{L}$

Krumholz et al; Padoan et al; Federrath et al

ϵ_{ff} depends in principle on $\alpha_{vir} \sim (t_{ff}/t_{dyn})^2$, v/c_s , v/v_A ; small if turbulence can disperse structures before they collapse

$$v \sim \epsilon_{ff} \frac{p_*}{m_*}$$

dynam. equil.
driving=dissipation
SF efficiency definition

$$L \sim \frac{GM}{v^2}$$

dynam. equil.

$$\dot{M}_* \sim \frac{GM^2}{L^2 p_* / m_*}$$

dynam. equil.
driving=dissipation

Gas-dominated starburst disk

Ostriker & Shetty (2011)

$$\dot{M}_* \sim \frac{GM^2}{L^2 p_* / m_*} \longrightarrow \frac{\dot{M}_*}{L^2} \sim \frac{GM^2}{L^4 (p_* / m_*)}$$

⇒

$$\Sigma_{\text{SFR}} \sim \frac{G\Sigma^2}{p_* / m_*}$$

- Star formation rate per unit area in disk is
 - *independent* of details of turbulence
 - *independent* of ϵ_{ff} on small scales
- Disk thickness and internal dynamical time must adjust until momentum feedback rate matches vertical weight of ISM

Gas-dominated starburst disk

Ostriker & Shetty (2011)

$$\dot{M}_* \sim \frac{GM^2}{L^2 p_* / m_*} \longrightarrow \frac{\dot{M}_*}{L^2} \sim \frac{GM^2}{L^4 (p_* / m_*)}$$

⇒

$$\Sigma_{\text{SFR}} \sim \frac{G\Sigma^2}{p_* / m_*}$$

- As for other strongly turbulent systems, expect low ϵ_{ff} , with self-consistent v related to feedback momentum by:

$$v \sim \epsilon_{\text{ff}} \frac{p_*}{m_*}$$

- L is disk thickness H :

$$H \sim \frac{GM}{v^2} \sim \frac{GH^2\Sigma}{v^2} \Rightarrow H \sim \frac{v^2}{G\Sigma}$$

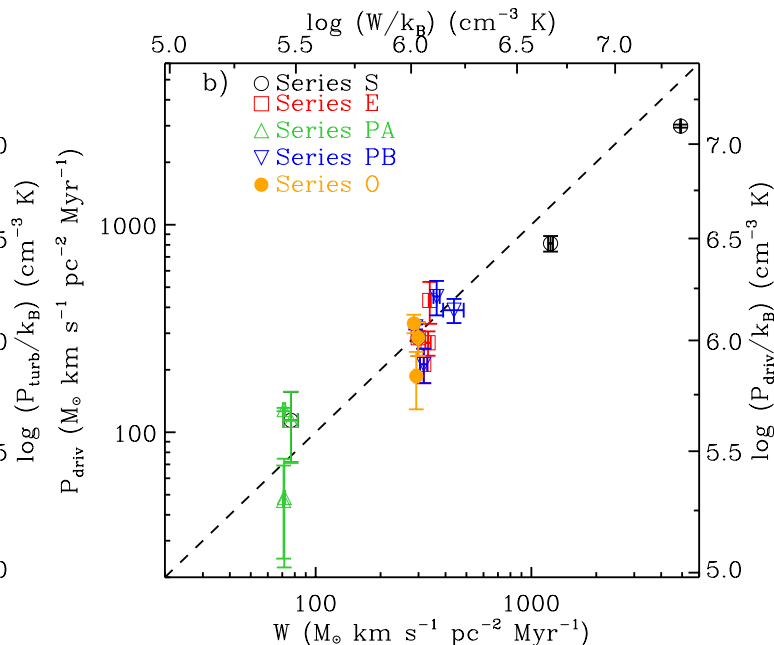
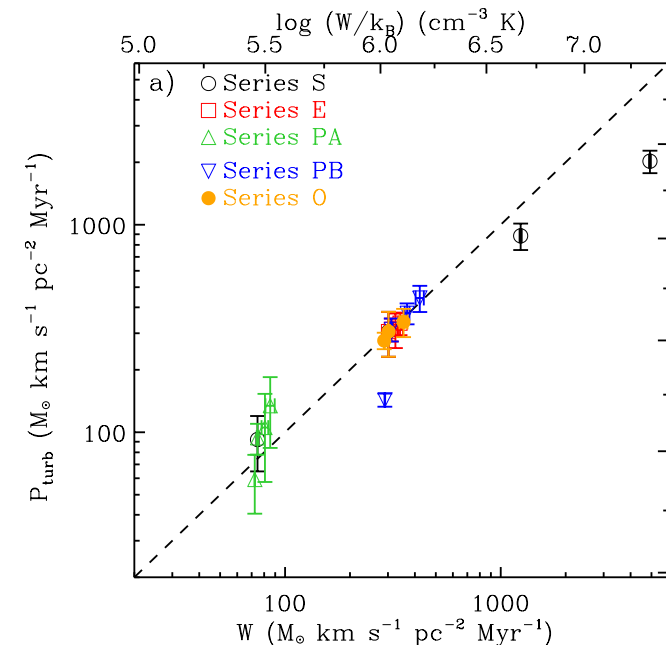
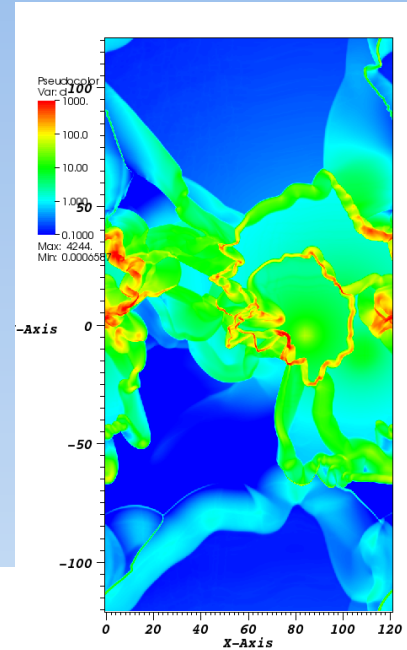
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Starburst regime simulations

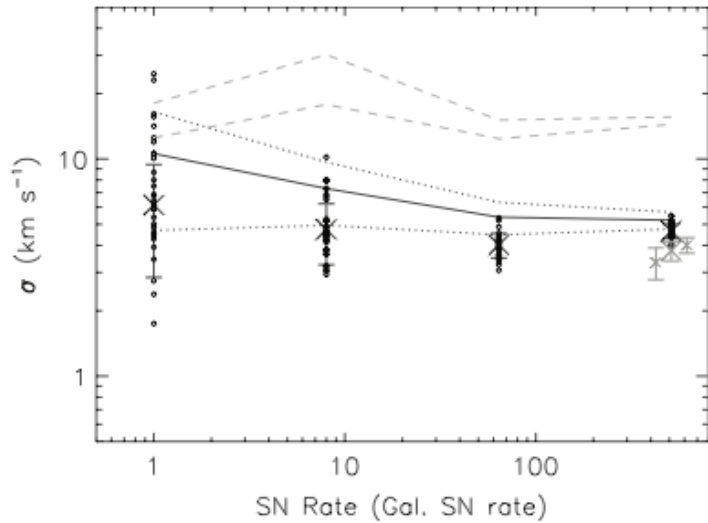
Shetty & Ostriker (2012)

- Feedback-driven, turbulence-dominated equilibrium:

- $P_{\text{turb}} \approx W \approx \pi G \Sigma^2 / 2 \approx (1/4)(p_*/m_*) \Sigma_{\text{SFR}}$
- Simulation yields $\epsilon_{\text{ff}}(\rho_0) \sim 0.005-0.01$ insensitive to other conditions
- Simulation yields $v_z \sim 5-10 \text{ km/s} \propto p_*/m_*$



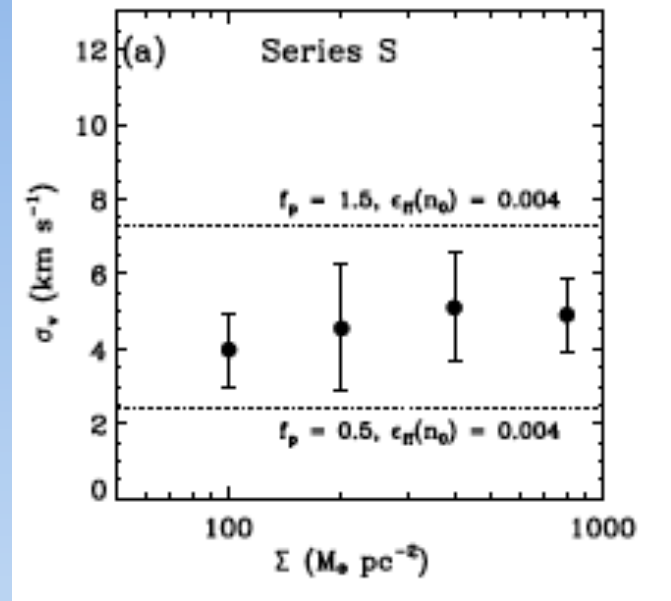
Supernova-driven turbulence



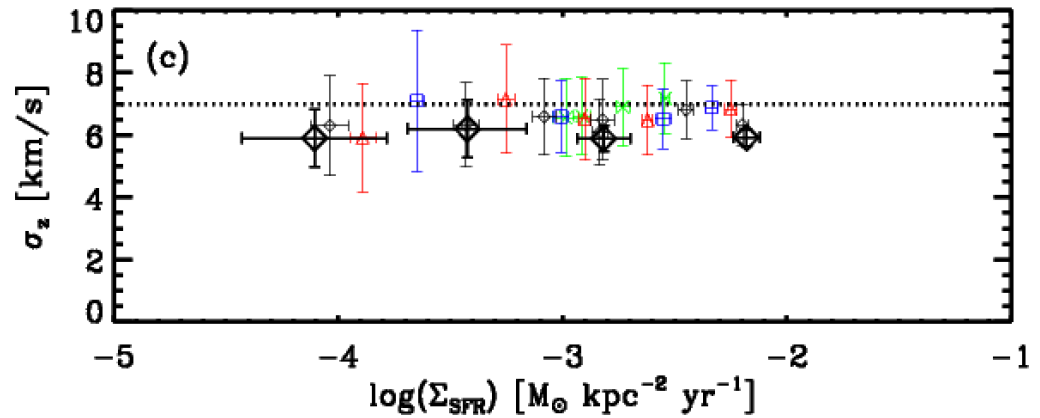
Energy-driven Joung et al (2009)

Radiative SN remnants (Cioffi et al 1988; Blondin et al 1998, Thornton et al 1998):

$$\frac{p_*}{m_*} \approx 3000 \text{ km s}^{-1} \left(\frac{E_{\text{SN}}}{10^{51} \text{ erg}} \right)^{0.94} \left(\frac{n_0}{1 \text{ cm}^{-3}} \right)^{-0.12} \left(\frac{m_*}{100 M_\odot} \right)^{-1}$$



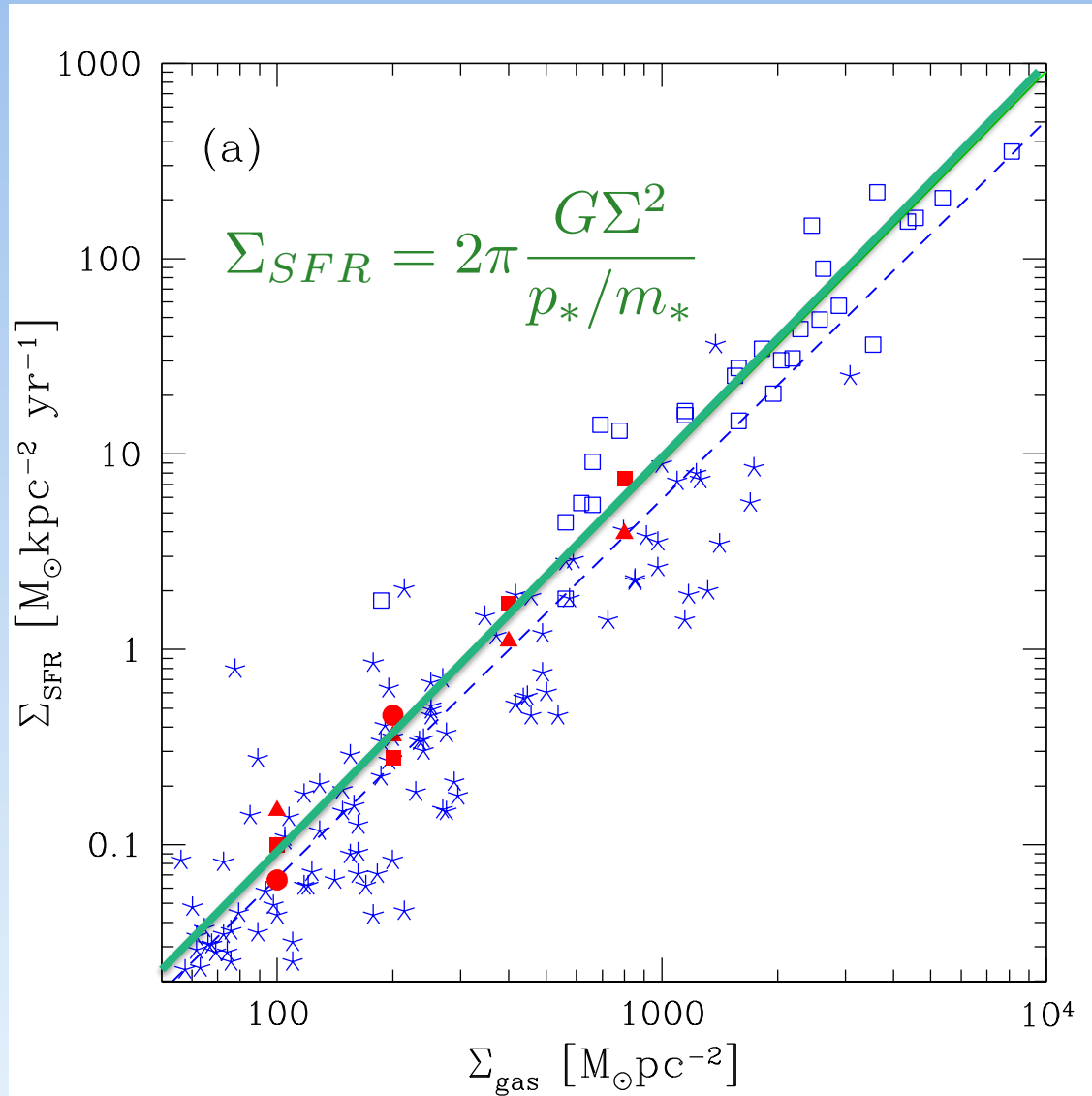
Momentum-driven



Kim et al (2011, 2013)

Starburst regime

Ostriker & Shetty (2011)



Data from Genzel et al (2010) sample

Self-regulation in externally-confined disks

Atomic-dominated regions of galactic disks are confined by the vertical *stellar* rather than vertical *gas* potential:

- Gravitational free-fall time (gas):

$$t_{\text{ff}} = \left(\frac{3\pi}{32G\rho} \right)^{1/2} = 43\text{Myr} \left(\frac{n_H}{1 \text{ cm}^{-3}} \right)^{-1/2}$$

- Dynamical crossing time:

$$t_{\text{dyn}} = \frac{1}{(4\pi G\rho_*)^{1/2}} = 13\text{Myr} (\rho_*/0.1M_\odot \text{ pc}^{-3})^{-1/2}$$

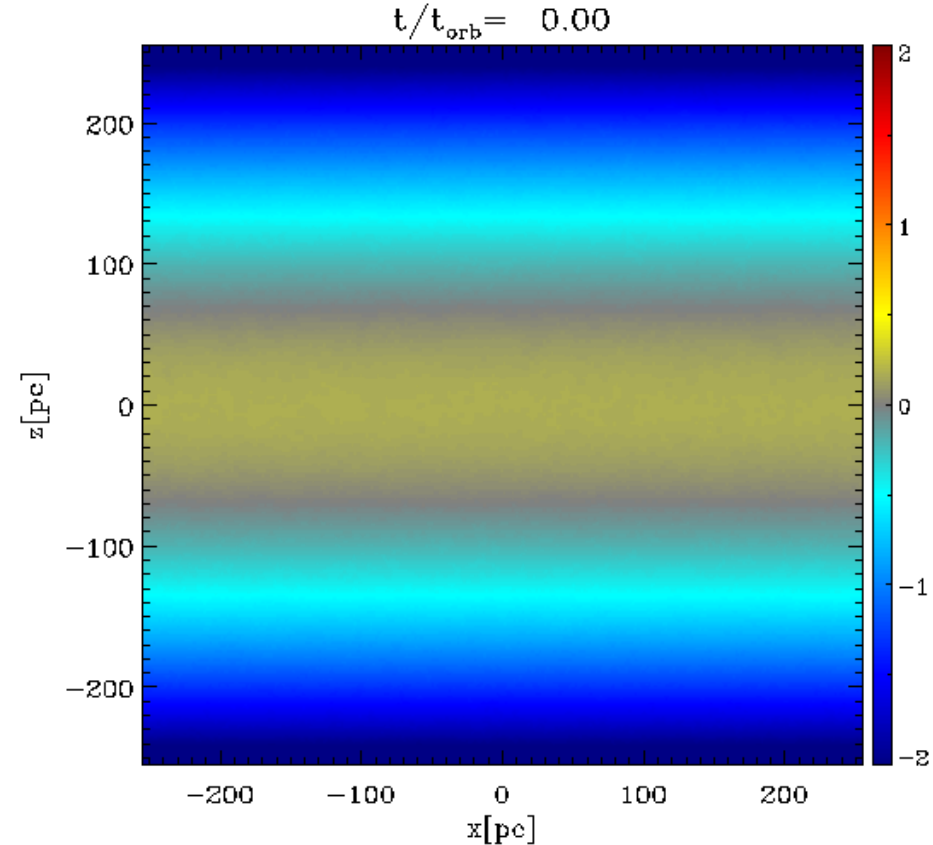
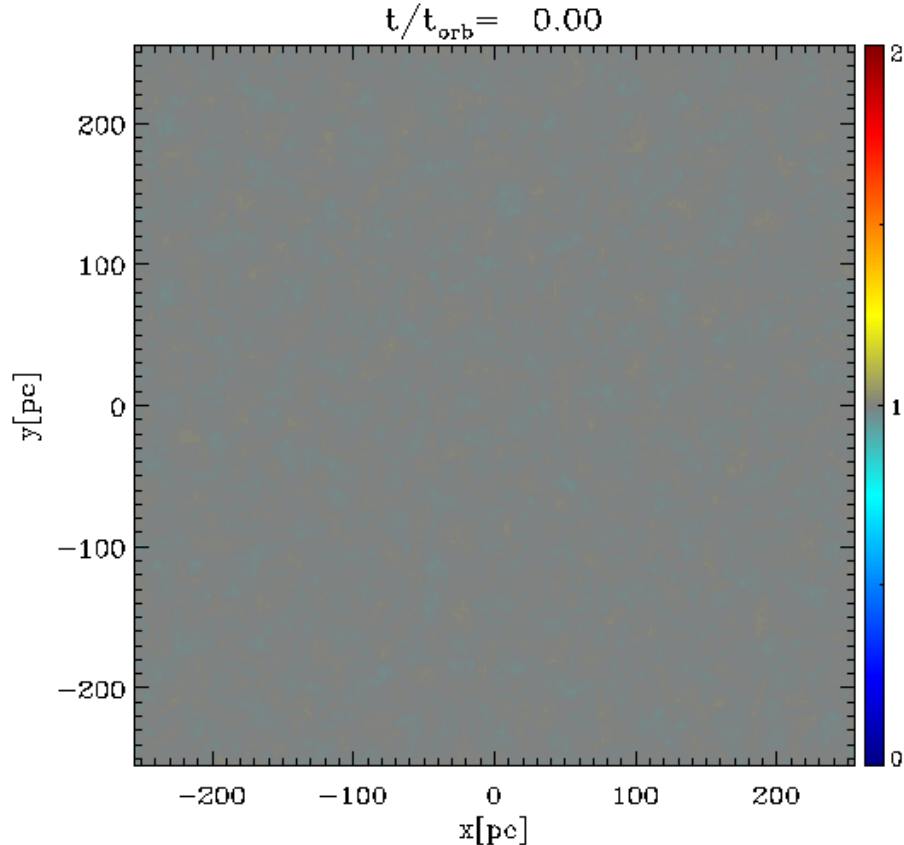
- Self-regulation via turbulent driving when $t_{\text{ff}} \gg t_{\text{dyn}}$:

$$\dot{M}_* \sim \frac{GM_{\text{tot}}M}{L^2 p_*/m_*} \longrightarrow \Sigma_{\text{SFR}} \sim \frac{G\Sigma_{\text{gas}}\rho_{\text{star}}H_{\text{gas}}}{p_*/m_*}$$

$$H_{\text{gas}} \sim v(G\rho_{\text{star}})^{-1/2} \quad \Sigma_{\text{SFR}} \sim \frac{\Sigma(G\rho_{\text{star}})^{1/2}v}{p_*/m_*}$$

Simulations with turbulent feedback and radiative heating and for outer-disk regime

- Kim, Kim, & Ostriker (2011); Kim, Ostriker, & Kim (2013)
 - include turbulent driving from SN (momentum injection)
 - include dependence of heating rate on star formation rate ($\Gamma \propto J_{\text{FUV}} \propto \Sigma_{\text{SFR}}$)



Self-regulation in externally-confined disks

Allowing for thermal as well as turbulent feedback to atomic gas, $P_{\text{th}} = \eta_{\text{th}} \Sigma_{\text{SFR}}$ and $P_{\text{turb}} = \eta_{\text{turb}} \Sigma_{\text{SFR}}$, with

$$P_{\text{th}} + P_{\text{turb}} = (\eta_{\text{th}} + \eta_{\text{turb}}) \Sigma_{\text{SFR}} = P_{\text{DE}} \text{ for}$$

$$P_{\text{DE}} = \frac{\Sigma}{2} g_z \approx \frac{\pi G \Sigma^2}{2} + \Sigma (2G\rho_*)^{1/2} \sigma_z$$

depending only on the gravity and total gas surface density of the disk form vertical dynamical equilibrium

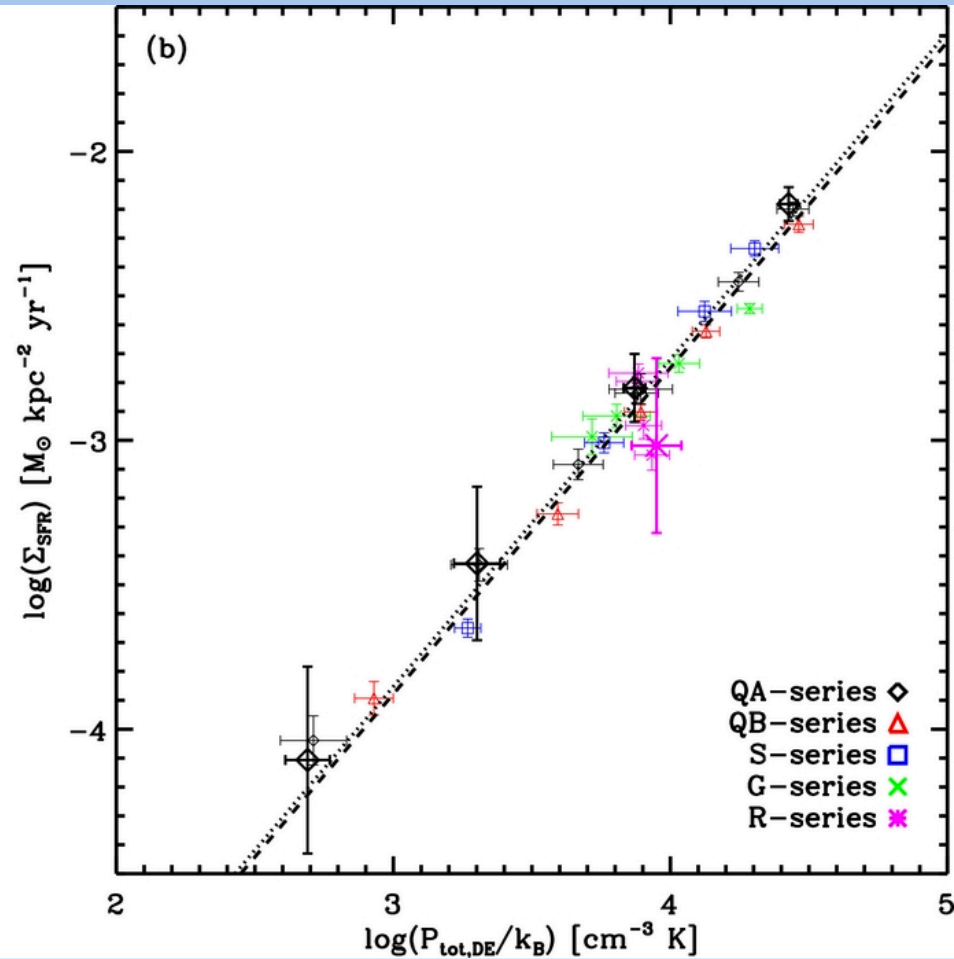
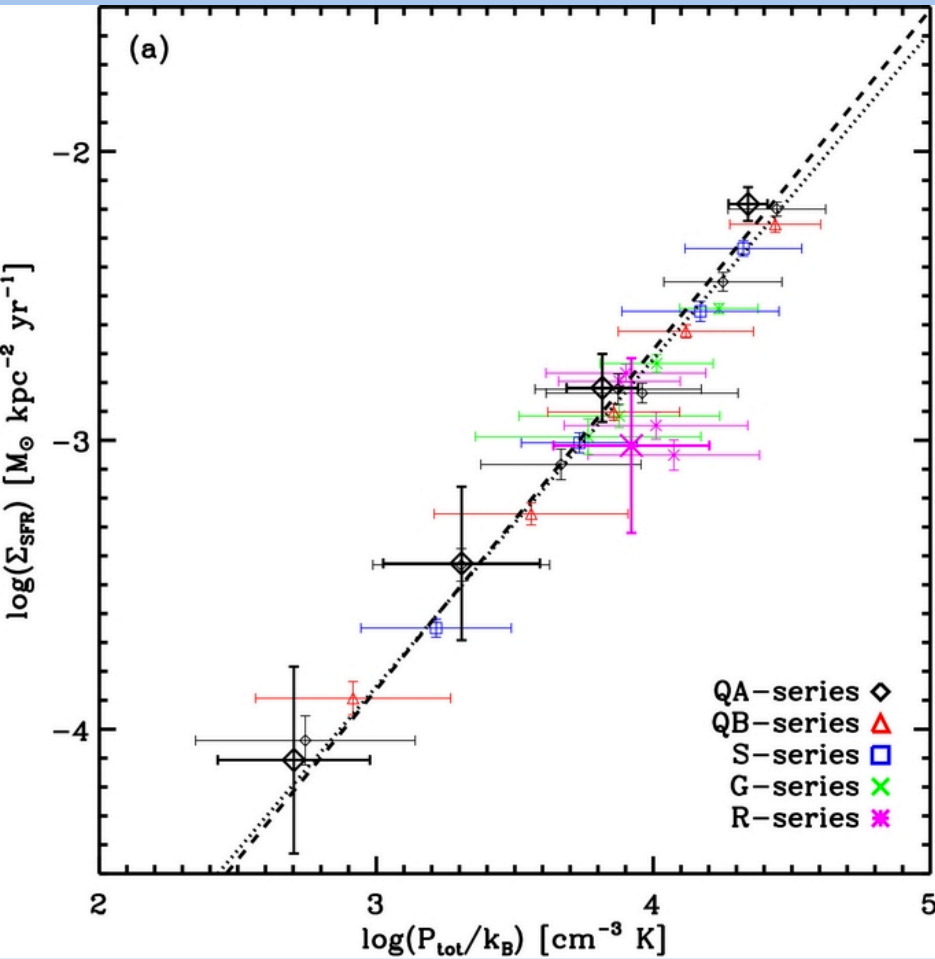
- General result is

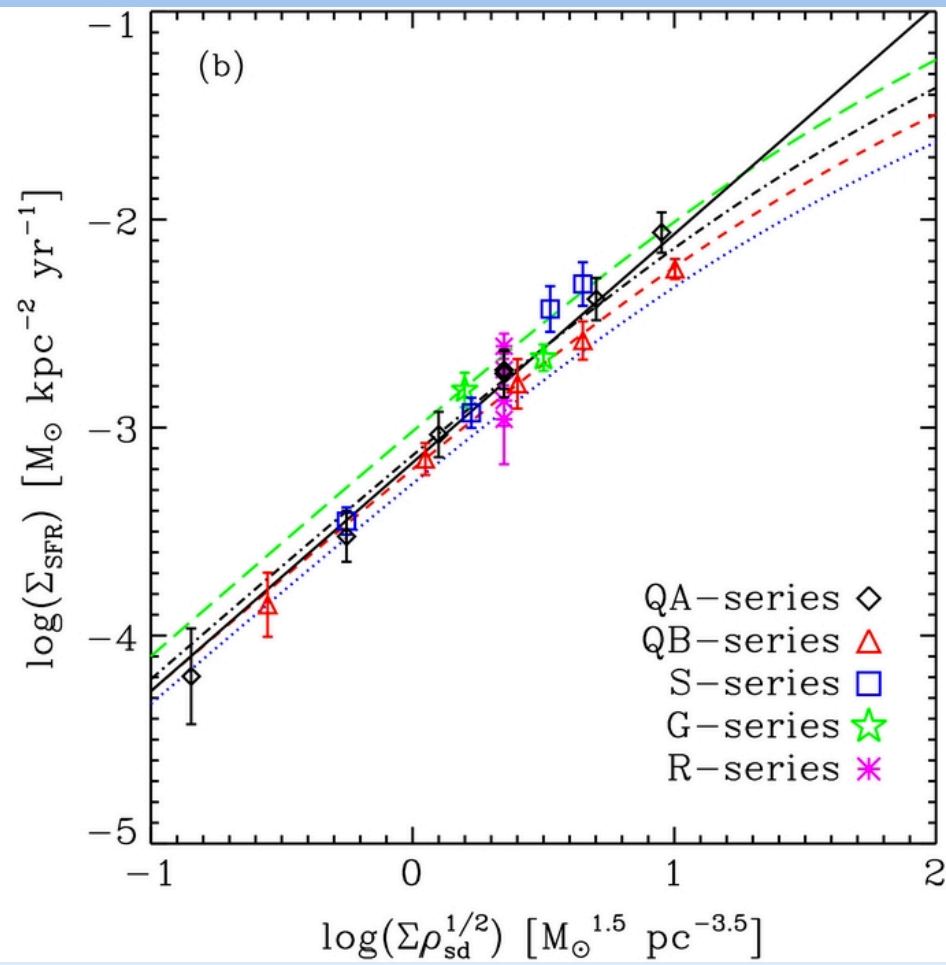
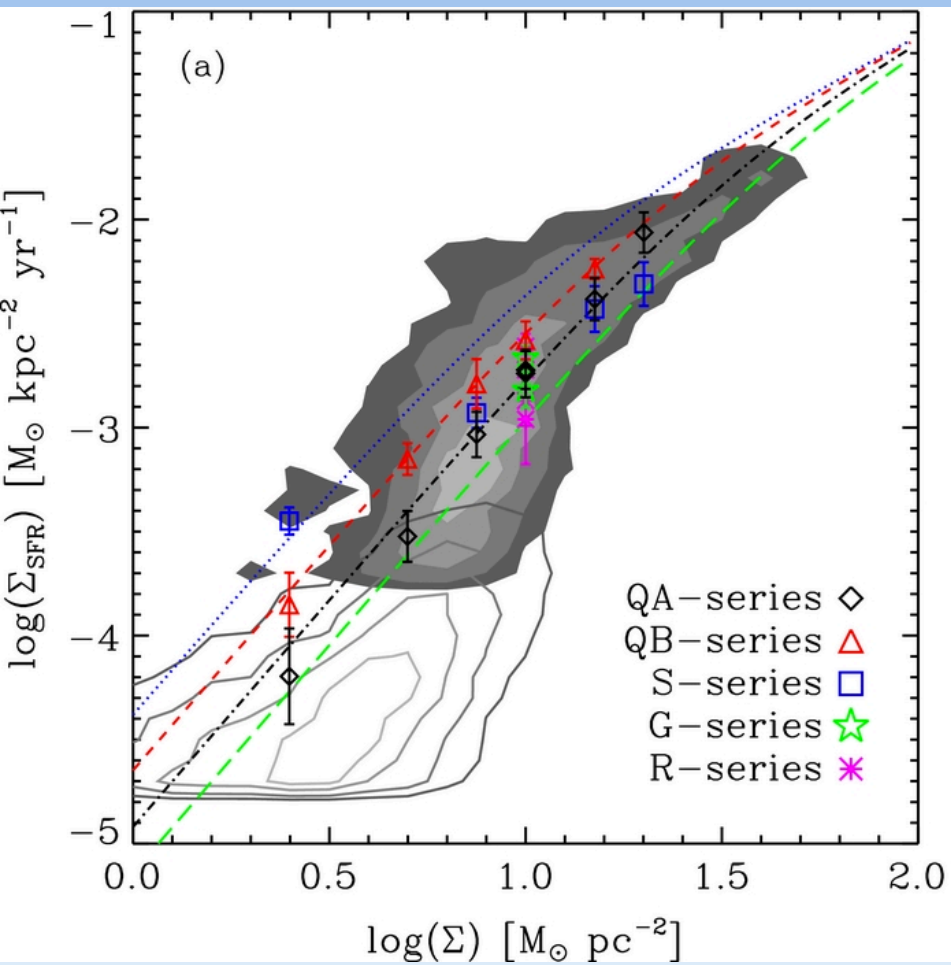
$$\Sigma_{\text{SFR}} = 2 \times 10^{-3} M_{\odot} \text{ kpc}^{-2} \text{ yr}^{-1} \left(\frac{P/k}{10^4 \text{ cm}^{-3} \text{ K}} \right)$$

and for outer disk regions:

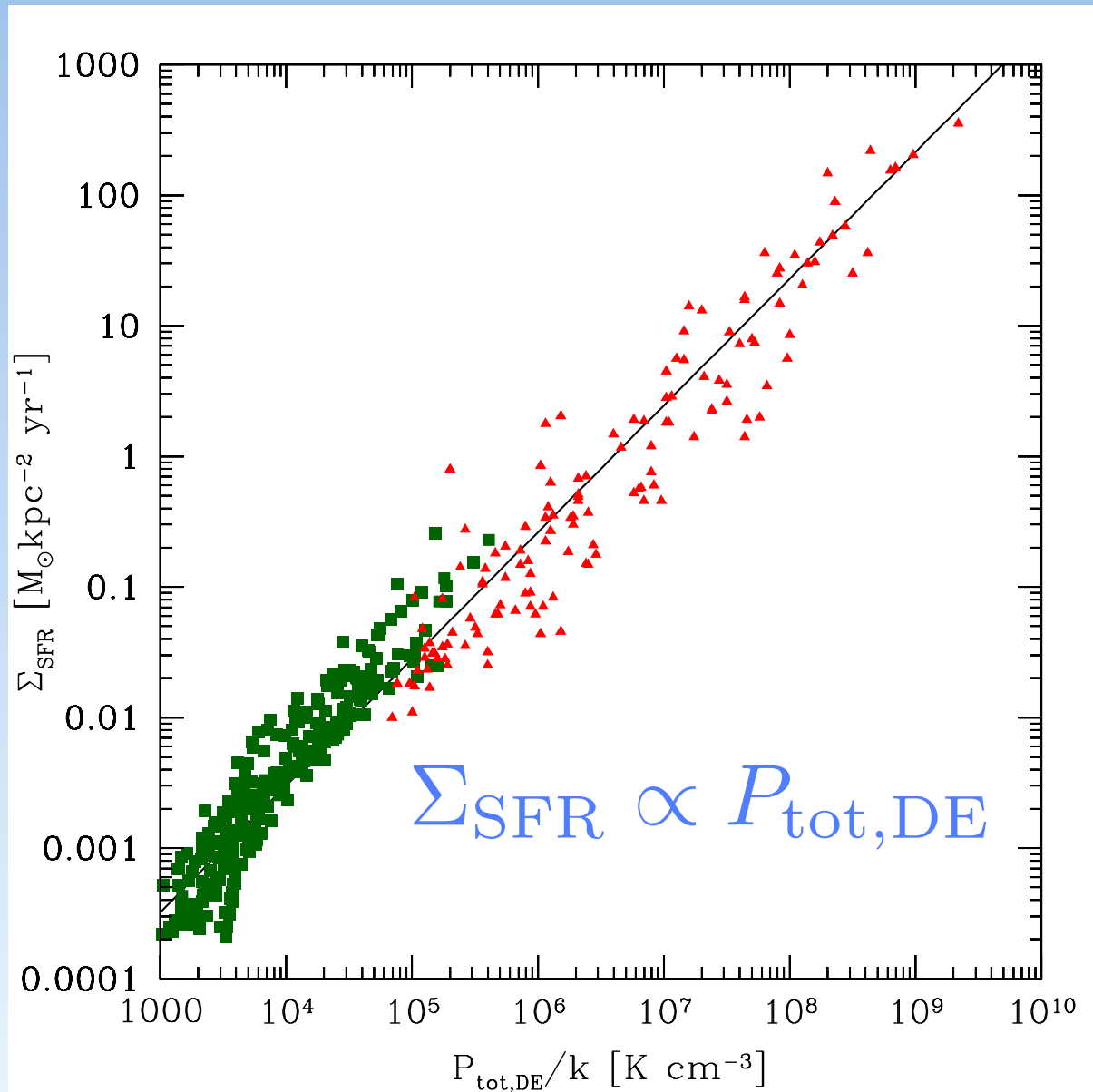
$$\Sigma_{\text{SFR}} = 2 \times 10^{-3} M_{\odot} \text{ kpc}^{-2} \text{ yr}^{-1} \left(\frac{\Sigma}{10 M_{\odot} \text{ pc}^{-2}} \right) \left(\frac{\rho_*}{0.1 M_{\odot} \text{ pc}^{-3}} \right)^{1/2}$$

$$P_{DE} = \frac{\Sigma}{2} g_z \approx \frac{\pi G \Sigma^2}{2} + \Sigma (2G \rho_*)^{1/2} \sigma_z$$

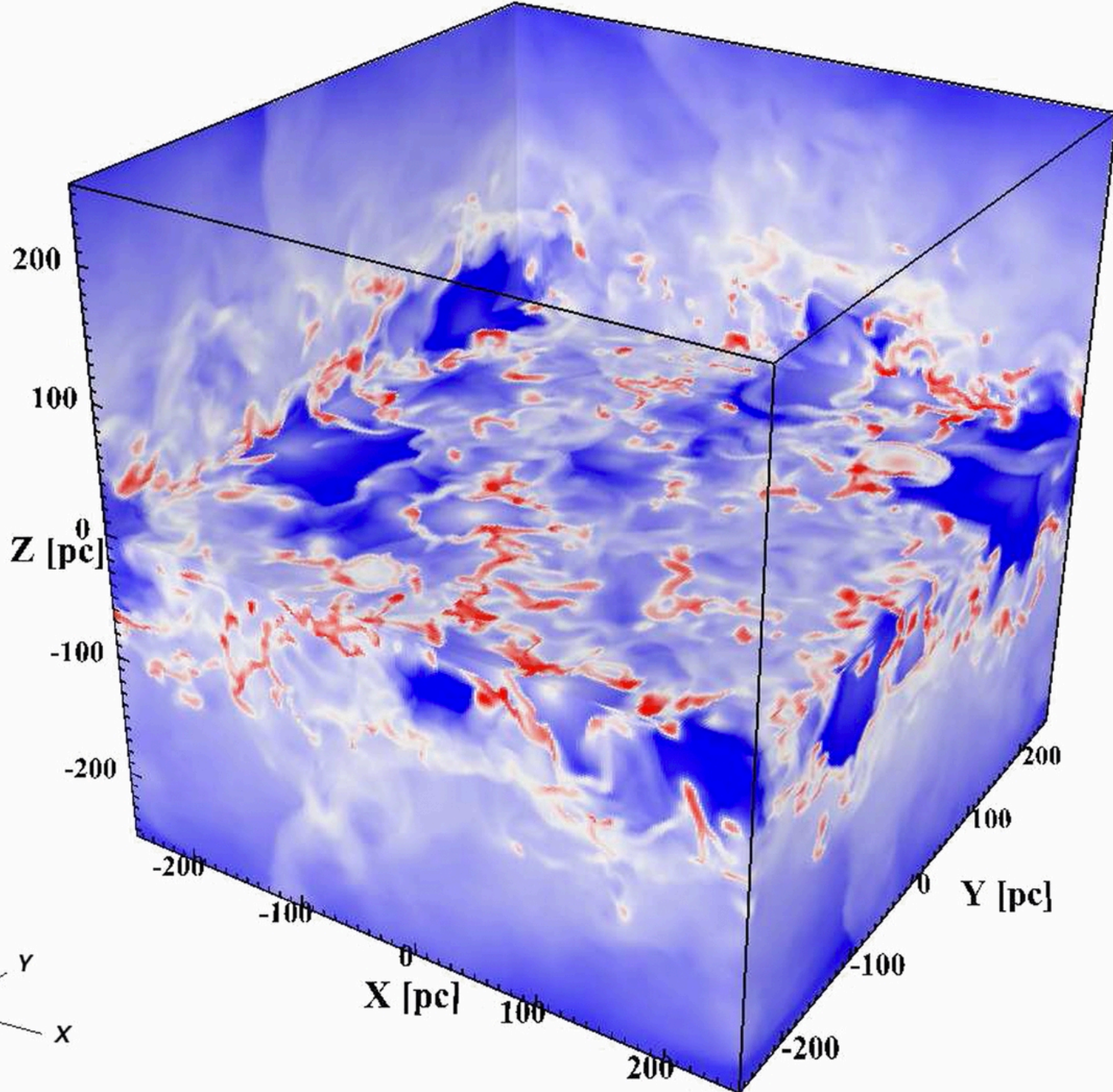


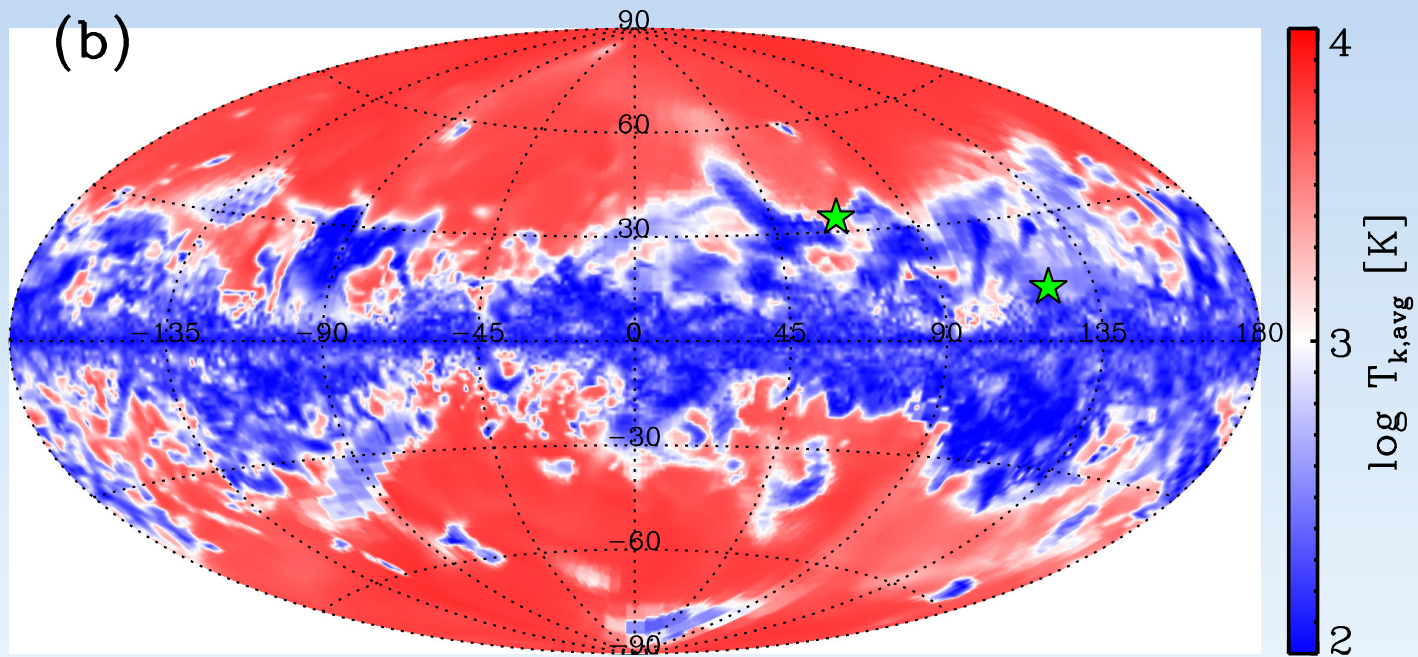
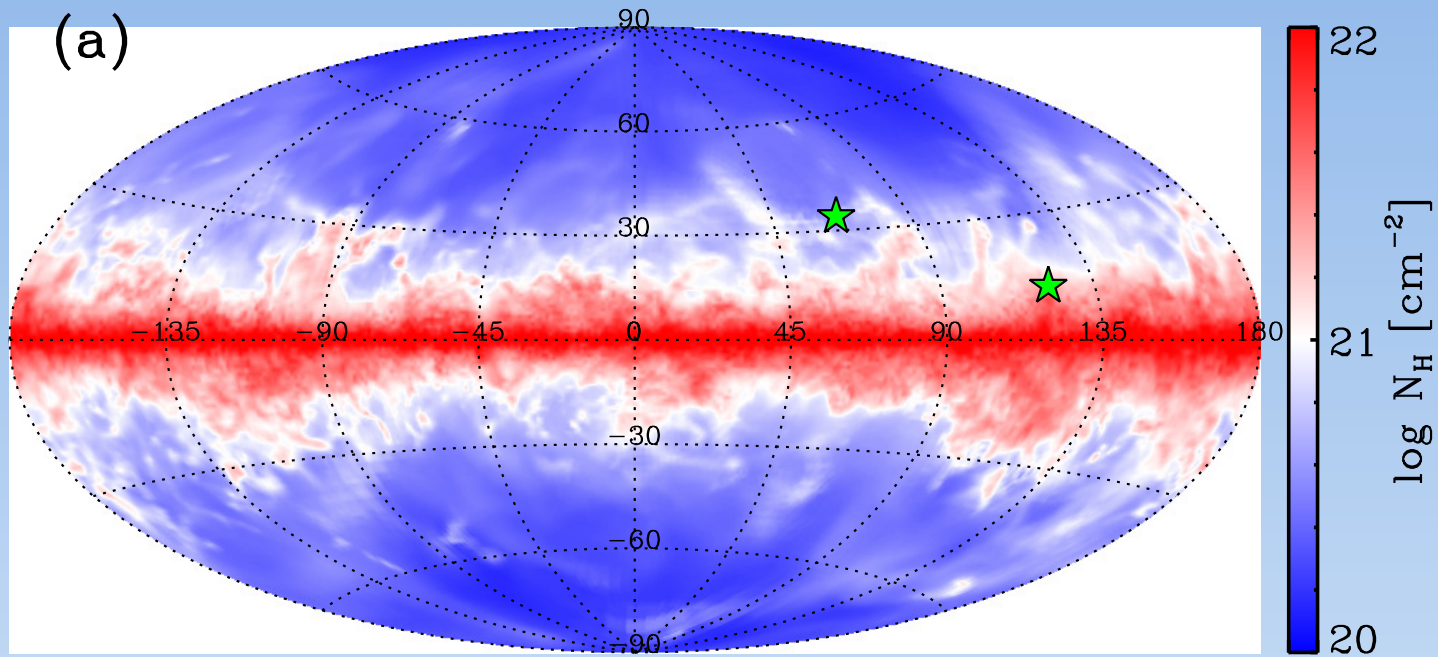


Σ_{SFR} vs. equilibrium pressure



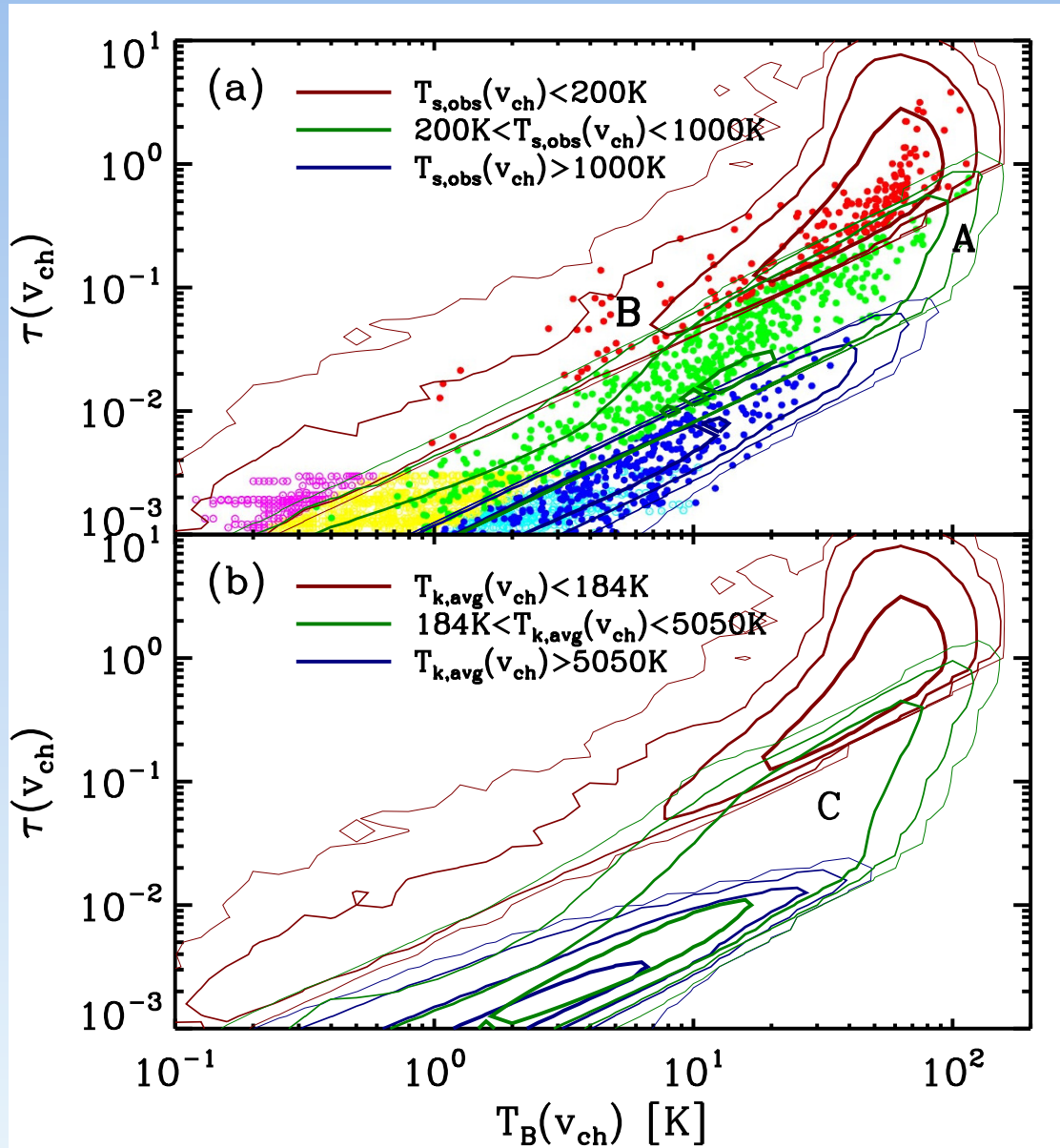
$\log(n)$





21 cm T_B , T_S , τ

Kim, Ostriker, Kim (2014)



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Radiation force

- Direct radiation

$$\dot{p} = \frac{\mathcal{L}_*}{c} \rightarrow \frac{\varepsilon_{nuc} c^2 \dot{M}_*}{c} \Rightarrow \frac{p_*}{m_*} = \frac{\dot{p}}{\dot{M}_*} = \varepsilon_{nuc} c \sim 180 \text{ km/s}$$

- Reprocessed radiation

Overall ISM: $\frac{p_*}{m_*} = \varepsilon_{nuc} c \kappa_{IR} \Sigma \sim 180 \text{ km/s} \times \tau_{IR}$

$$\tau_{IR} = 2 \kappa_{IR} \frac{\Sigma}{10^4 M_\odot \text{ pc}^{-2}}$$

Individual cluster-forming cloud: $\dot{p} \sim \frac{\mathcal{L}_* \tau_{IR}}{c} \rightarrow \frac{\Psi M_* \tau_{IR}}{c}$

$$\frac{p_*}{m_*} \sim \frac{\Psi \tau_{IR} t_{embed}}{c} \sim 20 \text{ km/s} \frac{t_{embed}}{\text{Myr}} \tau_{IR}$$

or $\frac{p_*}{m_*} \sim v \frac{\Psi \kappa_{IR} t_{embed}}{Gc} \frac{t_{embed}}{t_{dyn}} \sim v$

Turbulent, cluster-forming cloud with IR radiation

- Starting with $\epsilon_{ff} \sim \frac{v}{p_*/m_*}$, efficiency over cloud lifetime is

$$\epsilon \sim \epsilon_{ff} \frac{t_{life}}{t_{ff}} \sim \frac{v}{L} \frac{v}{p_*/m_*} t_{life} \sim \frac{v^2 M_*}{L \dot{p}} \sim \frac{GM}{L^2} \frac{M_*}{\dot{p}} \sim \frac{G \Sigma M_*}{\dot{p}}$$

- Momentum input rate from reprocessed IR is

$$\dot{p} \sim \frac{\mathcal{L}_* \tau}{c} \sim \frac{M_* \Psi \kappa_{IR} \Sigma}{c}$$

$$\Rightarrow \boxed{\epsilon \sim \frac{Gc}{\Psi \kappa_{IR}}}$$

cf. Murray et al (2010)

NB: for cluster with radiation-driven shell, exact result is:

$$\epsilon_{min} = \left[\frac{\Psi \kappa_{IR}}{2\pi Gc} - 1 \right]^{-1}$$

Ostriker & Shetty (2011) 23

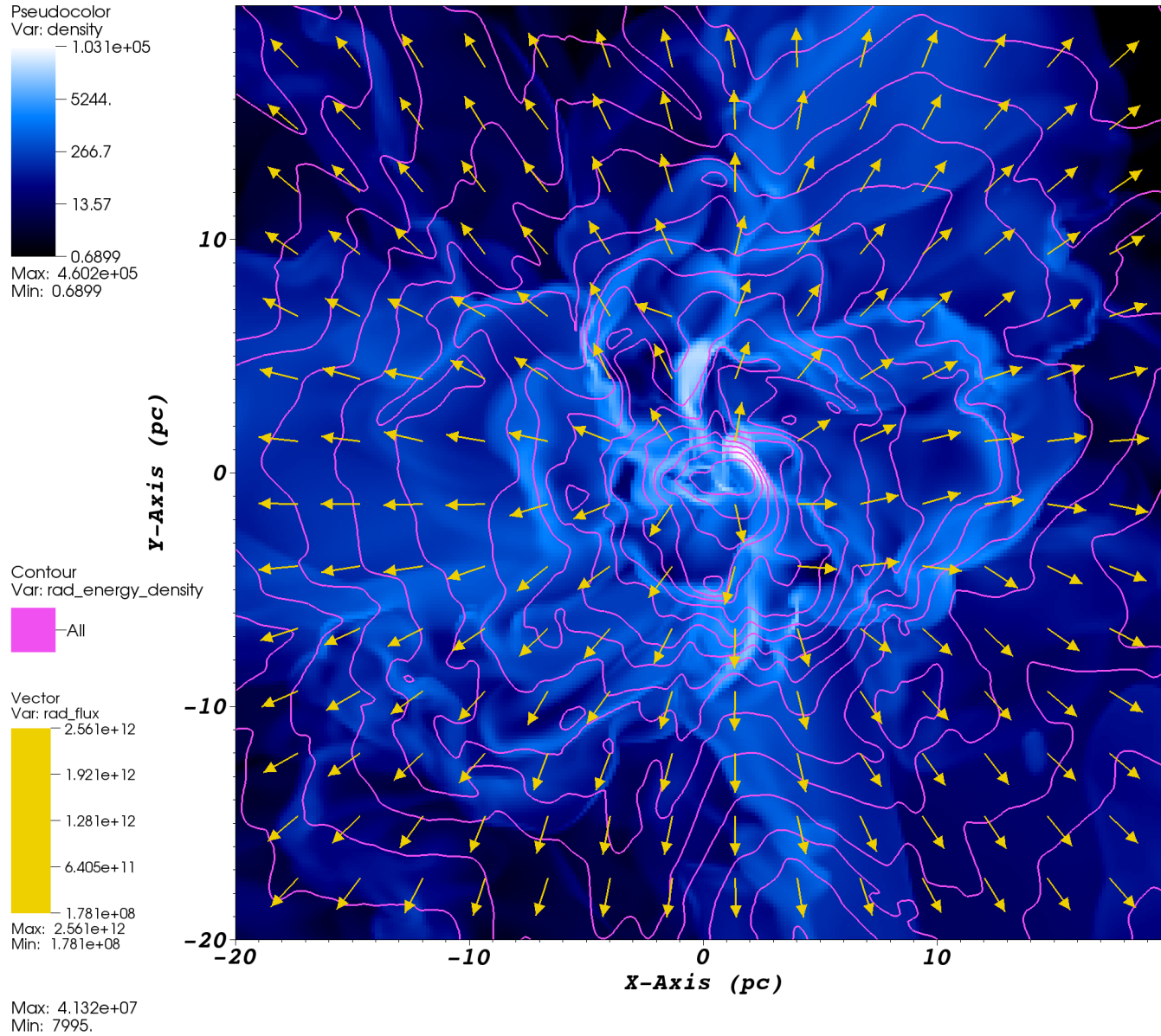
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Cluster-forming cloud with IR radiation: numerical simulation

Skinner & Ostriker (2014)

- Consider evolution of a turbulent, star-forming GMC, including effects of reprocessed radiation
- Solve evolution equations for first two radiation moments (energy and flux) using RSL method with M1 closure (Skinner & Ostriker 2013), combined with gas integration of *Athena*
- Sink particles (Gong & Ostriker 2013) model star (sub) clusters with luminosity $\mathcal{L}_* = \Psi M_*$
- Consider a range of κ , initial cloud mass M_{GMC} and radius R_{GMC}

$R=10$ pc, $M=1e6$ Msun, $\kappa=20$ cm² g⁻¹, $N=256$, $t/t_{ff}=3$



Cluster-forming cloud with IR radiation: numerical simulation

Skinner & Ostriker (2014)

- Measure gas mass converted to stars vs ejected
 - efficiency $\varepsilon_* = M_*/M_{\text{GMC}}$ and $\varepsilon_{\text{wind}} = M_{\text{ejected}}/M_{\text{GMC}}$
- Measure gas momentum ejected p_{ej}
 - compute $p_*/m_* = p_{\text{ej}}/M_*$ compared to $v = (GM/R)^{1/2}$
- Explore gas and radiation structure in cloud:

$$f_{\text{Edd}}(r) = \frac{\langle F_r \rho \kappa / c \rangle}{\langle g_r \rho \rangle}; \quad f_{\text{Edd}} = \frac{\int \langle F_r \rho \kappa / c \rangle r^2 dr}{\int \langle g_r \rho \rangle r^2 dr} \quad \textit{versus} \quad \frac{\Psi \kappa}{4\pi G c}$$

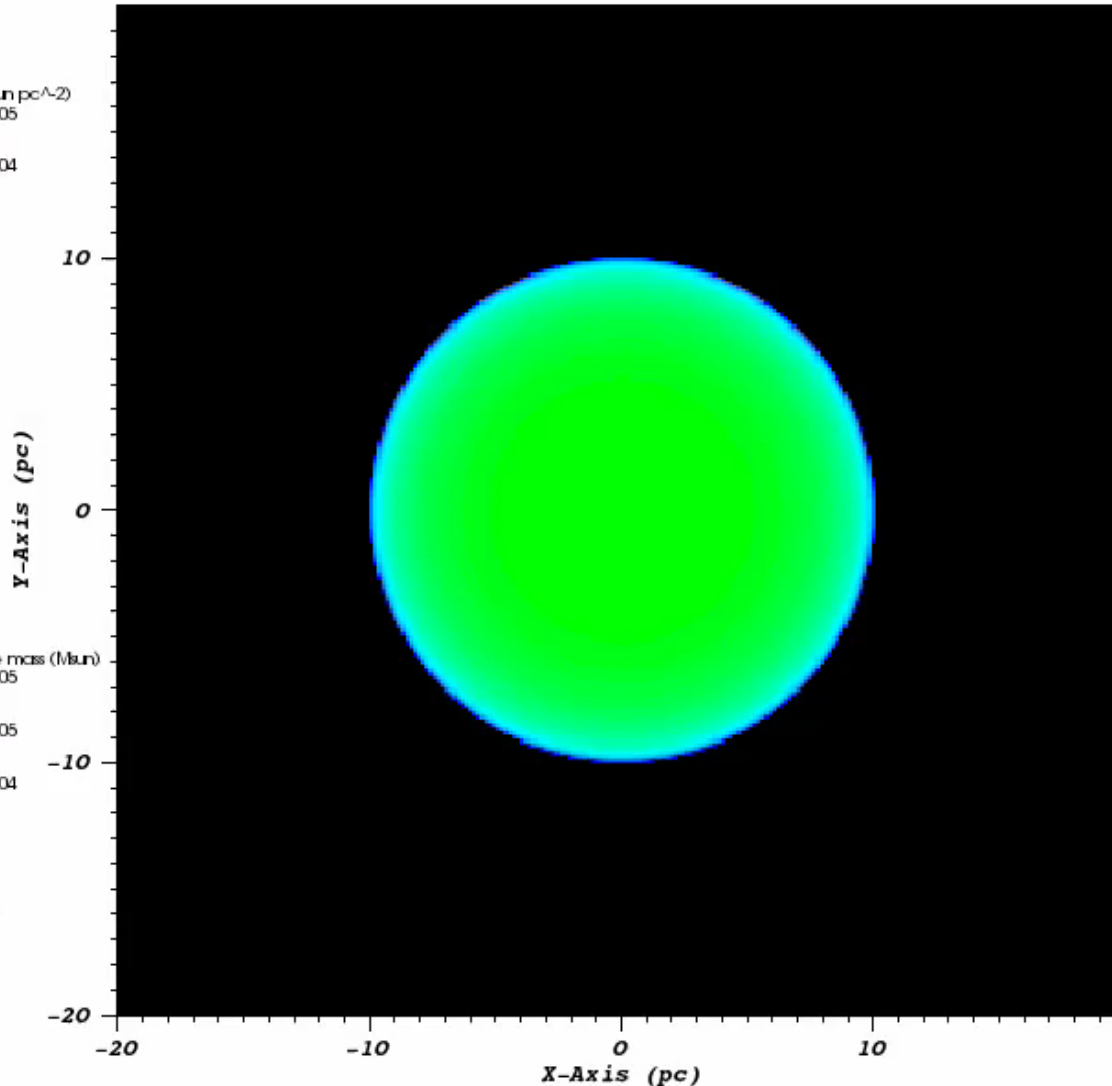
$$f_{\text{trap}} = \frac{\int \langle F_r \rho \kappa / c \rangle 4\pi r^2 dr}{\mathcal{L}/c} \quad \textit{versus} \quad \tau = \int \rho \kappa dr$$

cf. Krumholz & Thompson; Davis et al

Star-forming cloud with RHD

R=10 pc, M=1e6 Msun, kappa=10 cm² g⁻¹, N=256

Pseudocolor
Var: Sigma3 (Msun pc⁻²)
2.000e+05
2.991e+04
4472.
668.7
100.0
Max: 4822.
Min: 95.49



Pseudocolor
Var: star particle mass (Msun)
5.000e+05
1.057e+05
2.236e+04
4729.
1000.
Max: 5.000e+05
Min: 1000.

t/t_{ff} = 0

10⁶ M_⊙ initial
cloud with
sink particles
and RHD

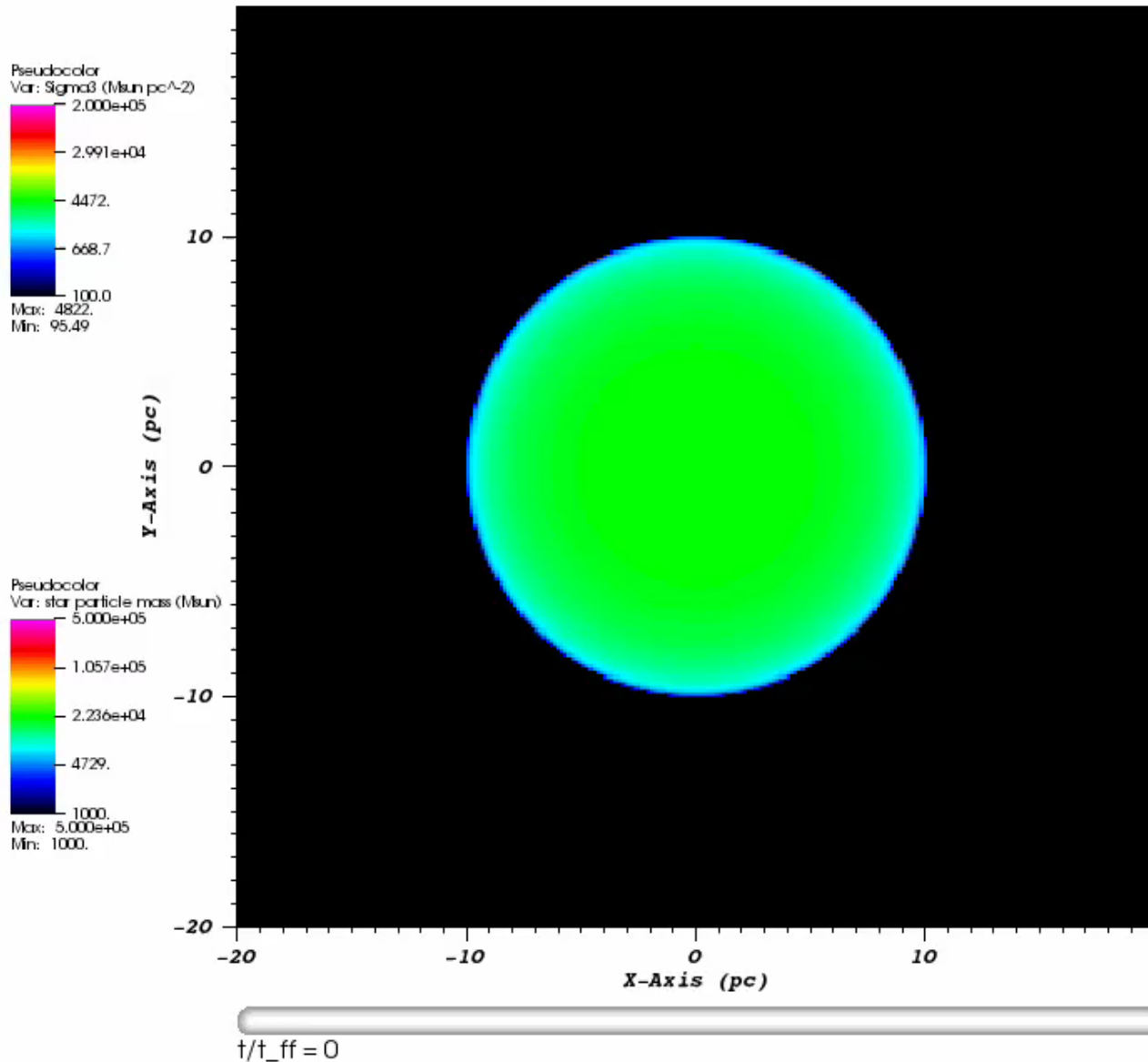
$\kappa=10 \text{ g/cm}^2$

$L_* = \Psi M_*$ for
subclusters;
 $\Psi=1700 \text{ erg/s/g}$

$t_{\text{ff}}=0.52 \text{ Myr}$

Star-forming cloud with RHD

$R=10 \text{ pc}$, $M=1e6 \text{ Msun}$, $\kappa=20 \text{ cm}^2 \text{ g}^{-1}$, $N=256$



$10^6 M_{\odot}$ initial cloud with sink particles and RHD

$\kappa=20 \text{ g/cm}^2$

$L_* = \Psi M_*$ for subclusters;
 $\Psi=1700 \text{ erg/s/g}$

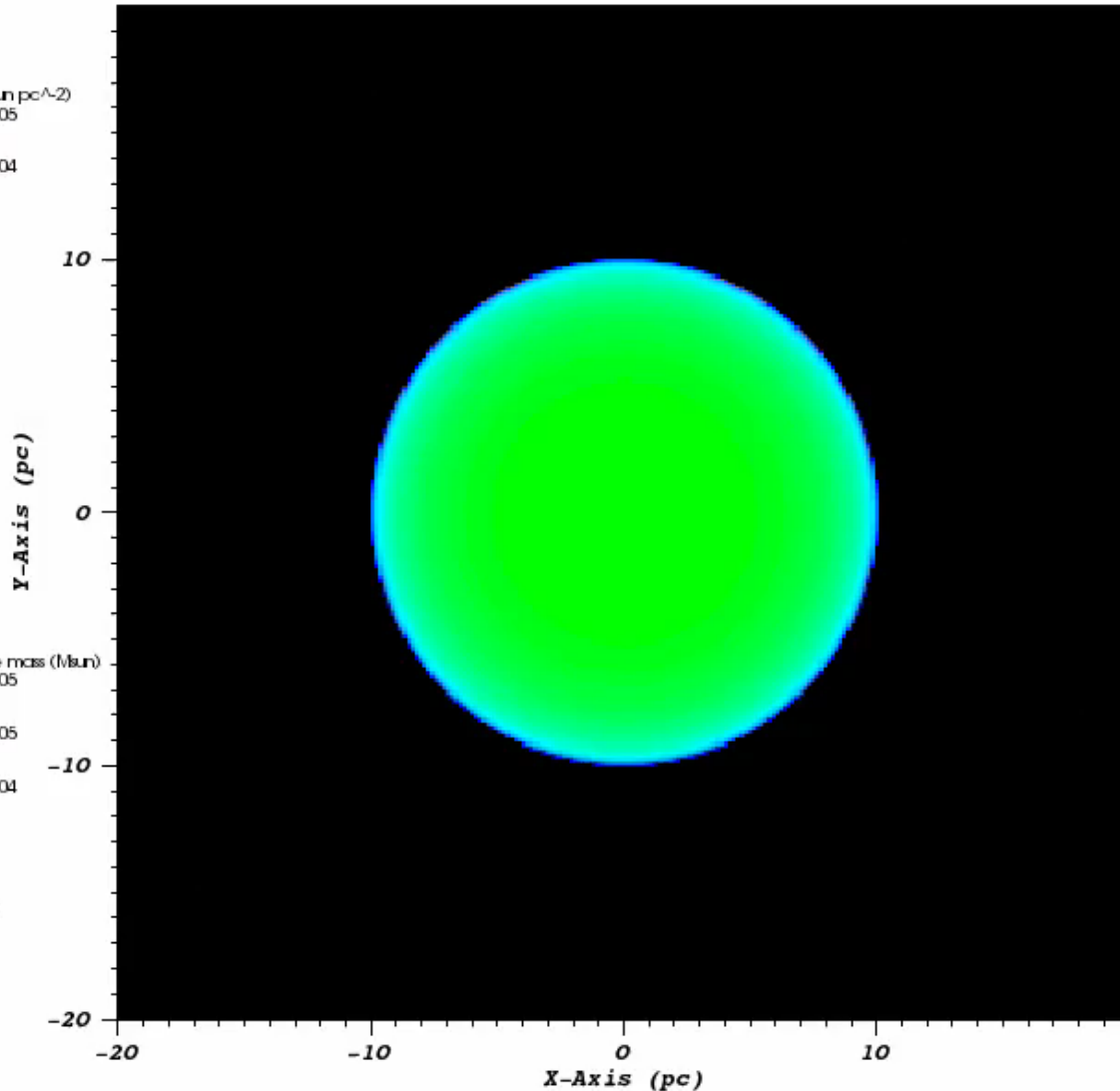
$t_{\text{ff}}=0.52 \text{ Myr}$

Skinner & Ostriker (2014)

Star-forming cloud with RHD

$R=10$ pc, $M=1e6$ Msun, $\kappa=40$ cm² g⁻¹, $N=256$

Pseudocolor
Var: Σ_{star} (Msun pc⁻²)
2.000e+05
2.991e+04
4472.
668.7
100.0
Max: 4822.
Min: 95.49



Pseudocolor
Var: star particle mass (Msun)
5.000e+05
1.057e+05
2.236e+04
4729.
1000.
Max: 5.000e+05
Min: 1000.

$10^6 M_{\odot}$ initial cloud with sink particles and RHD

$\kappa=40$ g/cm²

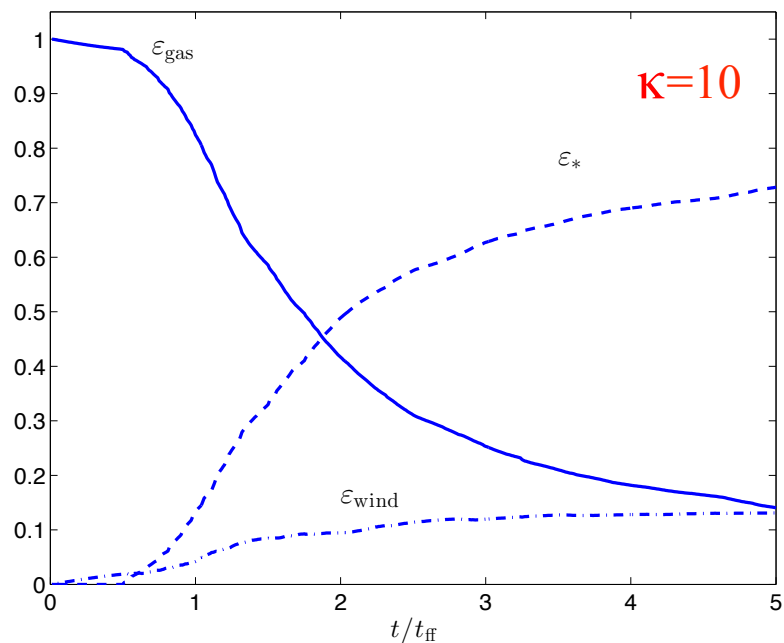
$L_* = \Psi M_*$ for subclusters;
 $\Psi=1700$ erg/s/g

$t_{\text{ff}}=0.52$ Myr

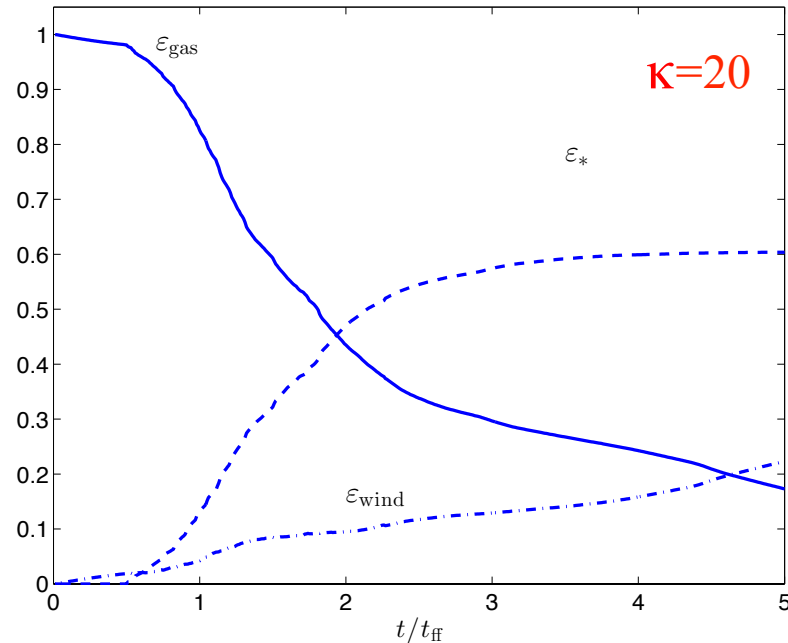
Skinner & Ostriker (2014)

$t/t_{\text{ff}} = 0$

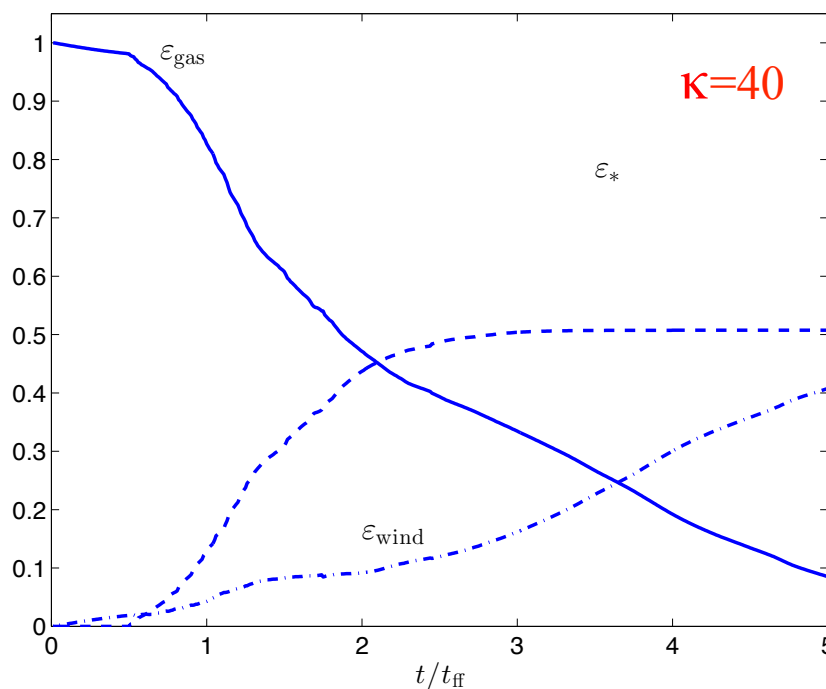
$R_{\text{GMC}} = 10.00 \text{ pc}$, $M_{\text{GMC}} = 1.00e + 06M_{\odot}$, $\kappa_{\text{IR}} = 10 \text{ cm}^2 \text{ g}^{-1}$, $N=256$



$R_{\text{GMC}} = 10.00 \text{ pc}$, $M_{\text{GMC}} = 1.00e + 06M_{\odot}$, $\kappa_{\text{IR}} = 20 \text{ cm}^2 \text{ g}^{-1}$, $N=256$



$R_{\text{GMC}} = 10.00 \text{ pc}$, $M_{\text{GMC}} = 1.00e + 06M_{\odot}$, $\kappa_{\text{IR}} = 40 \text{ cm}^2 \text{ g}^{-1}$, $N=256$

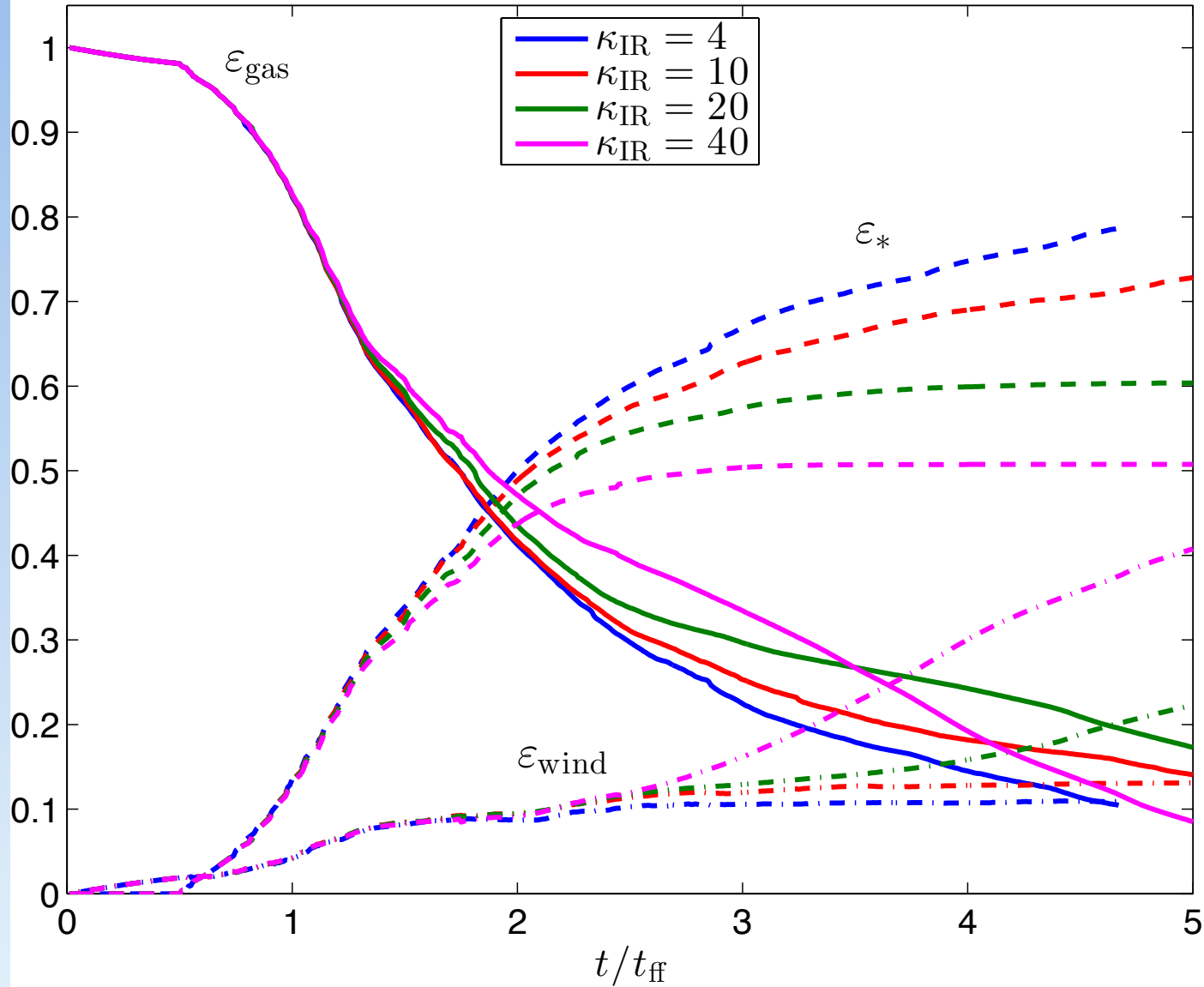


Skinner & Ostriker (2014)

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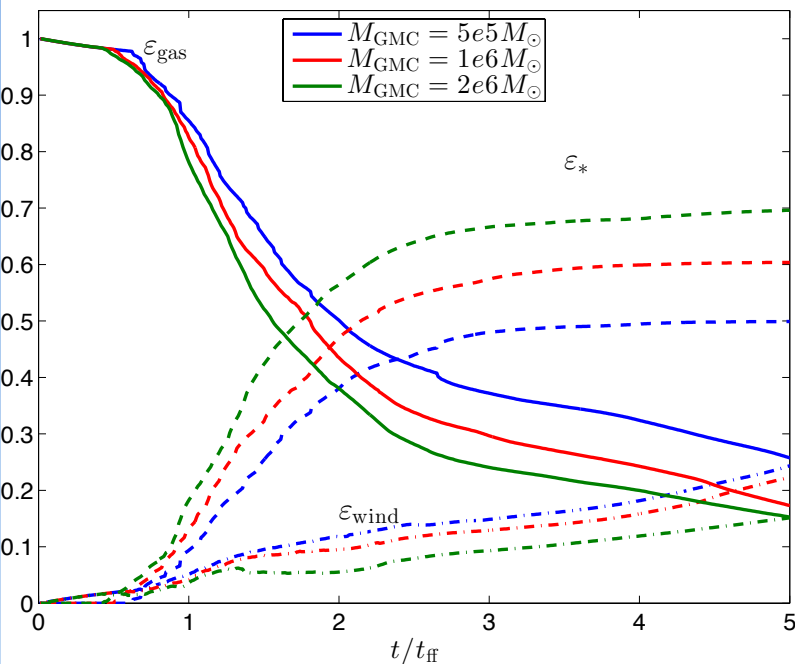
$t_{\text{ff}} = 0.52 \text{ Myr}$

$R_{\text{GMC}} = 10 \text{ pc}, M_{\text{GMC}} = 1e6 M_{\odot}, N = 256$

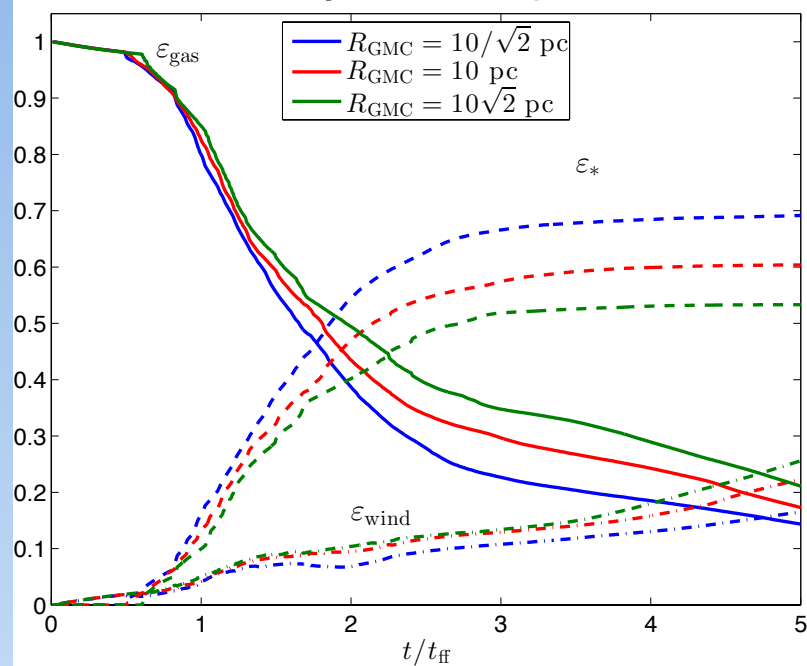


Skinner & Ostriker (2014)

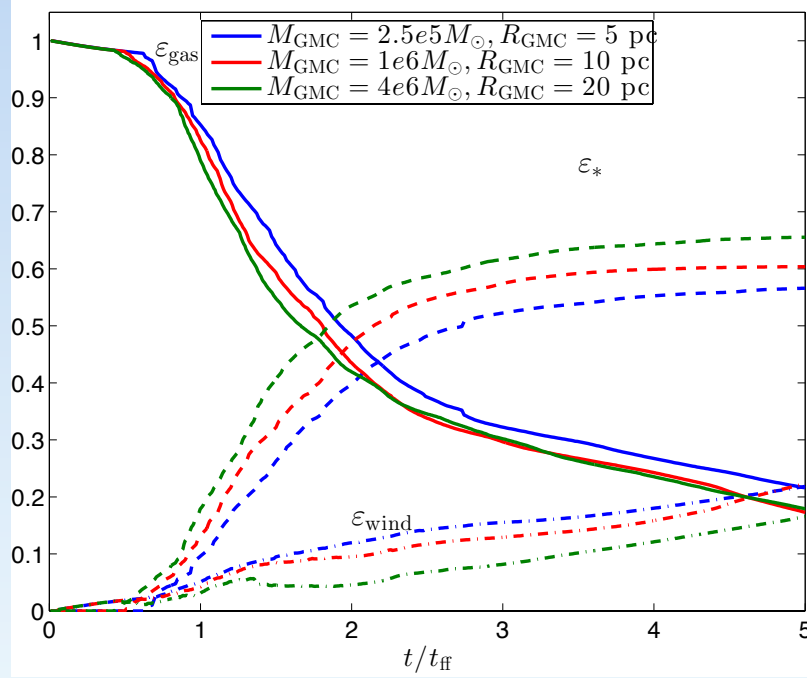
$R_{\text{GMC}} = 10 \text{ pc}$, $\kappa_{\text{IR}} = 20 \text{ cm}^2 \text{ g}^{-1}$, $N = 256$



$M_{\text{GMC}} = 1e6 M_{\odot}$, $\kappa_{\text{IR}} = 20 \text{ cm}^2 \text{ g}^{-1}$, $N = 256$



$\kappa_{\text{IR}} = 20 \text{ cm}^2 \text{ g}^{-1}$, $N = 256$



Fractional mass loss and net SF efficiency relatively insensitive to cloud mass and size

Skinner & Ostriker (2014)

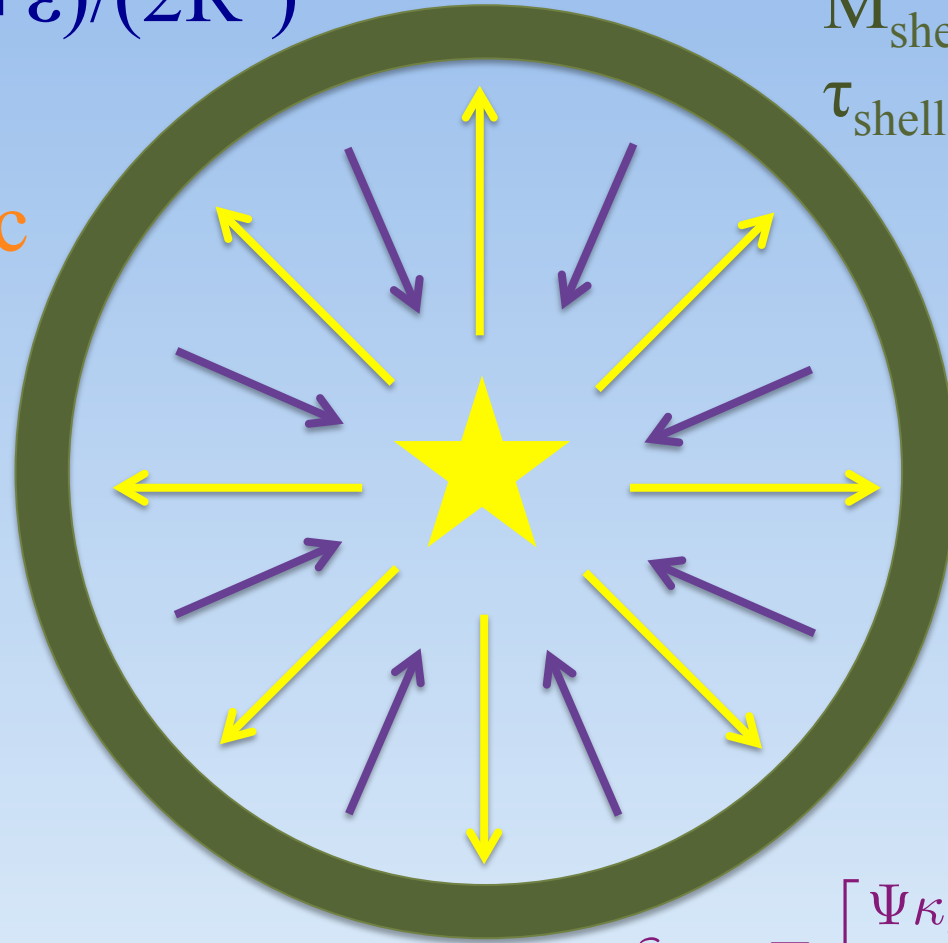
$t_{\text{ff}} = 0.52 \text{ Myr}$

$$F_{\text{grav}} = GM(1+\varepsilon)/(2R^2)$$

$$M_{\text{shell}} = (1-\varepsilon) M$$

$$\tau_{\text{shell}} = \kappa M_{\text{shell}} / (4\pi R^2)$$

$$F_{\text{rad}} = \mathcal{L}_* \tau_{\text{shell}} / c$$



$$M_* = \varepsilon M$$

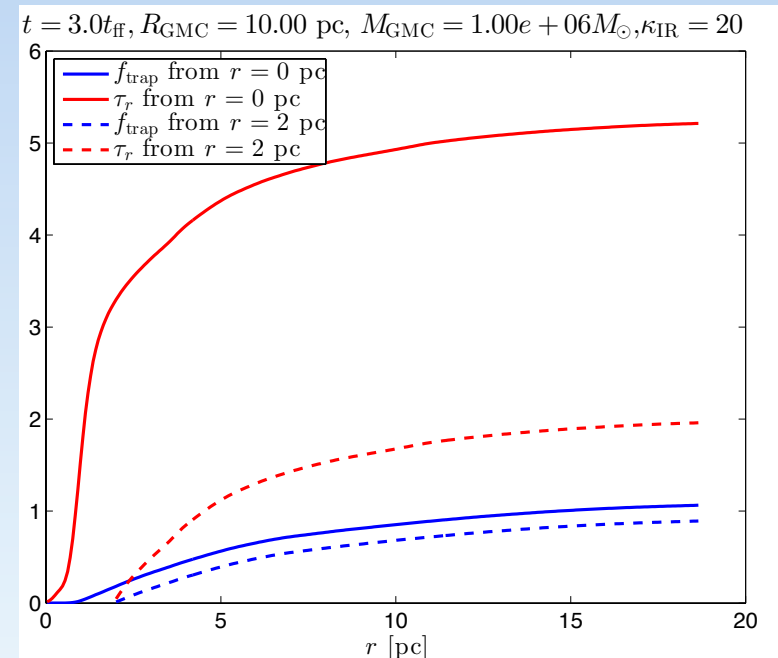
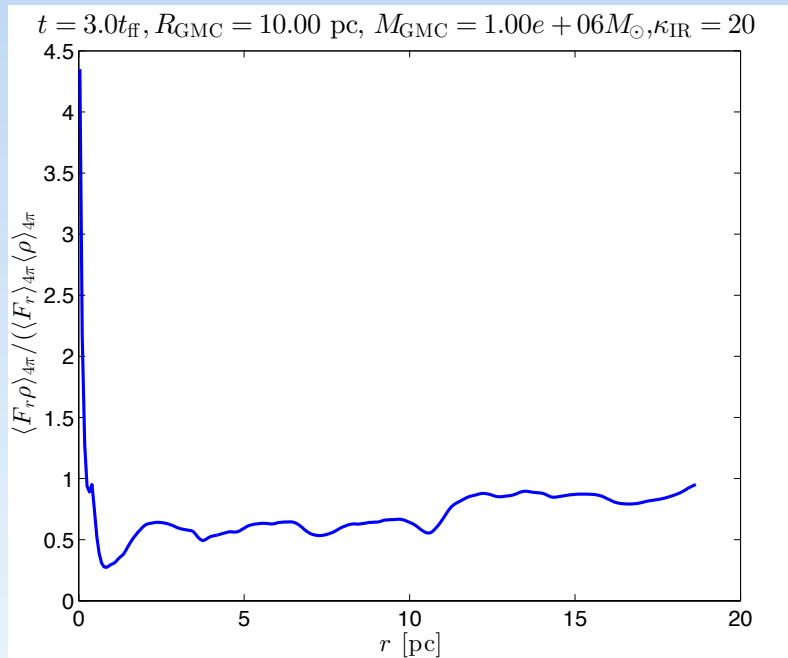
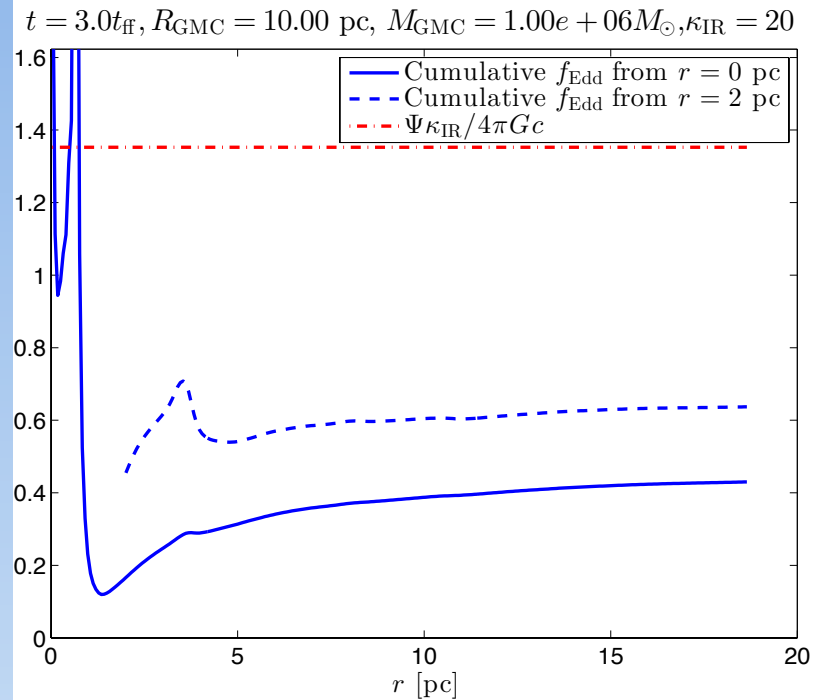
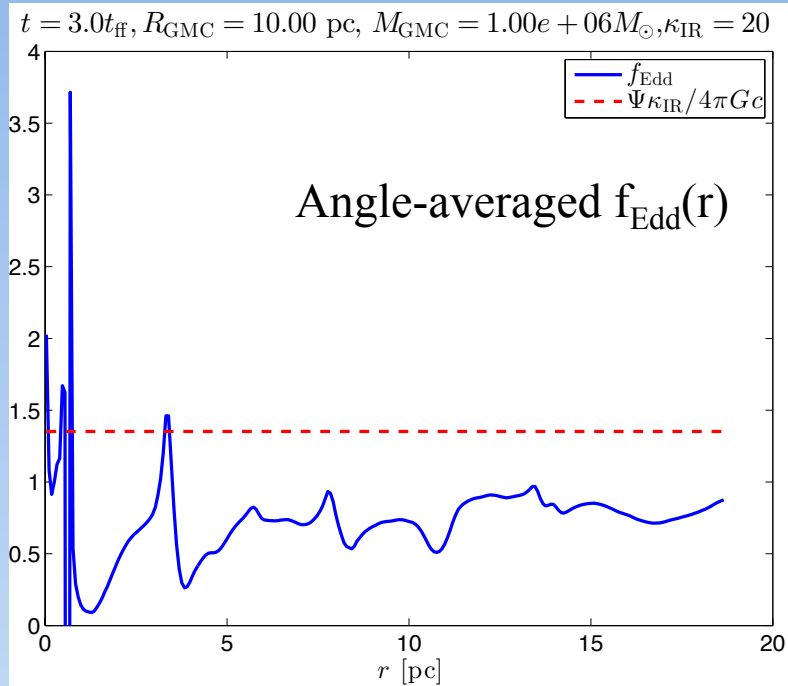
$$\mathcal{L}_* = \Psi M_*$$

$$\varepsilon_{\text{min}} = \left[\frac{\Psi \kappa_{IR}}{2\pi G c} - 1 \right]^{-1}$$

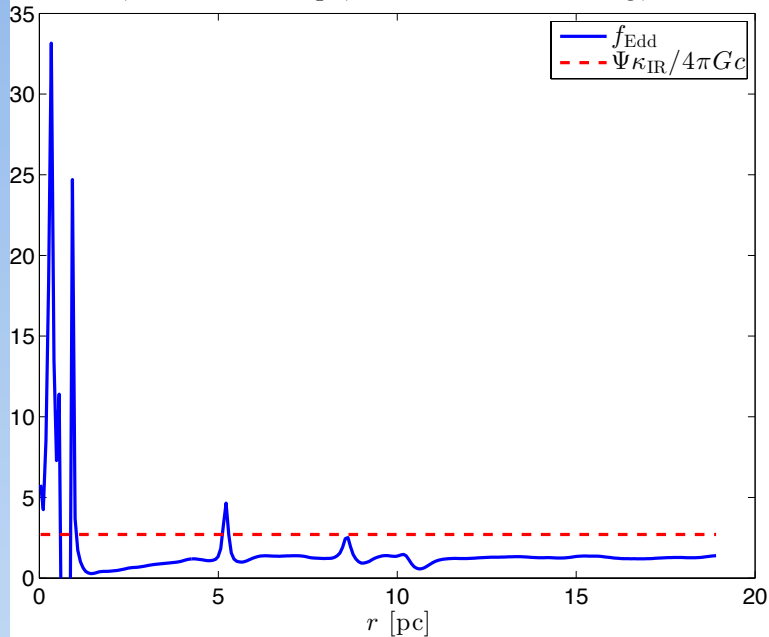
→ 2 for $\kappa = 10$ (*BOUND*)

→ 0.5 for $\kappa = 20$

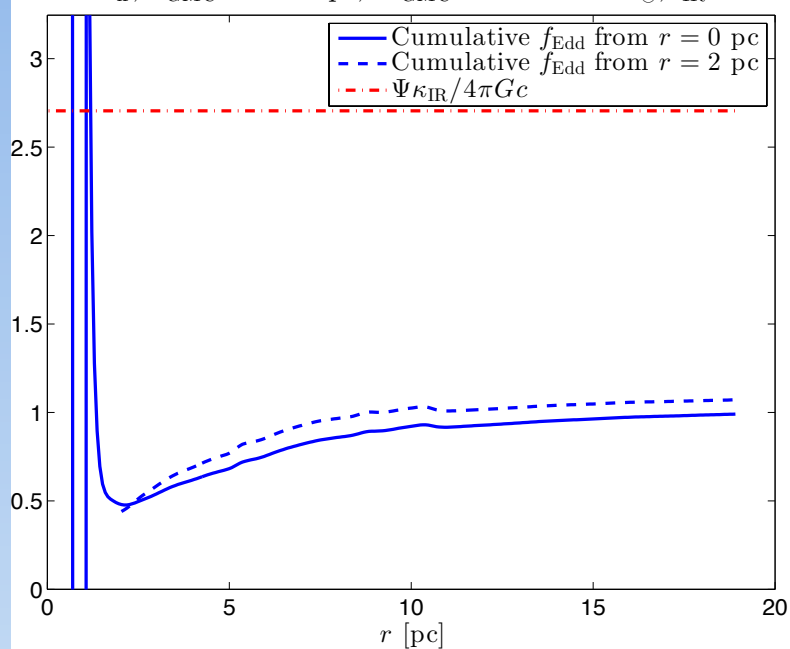
→ 0.2 for $\kappa = 40$



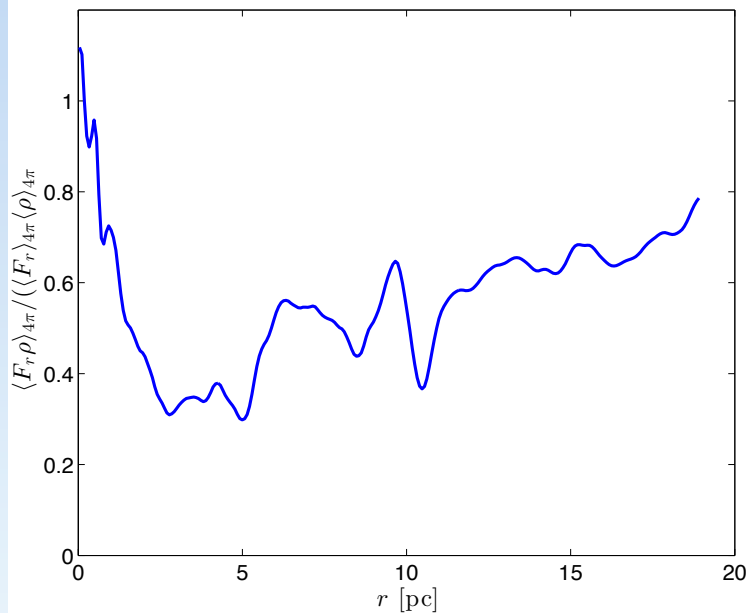
$t = 3.0t_{\text{ff}}, R_{\text{GMC}} = 10.00 \text{ pc}, M_{\text{GMC}} = 1.00e + 06M_{\odot}, \kappa_{\text{IR}} = 40$



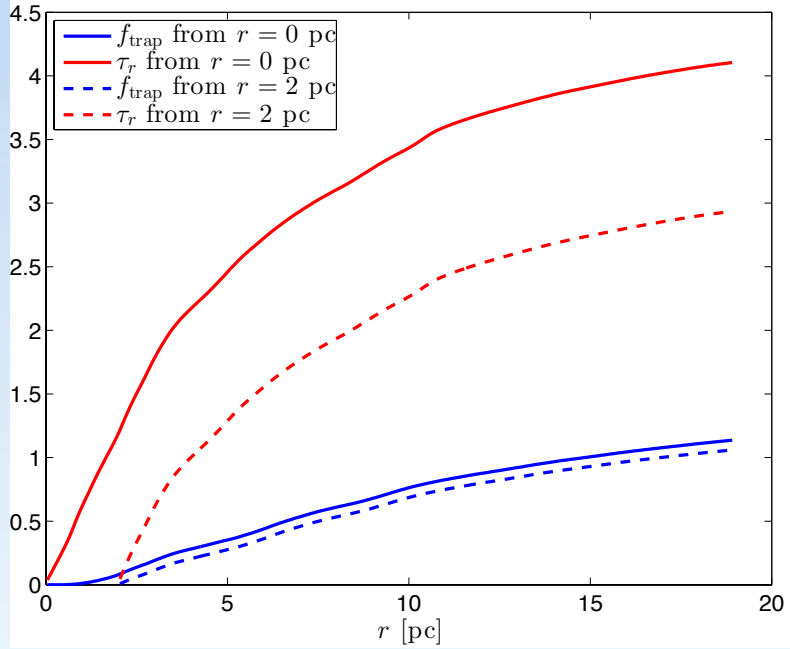
$t = 3.0t_{\text{ff}}, R_{\text{GMC}} = 10.00 \text{ pc}, M_{\text{GMC}} = 1.00e + 06M_{\odot}, \kappa_{\text{IR}} = 40$



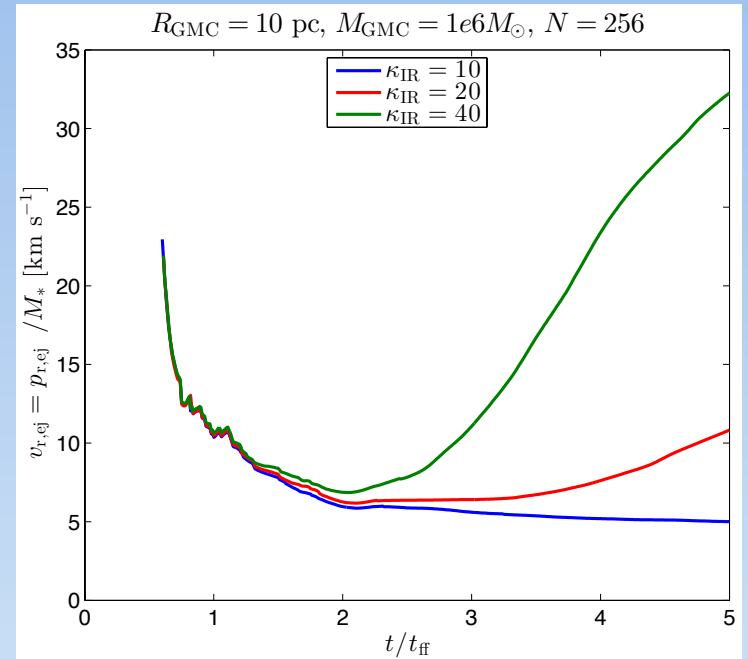
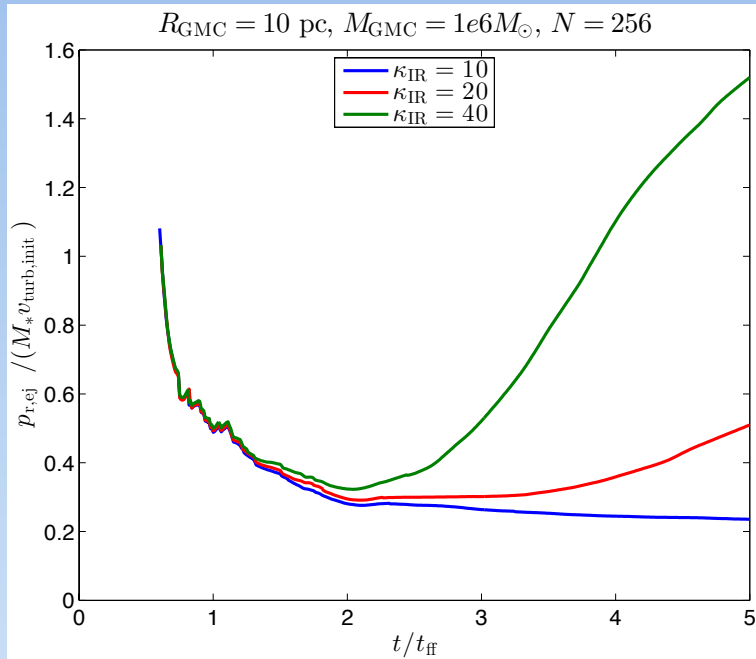
$t = 3.0t_{\text{ff}}, R_{\text{GMC}} = 10.00 \text{ pc}, M_{\text{GMC}} = 1.00e + 06M_{\odot}, \kappa_{\text{IR}} = 40$



$t = 3.0t_{\text{ff}}, R_{\text{GMC}} = 10.00 \text{ pc}, M_{\text{GMC}} = 1.00e + 06M_{\odot}, \kappa_{\text{IR}} = 40$



Momentum injection to ISM

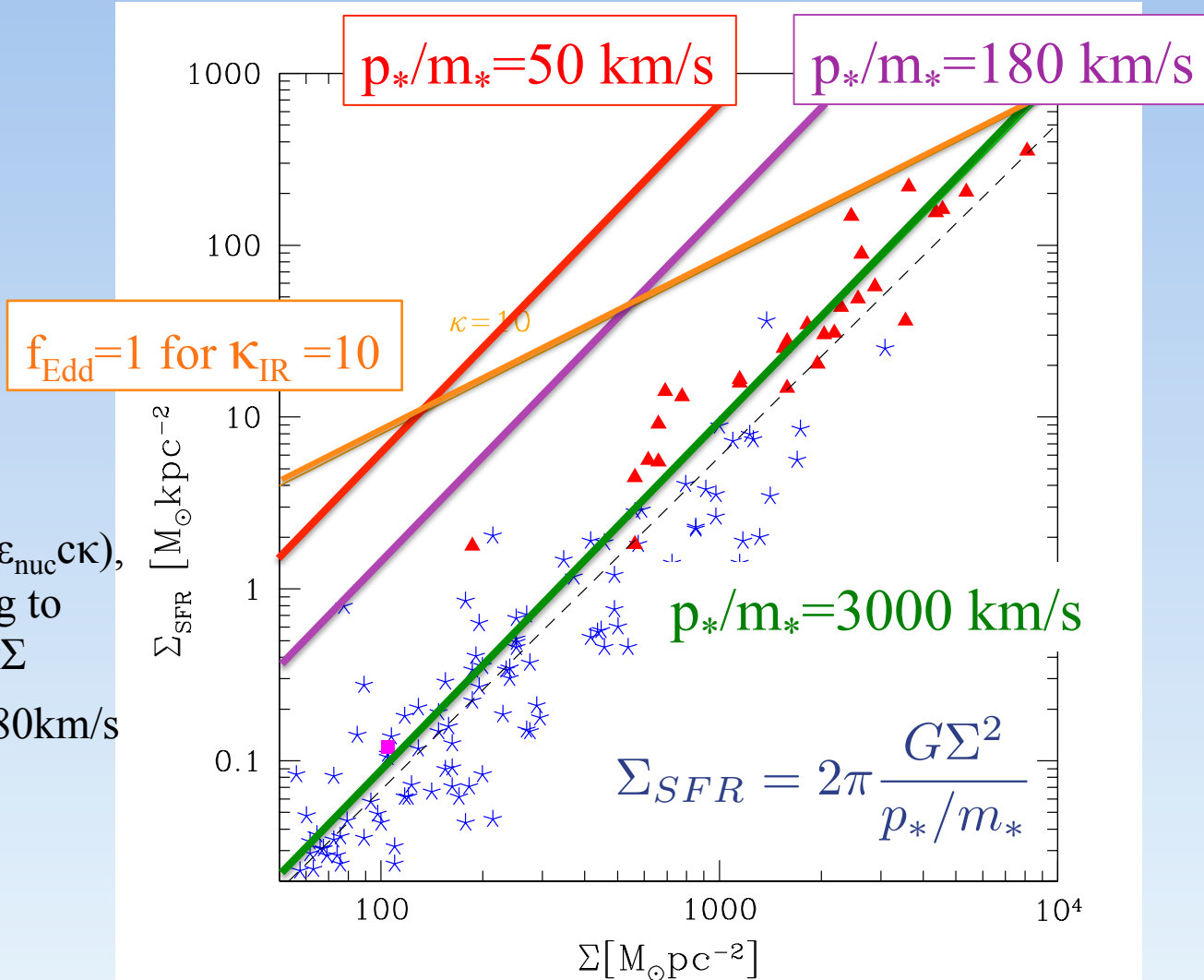


Momentum/mass given to the ISM:

$$p_*/m_* = p_{ej}/M_* \sim v_{\text{cloud}} < 100 \text{ km/s}$$

Radiation momentum feedback and large-scale SFRs

$f_{\text{Edd}}=1$:
 $\Sigma_{\text{SFR}}=2\pi G\Sigma/(\epsilon_{\text{nuc}}c\kappa)$,
 corresponding to
 $p_*/m_*=\epsilon_{\text{nuc}}c\kappa\Sigma$
 $=\tau_{\text{IR}}\times 180\text{km/s}$



Summary

- System in force balance and driving/dissipation balance for turbulence driven by feedback has

$$v \sim \varepsilon_{ff} \frac{p_*}{m_*} \quad L \sim \frac{GM_{tot}}{v^2} \quad \dot{M}_* \sim \frac{GM_{tot}M}{L^2 p_*/m_*}$$

- For gas-dominated disk system (starburst)

$$\Sigma_{SFR} \sim \frac{G\Sigma^2}{p_*/m_*}, \text{ or more generally: } \Sigma_{SFR} \sim \frac{\Sigma g_z}{p_*/m_*}$$

- Numerical simulations agree with simple theory and match observations of galactic Σ_{SFR} *provided p_*/m_* is large*
- Simulations of turbulent, star-forming clouds with IR radiation show that $\varepsilon = M_*/M_{cloud}$ is large (order-unity) and p_*/m_* is small (< 100 km/s) for realistic κ and v_{cloud}
- Further detailed studies are needed to quantify p_*/m_* from other sources (SNe with realistic ISM model, CR with realistic coupling,...) and connect to galactic winds